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FINITE ELEMENT SOLUTION OF A THREE-  
DIMENSIONAL NONLINEAR REACTOR  
DYNAMICS PROBLEM WITH FEEDBACK

Eulogio Conception Bermudes



# NAVAL POSTGRADUATE SCHOOL

## Monterey, California



# THESIS

FINITE ELEMENT SOLUTION OF A THREE-  
DIMENSIONAL NONLINEAR REACTOR  
DYNAMICS PROBLEM WITH FEEDBACK

by

Eulogio Conception Bermudes

December 1976

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DYNAMICS PROBLEM WITH FEEDBACK

by

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Submitted in partial fulfillment of the  
requirements for the degrees of

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## ABSTRACT

This work examines the three-dimensional dynamic response of a nonlinear fast reactor with temperature-dependent feedback and delayed neutrons when subjected to uniform and local disturbances. The finite element method was employed to reduce the partial differential reactor equation to a system of ordinary differential equations which can be numerically integrated. A program for the numerical solution of large sparse systems of stiff differential equations developed by Franke and based on Gear's method solved the reduced neutron dynamics equation. Although a study of convergence by refining element mesh sizes was not carried out, the crude finite element mesh utilized yielded the correct trend of neutron dynamic behavior.



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## I. INTRODUCTION

The nuclear reactor with delayed neutrons and temperature-dependent feedback is a nonlinear system, whose response to uniform and local disturbances differs greatly from that of a linear reactor. The prompt temperature feedback model and one average group of delayed neutrons were incorporated in this work. Practically all neutron dynamics analysis today deals with two-dimensional geometry, implying a symmetry in neutron dynamics behavior during transient [1,2]. This symmetry assumption, however, could be unrealistic in the safety analysis of nuclear reactors. A unique feature of this work is the consideration of the three-dimensional dependence of the neutron flux. No symmetry assumptions were imposed on the problem. This work investigated neutron dynamics under uniform and local disturbances in the core. A uniform initial flux throughout the interior of the reactor was imposed. The finite element method (FEM) was employed to solve this nonlinear initial-boundary value problem. The FEM is quite effective in handling discontinuous forcing functions thereby making it particularly suited for examining the effects of localized perturbations and space-dependent feedbacks.



## II. THE NUCLEAR REACTOR WITH TEMPERATURE DEPENDENT FEEDBACK AND DELAYED NEUTRONS

### A. THE PHYSICAL SYSTEM

The system under consideration is a fast reactor of cylindrical geometry that is composed of two different regions. The reactor core or region I is cylindrical in shape and is fueled by U-235. Region II or the reflector region completely surrounds the core and is composed of U-238. Both regions were assumed to be homogeneous. Table I lists the physical properties and geometry for each region. A schematic of the fast reactor geometry is shown in figure 1.

Temperature and delayed neutron effects were taken into account. Also, a one-velocity or one-group model was assumed, thereby making the velocity independent of spatial or temporal effects. The delayed neutrons were considered only in region I since region II was assumed to be a non-multiplying medium. The temperature effects were also assumed to be only in the core region.

In general form, the one-velocity neutron diffusion equation is [3]

$$\frac{\partial N(\vec{r}, t)}{\partial t} = vD\nabla^2 N - \Sigma_a vN + S(\vec{r}, t) \quad (1)$$

where  $\vec{r} = (x, y, z)$

$D$  = neutron diffusion coefficient

$\Sigma_a$  = macroscopic absorption cross-section

$N(\vec{r}, t)dV$  = number of neutrons in a volume element  $dV$  at a point  $\vec{r}$  at time  $t$

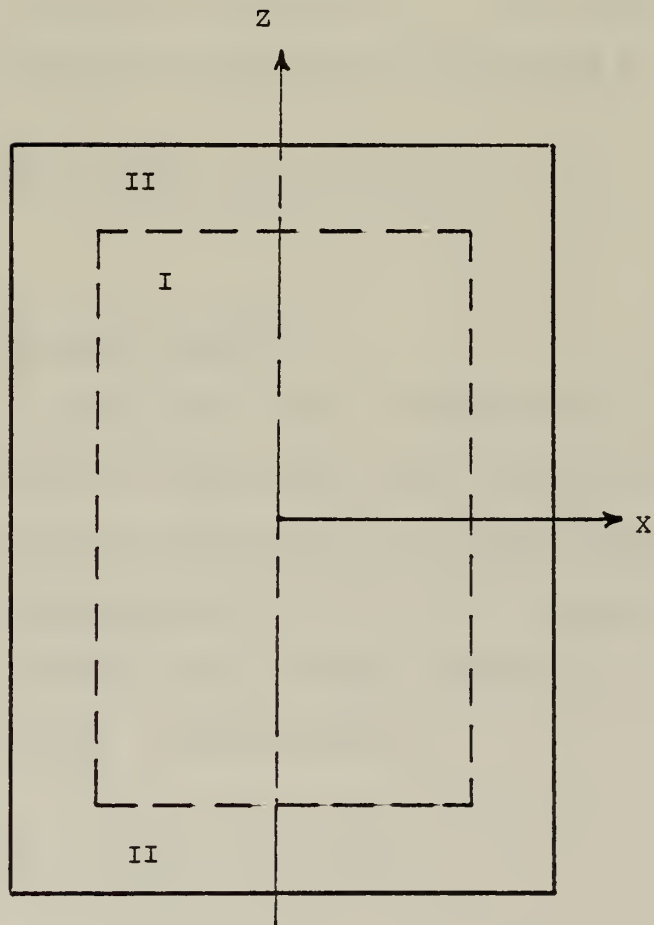
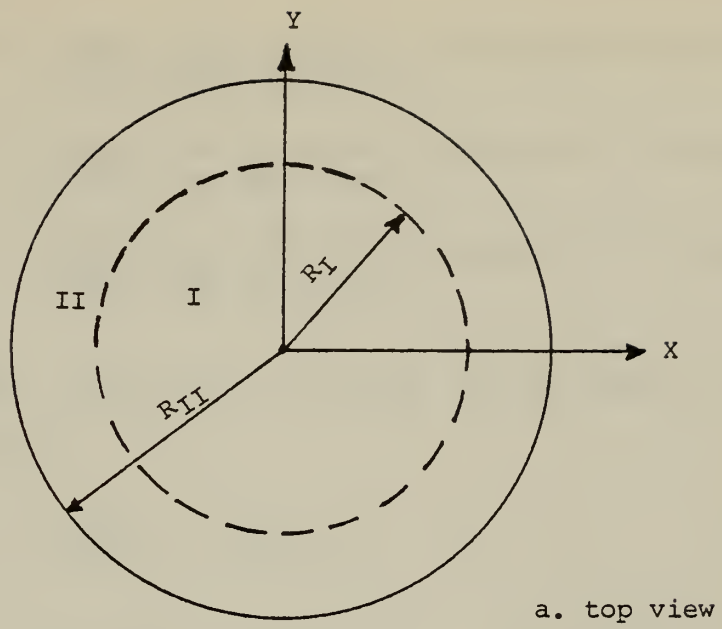


TABLE I  
Physical Constants

<u>SYMBOL</u>	<u>DEFINITION</u>	<u>VALUE</u>
$R_I$	radius of Region I	60 cm
$R_{II}$	total reactor radius	90 cm
$H_I$	height of Region I	160 cm
$H_{II}$	total reactor height	220 cm
$v$	neutron speed	$4.8 \times 10^7$ cm/sec
$D_I$	core neutron diffusion coefficient	0.913 cm
$D_{II}$	reflector neutron diffusion coefficient	1.200 cm
$\Sigma_{a_I}$	core neutron absorption cross section	$0.01401 \text{ cm}^{-1}$
$\Sigma_{a_{II}}$	reflector neutron absorption cross section	$0.0040 \text{ cm}^{-1}$
$\nu$	number of neutrons per fission	2.54
$\Sigma_f^*$	critical fission cross section	$0.005736 \text{ cm}^{-1}$
$\beta$	delayed neutron fraction; dollar reactivity	0.00642
$\epsilon$	fission energy	$7.652 \times 10^{-12}$ cal/fission
$\bar{h}(\frac{A}{V})$	modified convection heat transfer coefficient	0.0632 cal/(cm <sup>2</sup> sec °C)
$\alpha$	temperature coefficient	-0.004/°C
$\bar{\lambda}$	abundance-weighted mean decay constant	$0.4350 \text{ sec}^{-1}$







I - core  
II - reflector

Figure 1. Schematic of cylindrical reactor.



$vD\nabla^2 N dV$  = number of neutrons diffusing into  $dV$  per unit time at time  $t$

$\Sigma_a vN dV$  = number of neutrons absorbed in  $dV$  per unit time at time  $t$

$S(\bar{r},t)dV$  = number of neutrons produced in  $dV$  per unit time at time  $t$

The neutron number density is related to the neutron flux by the expression [4]

$$\phi(\bar{r},t) = vN(\bar{r},t) \quad (2)$$

where  $\phi(\bar{r},t)$  is the flux at time  $t$ . The neutron diffusion equation in terms of the flux is depicted by

$$\frac{1}{v} \frac{\partial \phi(\bar{r},t)}{\partial t} = D\nabla^2 \phi - \Sigma_a \phi + S(\bar{r},t) \quad (3)$$

#### B. PROMPT AND DELAYED NEUTRONS

The source or production term in equation (3) is composed of the contributions of the prompt and delayed neutrons. The majority of the fission neutrons are prompt neutrons that appear almost instantaneously (within  $10^{-7}$  second) on fission. Assuming a fast neutron non-leakage probability of unity, the prompt neutron source is described by

$$S_p(\bar{r},t) = (1-\beta)K_\infty(\bar{r},t) \phi(\bar{r},t) \quad (4)$$

where

$K_\infty(\bar{r},t)$  = infinite medium multiplication factor

$\beta$  = total fraction of delayed neutrons



The fraction of delayed neutrons is very small (note that  $\beta = 0.00642$  from Table I). However, they have a very significant effect on the reactivity because their mean lifetimes are long. Without delayed neutrons, reactor control would not be possible. These delayed neutrons are born in the decay by neutron emission of nuclei produced following the  $\beta$ -decay of certain fission fragments. For example, the  $\beta$ -decay of the fission fragment  $\text{Br}^{87}$  leads to  $\text{Kr}^{86}$  plus a neutron. Nuclei such as  $\text{Br}^{87}$  whose production in fission eventually leads to the emission of a delayed neutron are called delayed neutron precursors [4].

There are six main groups of delayed neutrons. Each group is classified according to its decay constant. The delayed neutron source term is portrayed by

$$S_d(\vec{r}, t) = \sum_{i=1}^6 C_i(\vec{r}, t) \lambda_i \quad (5)$$

where

$\lambda_i$  = decay constant of the  $i^{\text{th}}$  group

$C_i(\vec{r}, t)$  = density of the  $i^{\text{th}}$  precursor

Assuming that the fission fragments do not migrate appreciable distances and assuming a non-circulating fuel reactor [3], the precursor density is delineated by

$$\frac{\partial C_i(\vec{r}, t)}{\partial t} = \beta_i K_{\infty} \Sigma_a \phi - \lambda_i C_i \quad (6)$$

where  $\beta_i$  is the fraction of delayed neutrons of the  $i^{\text{th}}$



group. The solution to the precursor equation is in terms of a time integral expressed by

$$C_i(\bar{r}, t) = \beta_i \Sigma_a \int_0^t e^{-\lambda_i(t-t')} K_\infty(\bar{r}, t) \phi(\bar{r}, t) dt' \quad (7)$$

Inserting equation (7) in equation (5) yields the delayed neutron production term as

$$S_d(\bar{r}, t) = \sum_{i=1}^6 \beta_i \lambda_i \Sigma_a \int_0^t e^{-\lambda_i(t-t')} K_\infty(\bar{r}, t) \phi(\bar{r}, t) dt' \quad (8)$$

For convenience the six main groups of delayed neutrons were considered as one group. This was accomplished by using the abundance-weighted mean decay constant defined by

$$\bar{\lambda} = \frac{1}{\beta} \sum_{i=1}^6 \beta_i \lambda_i \quad (9)$$

This is a reasonable approximation since many of the important phenomena of nuclear reactor dynamics can be characterized satisfactorily by combining all the emitters or precursors into one, two, or, at most, three effective groups [3]. Replacing  $\lambda_i$  with the abundance-weighted mean decay constant showed the delayed neutron source term to be

$$S_d(\bar{r}, t) = \beta \bar{\lambda} \Sigma_a \int_0^t e^{-\bar{\lambda}(t-t')} K_\infty(\bar{r}, t) \phi(\bar{r}, t) dt' \quad (10)$$

### C. DOPPLER EFFECT AND TEMPERATURE FEEDBACK MODEL

The Doppler effect in fast reactors is due to the temperature broadening of many closely spaced high-energy resonances in both the fission and parasitic-absorption cross-sections.





These resonances basically mean that as the temperature increases, the number of neutrons that are absorbed increases and, similarly, the number of fission events increases.

These nonproductive and productive processes compete in a complicated manner, and the net effect may be either an increase or decrease in reactivity [3]. For most fast reactors, the effect is negative, and it will be so assumed in this work.

The reactivity change with respect to the fuel temperature change is modeled by [2]

$$\frac{dK_{\infty}}{dT} = aT^{-3/2} + bT^{-1} + cT^{(m-1)} \quad (11)$$

where  $a$ ,  $b$ , and  $c$  are parameters determined from experimental or neutronic calculations,  $m$  is an integer, and  $T$  is the current fuel temperature. For most fast reactor cores (ceramic-fueled cores),  $TdK_{\infty}/dT$  is very close to being constant over a wide range of temperatures. Therefore, it is assumed that  $a$  and  $c$  are zero and the Doppler coefficient is expressed by

$$b = T \frac{dK_{\infty}}{dT} \quad (12)$$

The following initial conditions were used:

$$K_{\infty}(\bar{r}, 0^+) = K_{\infty}^0 \quad (12a)$$

$$T(\bar{r}, 0^+) = T_0 \quad (12b)$$



The reactivity model is then expressed by

$$K_{\infty}(\bar{r}, t) = K_{\infty}^{\circ} + b \ln \left( \frac{T}{T_0} \right) \quad (13)$$

where

- $K_{\infty}^{\circ}$  = multiplication factor at steady-state
- $\nu$  = number of neutrons produced per fission
- $T_0$  = original fuel temperatures
- $b$  = Doppler constant

The definition of  $K_{\infty}^{\circ}$  is

$$K_{\infty}^{\circ} = \frac{\nu \Sigma_f^*}{\Sigma_a} \quad (14)$$

where

- $\Sigma_f^*$  = critical fission cross-section
- $\Sigma_a^c$  = macroscopic absorption cross-section of the core

The rise in fuel temperature is depicted by

$$\theta(\bar{r}, t) = T(\bar{r}, t) - T_0(\bar{r}) \quad (15)$$

This is also described by the integral [5]

$$\theta(\bar{r}, t) = \int_0^t f(t-t') \psi(\bar{r}, t) dt' \quad (16)$$

where  $f(t-t')$  is the feedback kernel and is dependent upon the type of temperature feedback model used.  $\psi(\bar{r}, t)$  is the



rise in neutron flux above the steady state and is expressed by

$$\psi(\vec{r},t) = \phi(\vec{r},t) - \phi_0(r) \quad (17)$$

where  $\phi_0(r)$  is the neutron flux at steady state.

There are three types of temperature feedback models that can be considered here. The first is known as "Newton's feedback" which determines the reactor temperature by Newton's law of cooling. The second temperature feedback model is called the adiabatic feedback model. This represents the temperature for the loss of coolant case. The third model is the prompt temperature feedback model in which the fuel temperature follows the behavior of the neutron flux without delay [5,6]. The prompt feedback model was employed in this analysis. The feedback kernel for this temperature feedback model is

$$f(t-t') = \frac{K}{\gamma} \delta(t-t') \quad (18)$$

where  $K$  is an energy production operator with units of °C per unit flux and  $\gamma$ , a dimensionless quantity, is related to the mean time for heat transfer to coolant.

Inserting equation (18) in equation (16) and performing the integration produced a rise in temperature of

$$\theta(\vec{r},t) = \frac{K}{\gamma} \psi(\vec{r},t) \quad (19)$$



From equation (19) the temperature of the fuel is

$$T(\bar{r}, t) = \frac{K}{\gamma} \psi(\bar{r}, t) + T_0(\bar{r}) \quad (20)$$

The ratio of the current temperature to the steady-state temperature is therefore

$$\frac{T}{T_0} = \frac{K}{\gamma T_0} \psi + 1 \quad (21)$$

In this work  $T_0$  was considered to be constant throughout the core. Incorporating equation (21) into equation (13) gave the reactivity model of

$$K_\infty(\bar{r}, t) = K_\infty^0 + b \ln \left[ \frac{K}{\gamma T_0} \psi(\bar{r}, t) + 1 \right] \quad (22)$$

#### D. FIELD EQUATIONS

Before establishing the field equations it is desirable to express equation (3) in terms of the rise in flux. Using equation (17) in equation (3) and grouping terms yielded the following diffusion equation:

$$\begin{aligned} \frac{1}{v} \frac{\partial \psi}{\partial t} = & [D \nabla^2 \psi - \Sigma_a \psi + (1-\beta) K_\infty \Sigma_a \psi + \beta \bar{\lambda} \Sigma_a \int_0^t e^{-\bar{\lambda}(t-t')} K_\infty \psi dt'] \\ & + [D \nabla^2 \phi_0 - \Sigma_a \phi_0 + (1-\beta) K_\infty \Sigma_a \phi_0 + \beta \bar{\lambda} \Sigma_a \int_0^t e^{-\bar{\lambda}(t-t')} K_\infty \phi_0 dt'] \end{aligned} \quad (23)$$

The second bracketed term in equation (23) is identically equal to zero since it is the steady state portion of the





diffusion equation. The rise in neutron flux above its steady state value is therefore expressed, for the core, by

$$\begin{aligned} \frac{1}{v} \frac{\partial \psi}{\partial t} = & D \nabla^2 \psi - \Sigma_a \psi + (1-\beta) K_{\infty} \Sigma_a \psi \\ & + \beta \bar{\lambda} \Sigma_a \int_0^t e^{-\bar{\lambda}(t-t')} K_{\infty} \psi dt' \end{aligned} \quad (24)$$

and, for the reflector region, by

$$\frac{1}{v} \frac{\partial \psi}{\partial t} = D \nabla^2 \psi - \Sigma_a \psi \quad (25)$$

Inserting the reactivity model into equation (24) yielded

$$\begin{aligned} \frac{1}{v} \frac{\partial \psi}{\partial t} = & D \nabla^2 \psi - \Sigma_a \psi + (1-\beta) \Sigma_a K_{\infty}^o \psi \\ & + (1-\beta) b \Sigma_a \left\{ \ln \left[ \frac{K \psi}{\gamma T_o} + 1 \right] \right\} \psi + \beta \bar{\lambda} K_{\infty}^o \int_0^t e^{-\bar{\lambda}(t-t')} \psi dt' \\ & + \beta \bar{\lambda} \Sigma_a b \left\{ \int_0^t e^{-\bar{\lambda}(t-t')} \left[ \ln \left( \frac{K \psi}{\gamma T_o} + 1 \right) \right] \psi dt' \right\} \end{aligned} \quad (26)$$

For the reflector, the last four terms of equation (26) are zero. The effects of the temperature on the delayed neutrons were neglected in this work. The field equations can now be expressed, for the core, as

$$\begin{aligned} \frac{\partial \psi}{\partial t} - v D \nabla^2 \psi + [v \Sigma_a - v(1-\beta) v \Sigma_f] \psi \\ + [-(1-\beta) v \Sigma_a b] \left[ \ln \left( \frac{K \psi}{\gamma T_o} + 1 \right) \right] \psi \\ + [-\beta \bar{\lambda} v \Sigma_f] \left[ \int_0^t e^{-\bar{\lambda}(t-t')} \psi dt' \right] = 0 \end{aligned} \quad (27)$$



and, for the reflector, as

$$\frac{\partial \psi}{\partial t} - v D \nabla^2 \psi + v \Sigma_a \psi = 0 \quad (28)$$

The non-linear terms of the core field equation will be linearized accordingly. In more compact form, equations (27) and (28) became

$$\frac{\partial \psi}{\partial t} - c1 \nabla^2 \psi + c2 \psi + c4 \left[ \ln \left( \frac{K \psi}{\gamma T_0} + 1 \right) \right] \psi + c5 \left[ \int_0^t e^{-\bar{\lambda}(t-t')} \psi dt' \right] = 0 \quad (29)$$

and

$$\frac{\partial \psi}{\partial t} - c1 \nabla^2 \psi + c2 \psi = 0 \quad (30)$$

where the meanings of the coefficients  $c1$ ,  $c2$ , etc., are obvious for the core and reflector.

Equations (29) and (30) were subjected to the following conditions:

boundary condition:

$$\psi_R(\bar{r}_B, t) = 0 \quad (31)$$

where  $\bar{r}_B$  are coordinates of points on the outer surface of the reflector and the subscript  $R$  refers to the reflector

continuity of flux:

$$\psi_R(\bar{r}, t) = \psi_C(\bar{r}_I, t) \quad (32)$$

where  $\bar{r}_I$  are coordinates of points on the core-reflector interface and the subscript  $c$  refers to the core.



### III. FINITE ELEMENT

#### A. INTRODUCTION

The application of the finite elements of nonlinear continua has been mostly in the field of solid mechanics. Prior work [1] using the finite elements of nonlinear continua on a nonlinear reactor dynamics problem has been successful. The FEM in this work was utilized to reduce a nonlinear partial differential equation of the nuclear reactor to a system of nonlinear ordinary differential equations in time. The time integration was accomplished by using a computer program for the numerical solution of stiff differential equations developed by Franke [7]. In order to minimize computer storage requirements, an optimum compacting scheme (OCS), described by Ref. 8, was adopted.

The finite element models of operator equations are generally classified into three categories: a) the variational finite element models such as the Ritz method, b) the weighted residuals method such as the method of Galerkin, and c) the direct finite element models which are not based on functional minimization. From experience in structural mechanics, the most effective method for generating acceptable finite element models of nonlinear equations is the Galerkin method [1]. This work adopted the method of Galerkin in a finite element approximation over the spatial domain of the field equations.



## B. THE METHOD OF GALERKIN

The Galerkin method is a special case of the method of weighted residuals. It involves a rational choice of weighting function that is consistent with the type of finite element approximation considered. Indeed, the weighting functions chosen are the basis or shape functions employed in the finite element approximation. There are two favorable characteristics of the Galerkin method which makes it attractive. The first attribute is its amenability to integration by parts. This supplied the freedom of using a lower order finite element than might be otherwise possible. The second favorable characteristic of the Galerkin method is that the symmetric operators in the field equations transform into symmetric matrix operators. Both these attributes are attractive for computational purposes.

Consider the initial-boundary-value problem

$$\frac{\partial \psi}{\partial t} = \mathcal{L}\psi - f(\bar{r}, t) \quad (33)$$

where  $\mathcal{L}$  contains the nonlinear operators. According to the spatial finite element discretization, the solution of equation (33) is in the form of an union of an  $\bar{N}$ -term approximation given by

$$\psi(\bar{r}, t) \approx \tilde{\psi}(\bar{r}, t) = \bigcup_{j=1}^{\bar{N}} \psi_j(t) G_j(\bar{r}) , \quad j=1, 2, \dots, \bar{N} \quad (34)$$

where  $\bar{N}$  is the number of system degrees of freedom (i.e., number of coordinates), and  $G_j(\bar{r})$  are the system or global basis functions which span the space of the approximate





solution  $\tilde{\psi}(\bar{r}, t)$  [1]. The Einstein summation is used. The global basic functions are "pyramid functions", each of which has a prescribed functional description over a sub-domain of the system and is zero elsewhere [9]. The unknown coordinate functions  $\psi_j(t)$  are the time-dependent magnitudes of the approximated flux  $\tilde{\psi}(\bar{r}, t)$  and/or its derivatives at discrete nodal points [10].

The residual function,  $R(\bar{r}, t)$ , is defined such that it is identical to zero when  $\tilde{\psi}(\bar{r}, t)$  is equal to the exact solution. The residual function is expressed by

$$R(\bar{r}, t) = \frac{\partial \tilde{\psi}}{\partial t} - \mathcal{L} \tilde{\psi} - f \quad (35)$$

The Galerkin orthogonality condition (using the basis functions as weight functions), when applied to the residual function, requires that

$$\int_{Vol} G_I R(\bar{r}, t) dVol = 0, \quad I=1, 2, \dots, \bar{N} \quad (36)$$

From the field equations, the residual function for the core is

$$\begin{aligned} R(\bar{r}, t) = & \frac{\partial \tilde{\psi}}{\partial t} - c_1 \nabla^2 \tilde{\psi} + c_2 \tilde{\psi} + c_4 \left[ \ln \left( \frac{K \tilde{\psi}}{\gamma T_0} + 1 \right) \right] \tilde{\psi} \\ & + c_5 \left[ \int_0^t e^{-\bar{\lambda}(t-t')} \tilde{\psi} dt' \right] \end{aligned} \quad (37)$$

and for the reflector

$$R(\bar{r}, t) = \frac{\partial \tilde{\psi}}{\partial t} - c_1 \nabla^2 \tilde{\psi} + c_2 \tilde{\psi} \quad (38)$$



Using equation (34) and applying Galerkin's orthogonality condition produced a core equation of

$$\begin{aligned}
 \int_{Vol} G_I \{ G_J \dot{\psi}_J(t) - \nabla^2 G_J(c1 \cdot \psi)_J + G_J(c2 \cdot \psi)_J \\
 + [\ln(\frac{K}{\gamma T_0} G_K \psi_K + 1)] G_J(c4 \cdot \psi)_J \\
 + [\int_0^t e^{-\bar{\lambda}(t-t')} (c5 \cdot \psi)_J dt'] G_J \} = 0 \quad (39)
 \end{aligned}$$

where  $I = 1, 2, \dots, \bar{N}$   
 $J = 1, 2, \dots, \bar{N}$   
 $K = 1, 2, \dots, \bar{N}$

The reflector has a similar equation without the nonlinearity and is expressed by

$$\int_{Vol} G_I \{ G_J \dot{\psi}_J(t) - \nabla^2 G_J(c1 \cdot \psi)_J + G_J(c2 \cdot \psi)_J \} dVol = 0 \quad (40)$$

### C. THE ELEMENT

A three-dimensional quadratic isoparametric element was employed in this work. The parent element is a triangular prism or solid wedge with straight sides. The element shape functions are expressed in terms of area coordinates in the plane of the triangle and by an isoparametric coordinate



along the prism axis. Figure 2 shows the parent element. This element was chosen because of the ease with which it fits the cylindrical structure when it is transformed into a curved element. This type of element has been used before as filler elements [11].

The area coordinates are defined by area ratios. Consider the triangle shown in figure 3. An arbitrary point P within the triangle defines three subareas designated by  $A_1$ ,  $A_2$ , and  $A_3$ . The ratio of each of the subareas to the total area is known as an area coordinate. In equation form, the area coordinates are

$$L_1 = A_1/A \quad (41a)$$

$$L_2 = A_2/A \quad (41b)$$

$$L_3 = A_3/A \quad (41c)$$

where  $A$  is total area of the triangle.  $L_1$ ,  $L_2$ , and  $L_3$  are the natural coordinates for a triangle. The requirement that the sum of the subareas be equal to the total area is obviously satisfied by the identity

$$L_1 + L_2 + L_3 = 1 \quad (42)$$

In the plane of the triangle, only two of the area coordinates are independent.

The isoparametric coordinates are best visualized by considering the rectangular prism shown in figure 4. Isoparametric coordinates are normalized coordinates such that their values on the faces of the rectangle are  $\pm 1$ . The  $\xi\eta\zeta$  axes are in general not orthogonal. They are orthogonal



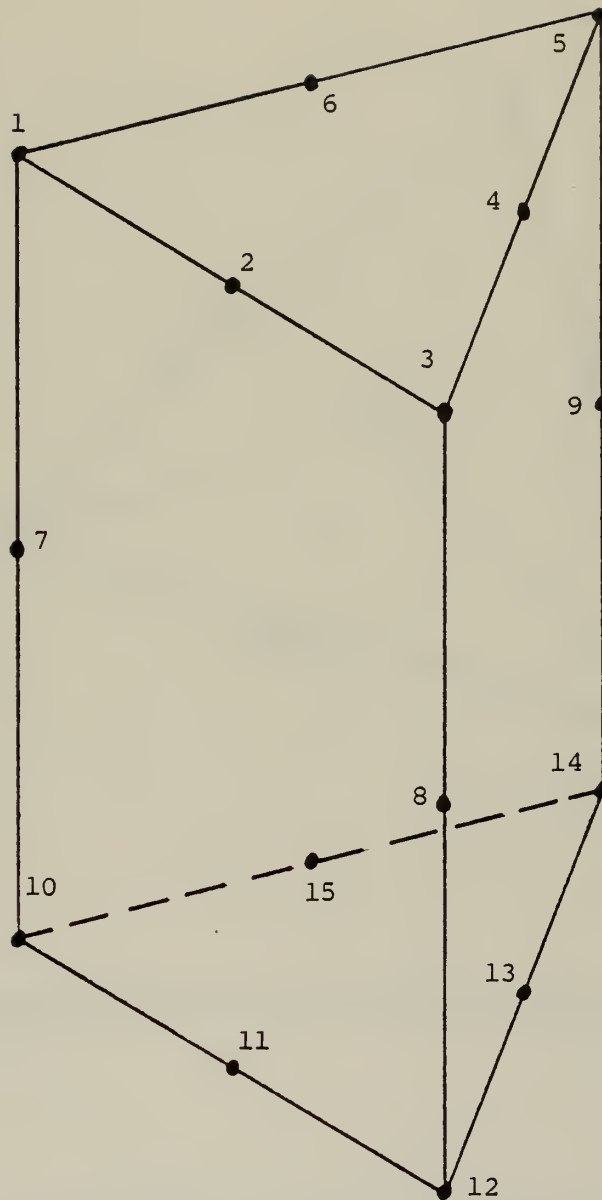


Figure 2. Quadratic triangular prism parent element





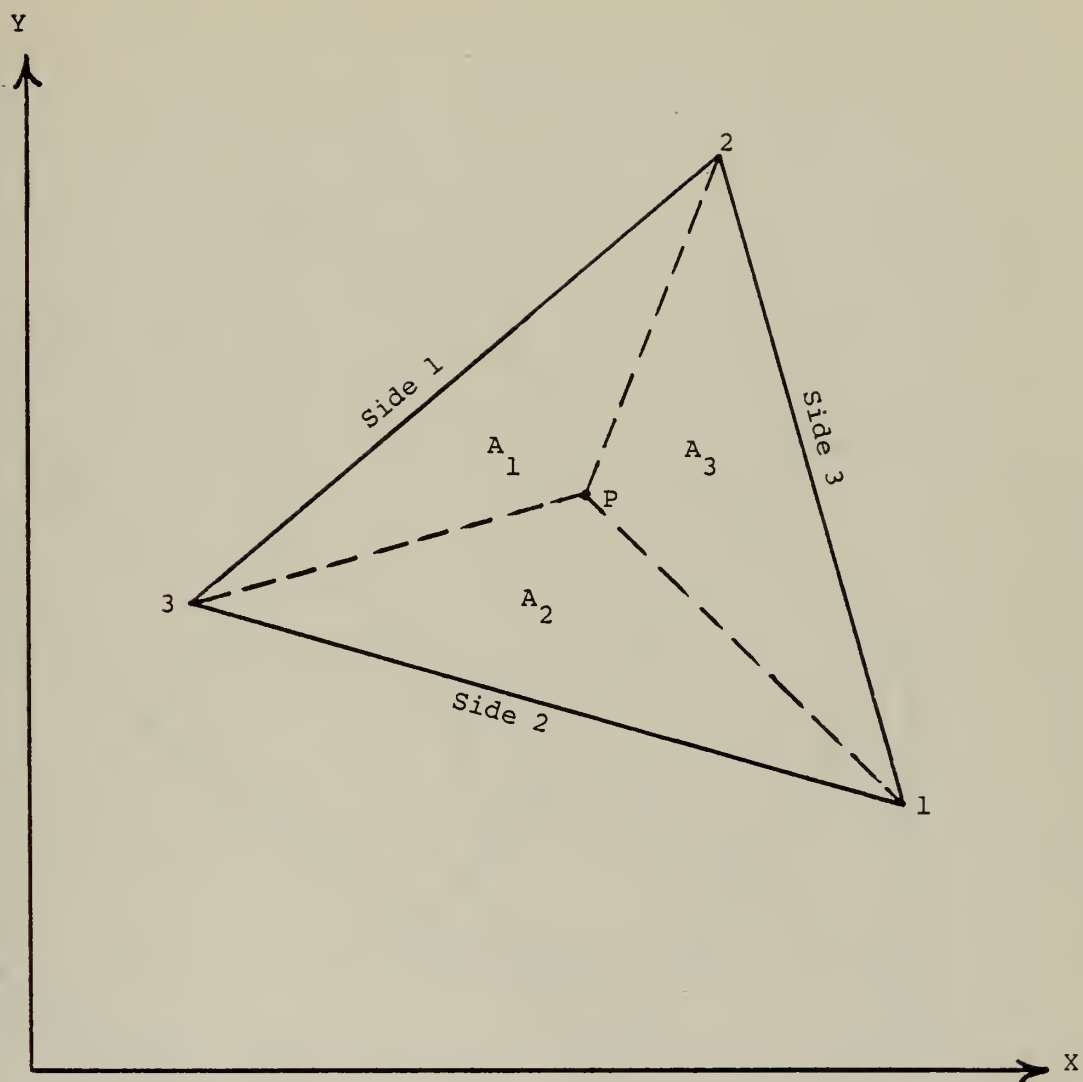


Figure 3. Definition of area coordinates



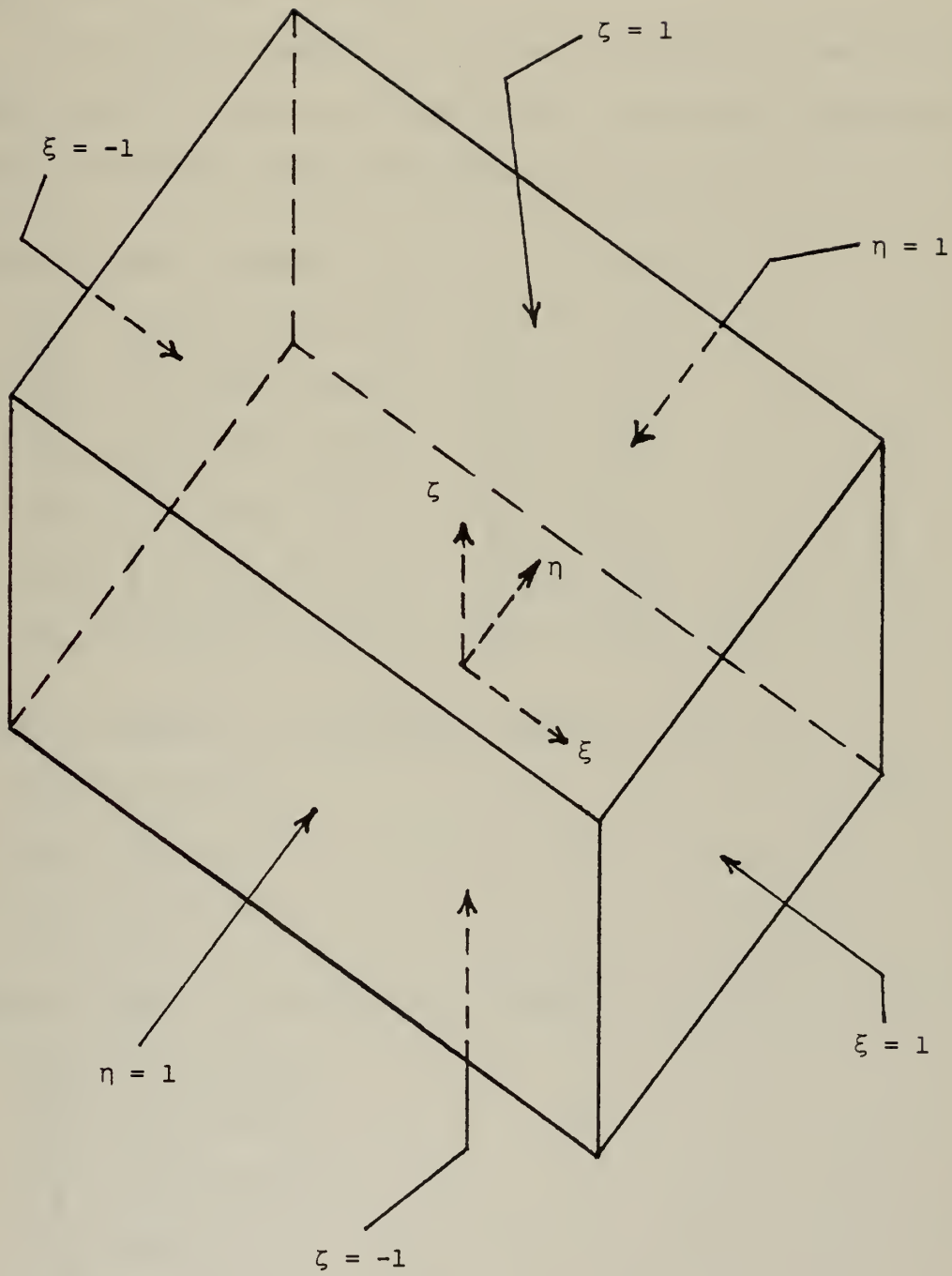


Figure 4. Isoparametric coordinates



only in the special case of a rectangular prism element [12]. The element basis functions use the  $\zeta$  coordinate in the prism axis and the  $L_1$ ,  $L_2$ , and  $L_3$  coordinates in the plane of the triangle.

The parent element, as shown in figure 2, has 15 local nodal points around its periphery. As such, there are 15 basis functions which are given below [11]:

Corner nodes: (nodes 1, 3, 5, 10, 12, 14)

$$N_1 = \frac{1}{2} L_1 (2L_1 - 1)(1 + \zeta) - \frac{1}{2} L_1 (1 - \zeta^2)$$

$$N_3 = \frac{1}{2} L_2 (2L_2 - 1)(1 + \zeta) - \frac{1}{2} L_2 (1 - \zeta^2)$$

$$N_5 = \frac{1}{2} L_3 (2L_3 - 1)(1 + \zeta) - \frac{1}{2} L_3 (1 - \zeta^2)$$

$$N_{10} = \frac{1}{2} L_1 (2L_1 - 1)(1 - \zeta) - \frac{1}{2} L_1 (1 - \zeta^2)$$

$$N_{12} = \frac{1}{2} L_2 (2L_2 - 1)(1 - \zeta) - \frac{1}{2} L_2 (1 - \zeta^2)$$

$$N_{14} = \frac{1}{2} L_3 (2L_3 - 1)(1 - \zeta) - \frac{1}{2} L_3 (1 - \zeta^2)$$

Midside nodes of rectangles: (nodes 7, 8, 9)

$$N_7 = L_1 (1 - \zeta^2)$$

$$N_8 = L_2 (1 - \zeta^2)$$

$$N_9 = L_3 (1 - \zeta^2)$$

Midside nodes of triangles: (nodes 2, 4, 6, 11, 13, 15)

$$N_2 = 2L_1 L_2 (1 + \zeta)$$

$$N_4 = 2L_2 L_3 (1 + \zeta)$$

$$N_6 = 2L_3 L_1 (1 + \zeta)$$

$$N_{11} = 2L_1 L_2 (1 - \zeta)$$

$$N_{13} = 2L_2 L_3 (1 - \zeta)$$

$$N_{15} = 2L_3 L_1 (1 - \zeta)$$



The coordinates of each local node, in terms of  $L_1$ ,  $L_2$ ,  $L_3$ , and  $\zeta$  are listed in table II. These element basis functions  $\langle N \rangle$  define the geometry of the element. Note that they satisfy the relationship

$$N_i = \begin{cases} 1, & \text{at node } i \\ 0, & \text{at other nodes} \end{cases} \quad (43)$$

On the element level, the variation of the unknown function  $\psi$  is approximated by

$$\tilde{\psi}^e = \langle N' \rangle \{ \psi \}^e \quad (44)$$

where  $\langle N' \rangle$  = row vector of element shape functions  
 $\{ \psi \}^e$  = column vector of time-dependent nodal values of  $\tilde{\psi}^e$

To satisfy continuity requirements, the shape functions  $\langle N' \rangle$  have to be such that the continuity of the unknown function  $\psi$  is preserved in the parent coordinates [11].

The shape functions  $\langle N \rangle$  which characterize the element geometry and the shape functions  $\langle N' \rangle$  which describe the unknown function do not necessarily have to be the same. There is no requirement that the nodal values be associated with the same nodes which were used to define the element geometry, though in practice it is often the case. Consider for example, the illustrations in figure 5. If the nodes defining the element geometry and the nodes defining the



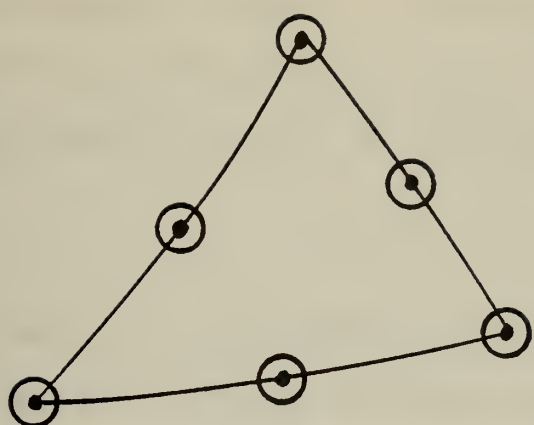


TABLE II

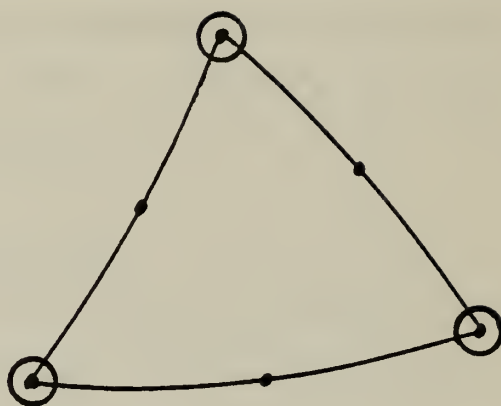
Coordinates of Local Nodal Points

Local Node	$L_1$	$L_2$	$L_3$	5
1	1	0	0	1
2	$\frac{1}{2}$	$\frac{1}{2}$	0	1
3	0	1	0	1
4	0	$\frac{1}{2}$	$\frac{1}{2}$	1
5	0	0	1	1
6	$\frac{1}{2}$	0	$\frac{1}{2}$	1
7	1	0	0	0
8	0	1	0	0
9	0	0	1	0
10	1	0	0	-1
11	$\frac{1}{2}$	$\frac{1}{2}$	0	-1
12	0	1	0	-1
13	0	$\frac{1}{2}$	$\frac{1}{2}$	-1
14	0	0	1	-1
15	$\frac{1}{2}$	0	$\frac{1}{2}$	-1

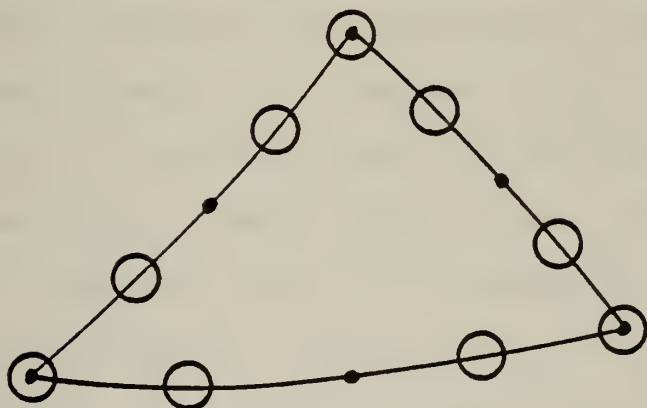




(a) Isoparametric



(b) Super-parametric



(c) Sub-parametric

- - geometry nodes
- - variable function nodes

Figure 5. Element classification



unknown function are identical, the element is known as an isoparametric element. This means that the shape functions describing the geometry and the shape functions describing the unknown function  $\psi$  are equal, or

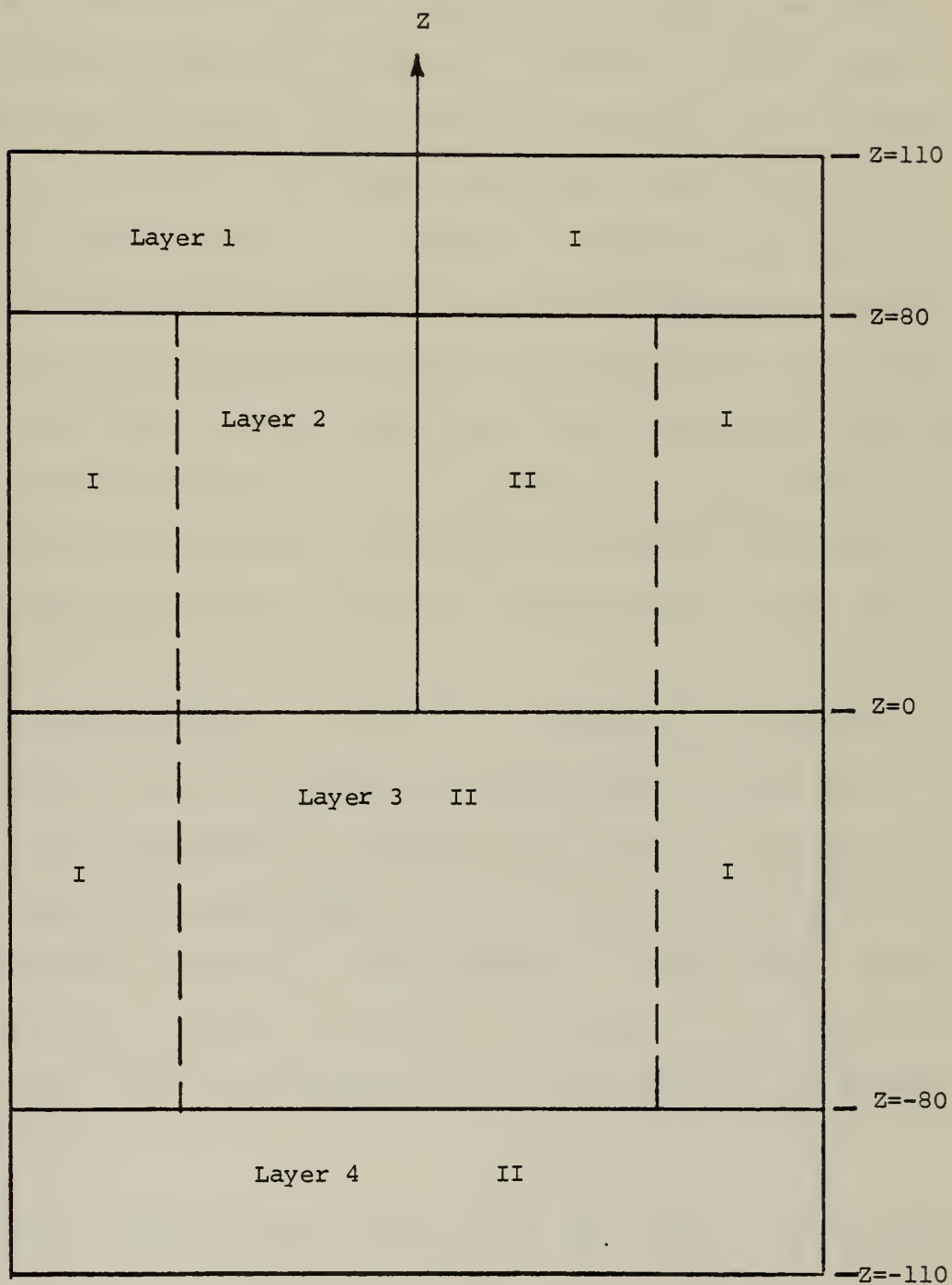
$$\langle N' \rangle = \langle N \rangle \quad (45)$$

If there are more nodes defining the geometry than nodes defining the variable function, the element is called a superparametric element. Using more nodal points to define the unknown function than to describe the geometry of the element results in a subparametric element [11]. This work utilized the isoparametric element classification.

#### D. DIVISION OF THE SYSTEM INTO ELEMENTS

In three-dimensional space, the division of the system into discrete elements is difficult to visualize. It is virtually impossible to show every nodal point of the system in one schematic. To present a clear view of the discretized domain, a "layer" approach was adopted. The first finite element grid or mesh employed here consisted of 128 elements. Under this grid (mesh I), the reactor was divided into four layers as shown in figure 6. Each layer was composed of 32 elements. The first and fourth layers were each 30 cm in height and each contained entirely reflector elements. The second and third layers were each 80 cm in height and together they encompassed the entire core plus the remaining reflector elements. Each layer, in turn, was partitioned into three horizontal (xy) planes. The top plane included all the global





I - core      II - reflector

Figure 6. Layers of mesh I





or system nodes corresponding to local or element nodes 1 through 6. The middle plane contained all the system nodes corresponding to the local nodes 7, 8, and 9. System nodes corresponding to element nodal points 10 through 15 comprised the bottom plane. To fix ideas, the three nodal planes of the first layer of mesh I are shown in figures 7, 8, and 9.

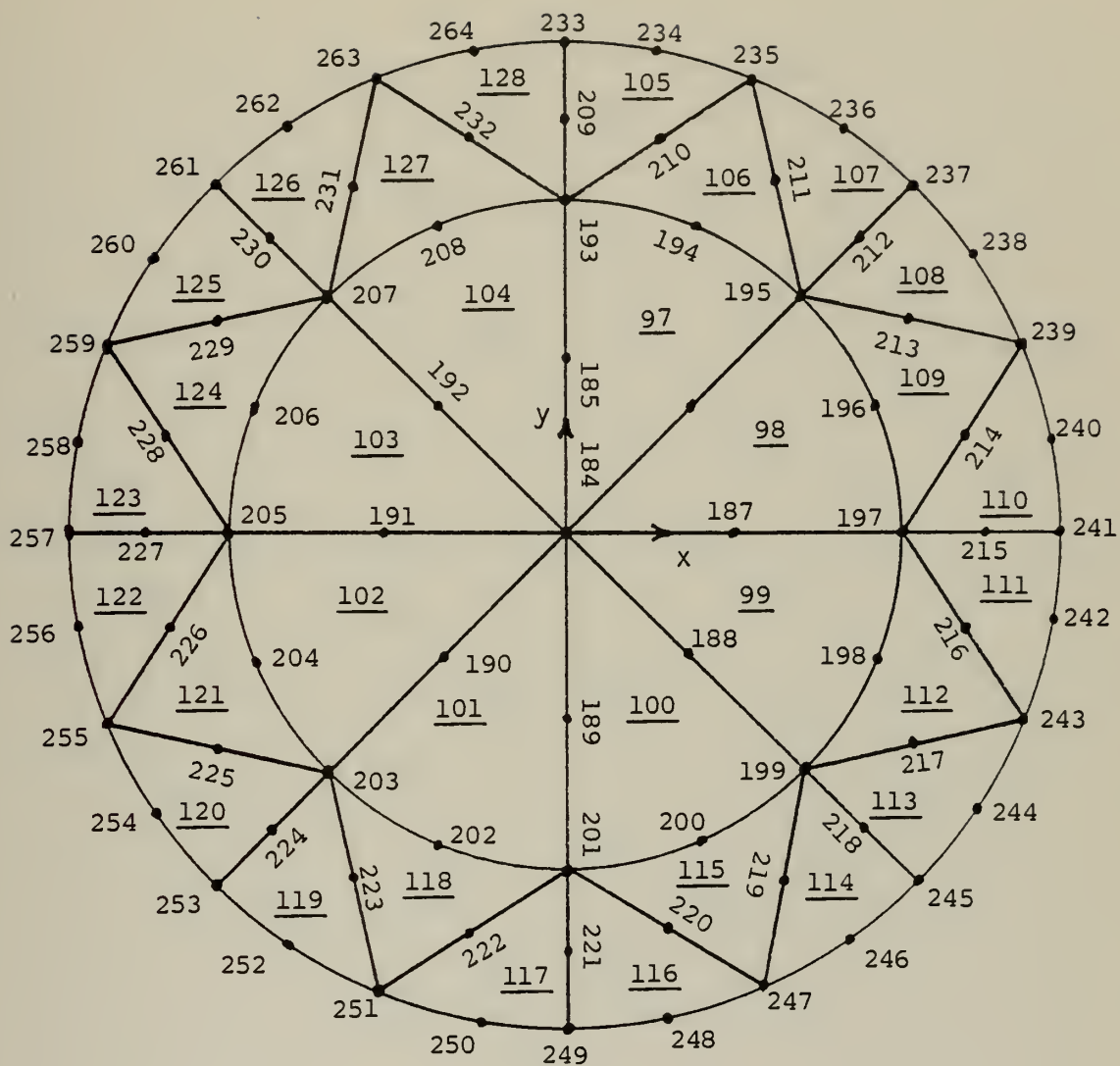
In this work there was only one curved side which was in the plane of the triangle, as shown in figure 10. The first, seventh, and tenth element nodes were each arbitrarily assigned to have as an opposite side the curved side of the triangle. The remaining local nodes in each of the respective planes of the element were then numbered consecutively in the counter-clockwise direction.

At this point it is appropriate to define connectivity. The connectivity of an element is a row vector array that relates the local nodes to system nodal points. The connectivity lists the system nodal points that are "connected" to form the element domain in the sequence of local nodal numbering. Thus, the connectivity matrix for mesh I is a matrix of size  $128 \times 15$ . To illustrate, the connectivity of element number 106 is

<235, 210, 193, 194, 195, 211, 266, 176, 177, 283, 152, 10, 11, 12, 53>

A second finite element mesh (mesh II) consisting of 192 elements was developed. It contained the same number of elements per layer as mesh I. However, mesh II has six layers. The second and third layers of mesh I were each divided in

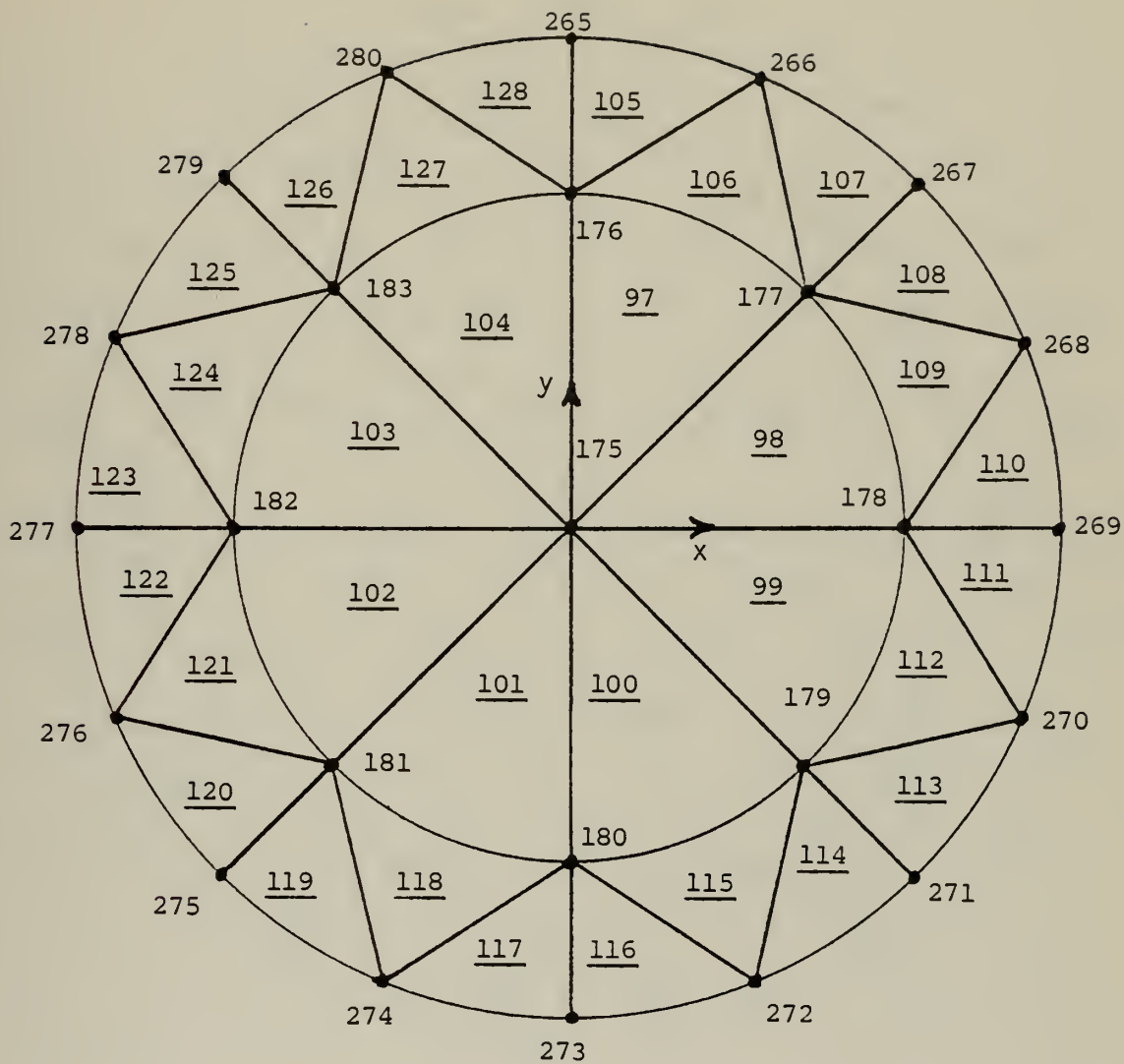




Number - element number

Figure 7. Top nodal plane of the first layer of mesh I,  
 $z = 110$  cm

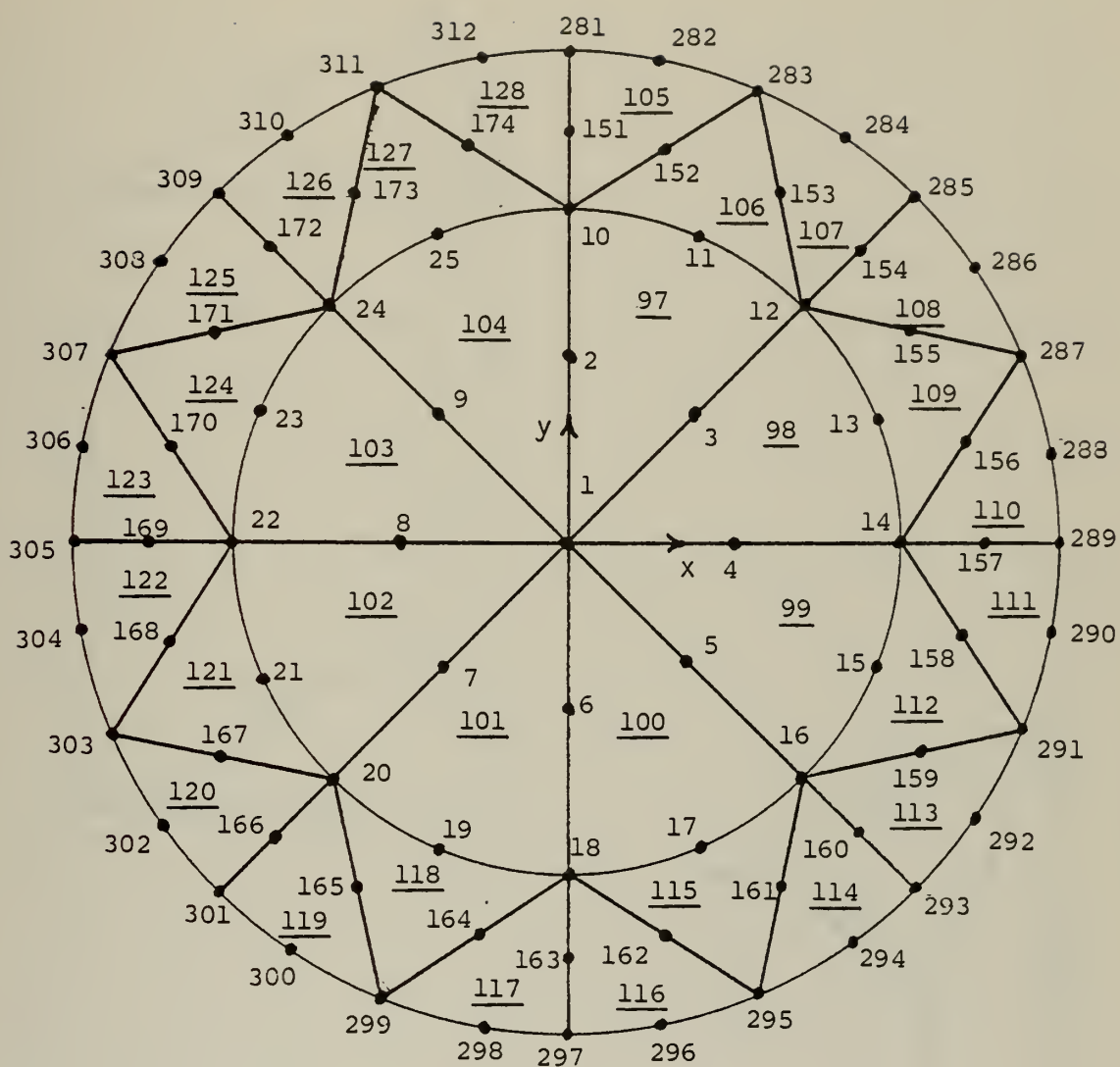




Number - element number

Figure 8. Middle nodal plane of the first layer of mesh I,  $z = 95$  cm





Number - element number

Figure 9. Bottom nodal plane of the first layer of mesh I,  $z = 80$  cm





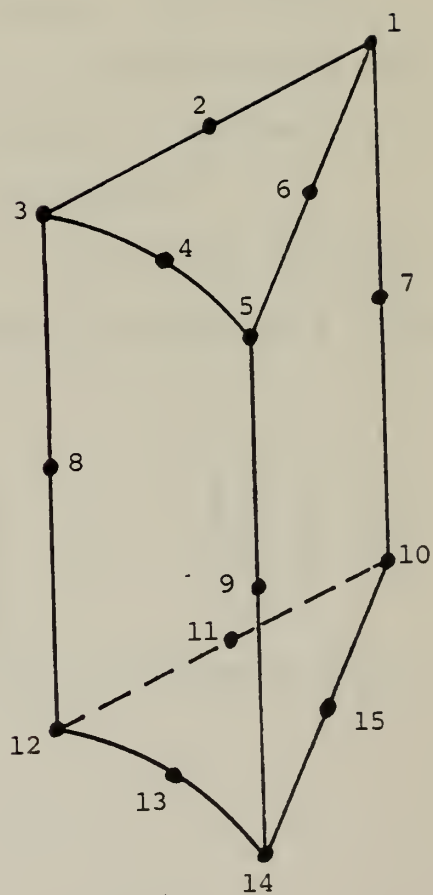
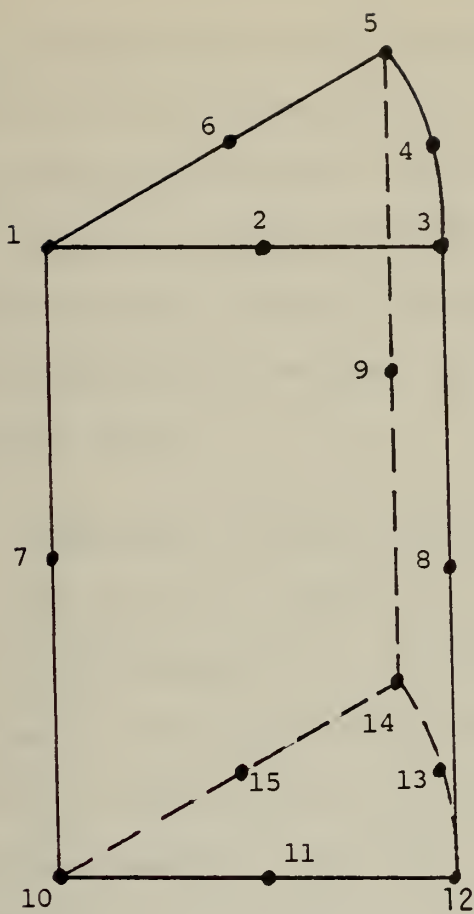


Figure 10. Local nodal numbering of curved elements



half, thus forming the two additional layers of mesh II. The connectivity matrix and the coordinates of each global node for mesh I are given in Appendix A. Appendix B lists the connectivity matrix and nodal coordinates of mesh II. The reader can construct the different layers and elements without much difficulty by using the nodal coordinates and the connectivity matrix. A mesh generator was not utilized in this work.

#### E. COORDINATE TRANSFORMATION

The use of a quadratic or higher order element permits the transformation or mapping of the straight-sided parent element into an element with curved sides. Distorted or curved elements provide a better fit to curved domains than linear elements, and thus a smaller number of elements is required to represent the structure adequately.

The transformation from cartesian coordinates to curvilinear coordinates can be accomplished by employing a one-to-one correspondence defined by [11]

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = f \begin{pmatrix} \xi \\ \eta \\ \zeta \end{pmatrix} \quad \text{or} \quad f \begin{pmatrix} L_1 \\ L_2 \\ L_3 \\ \zeta \end{pmatrix} \quad (46)$$

The element shape functions are utilized to achieve this transformation via the relation

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \sum_i N_i x_i \\ \sum_i N_i y_i \\ \sum_i N_i z_i \end{pmatrix}, \quad i=1,2,\dots,n^e \quad (47)$$



where  $n^e$  is the number of element nodes, and the element shape functions  $N_i$  are in terms of local coordinates. Each set of local coordinates corresponds to only one set of cartesian coordinates.

In performing the transformation, the compatibility requirement must be met. The transformation into the new, curved elements should leave no gaps between adjacent elements. If two adjacent elements are generated from parents in which the element shape functions satisfy continuity requirements, then the curved elements will be contiguous [11]. For the isoparametric element, uniqueness of coordinates ensures compatibility. Continuity is assured when adjacent elements are given the same sets of coordinates at common nodes.

Since the element shape functions are in terms of local coordinates, an element of volume,  $dx dy dz$ , must be transformed into an element of volume expressed in local coordinates. This is achieved through the use of the Jacobian matrix defined below. Using the chain rule, the relationship between  $\xi, \eta, \zeta$  and a corresponding set of cartesian coordinates  $x, y, z$  is

$$\begin{Bmatrix} \frac{\partial N_i}{\partial \xi} \\ \frac{\partial N_i}{\partial \eta} \\ \frac{\partial N_i}{\partial \zeta} \end{Bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} & \frac{\partial z}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} & \frac{\partial z}{\partial \eta} \\ \frac{\partial x}{\partial \zeta} & \frac{\partial y}{\partial \zeta} & \frac{\partial z}{\partial \zeta} \end{bmatrix} \begin{Bmatrix} \frac{\partial N_i}{\partial x} \\ \frac{\partial N_i}{\partial y} \\ \frac{\partial N_i}{\partial z} \end{Bmatrix} \quad (48)$$



The Jacobian matrix  $[J]$  is defined as

$$[J] = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} & \frac{\partial z}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} & \frac{\partial z}{\partial \eta} \\ \frac{\partial x}{\partial \zeta} & \frac{\partial y}{\partial \zeta} & \frac{\partial z}{\partial \zeta} \end{bmatrix} \quad (49)$$

From equation (47), the Jacobian matrix becomes

$$[J] = \begin{bmatrix} \sum_i \frac{\partial N_i}{\partial \xi} x_i & \sum_i \frac{\partial N_i}{\partial \xi} y_i & \sum_i \frac{\partial N_i}{\partial \xi} z_i \\ \sum_i \frac{\partial N_i}{\partial \eta} x_i & \sum_i \frac{\partial N_i}{\partial \eta} y_i & \sum_i \frac{\partial N_i}{\partial \eta} z_i \\ \sum_i \frac{\partial N_i}{\partial \zeta} x_i & \sum_i \frac{\partial N_i}{\partial \zeta} y_i & \sum_i \frac{\partial N_i}{\partial \zeta} z_i \end{bmatrix} \quad (50)$$

The determinant of the Jacobian is used to transform the volume of element in cartesian coordinates to local coordinates. For a volume of element [11],

$$dx dy dz = \det [J] d\xi d\eta d\zeta \quad (51)$$

Note that the determinant of the Jacobian is a variable for elements of curved geometry. Only in the case of straight-sided elements is the determinant of the Jacobian a constant.

In the plane of the triangle, the area coordinates ( $L_1, L_2, L_3$ ) number one more than the cartesian coordinates ( $x, y$ ). Thus,  $L_3$  is defined as a dependent variable. This establishes the origin of the  $\xi\eta$  coordinate system at corner point 3, as





illustrated in figure 11. Recall that the  $\xi\eta$  axes need not be orthogonal. As such

$$\xi = L_1 \quad (52a)$$

$$\eta = L_2 \quad (52b)$$

Using equation (42),

$$L_3 = 1 - \xi - \eta \quad (52c)$$

Applying the chain rule yields

$$\frac{\partial N_i}{\partial \xi} = \frac{\partial N_i}{\partial L_1} \frac{\partial L_1}{\partial \xi} + \frac{\partial N_i}{\partial L_2} \frac{\partial L_2}{\partial \xi} + \frac{\partial N_i}{\partial L_3} \frac{\partial L_3}{\partial \xi} \quad (53a)$$

$$\frac{\partial N_i}{\partial \eta} = \frac{\partial N_i}{\partial L_1} \frac{\partial L_1}{\partial \eta} + \frac{\partial N_i}{\partial L_2} \frac{\partial L_2}{\partial \eta} + \frac{\partial N_i}{\partial L_3} \frac{\partial L_3}{\partial \eta} \quad (53b)$$

Using equations (52a), (52b), and (52c) in equations (53a) and (53b) gives

$$\frac{\partial N_i}{\partial \xi} = \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \quad (54a)$$

$$\frac{\partial N_i}{\partial \eta} = \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \quad (54b)$$



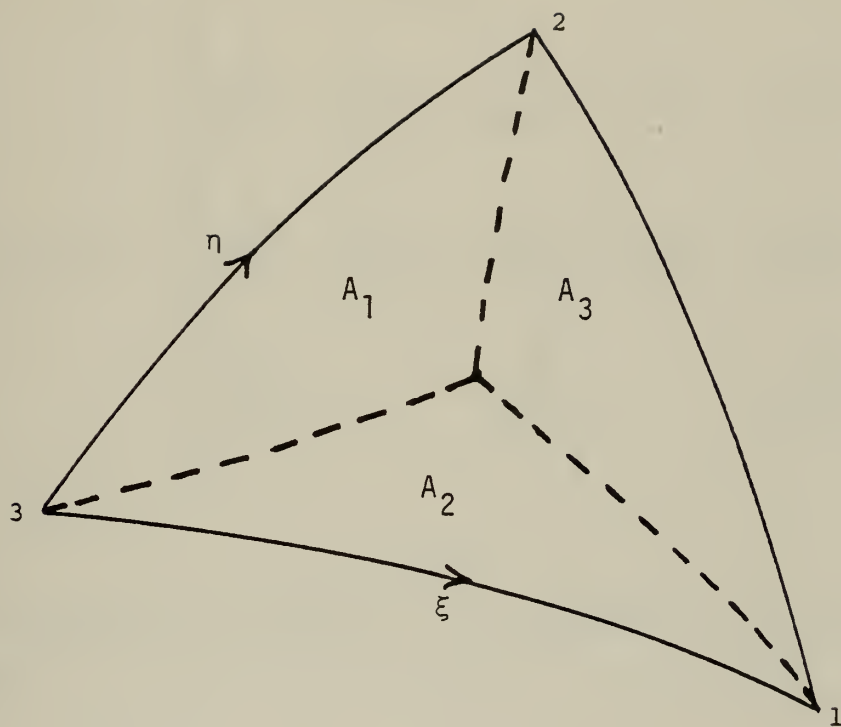


Figure 11.  $\eta\xi$  coordinates in a triangle



Now the Jacobian matrix can be evaluated using the shape functions  $N_i = N_i(L_1, L_2, L_3, \zeta)$ . Inserting equations (54a) and (54b) in equation (50) produces

$$[J(L_1, L_2, L_3, \zeta)] = \begin{bmatrix} \sum_i \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) x_i & \sum_i \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) y_i & \sum_i \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) z_i \\ \sum_i \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right) x_i & \sum_i \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right) y_i & \sum_i \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right) z_i \\ \sum_i \left( \frac{\partial N_i}{\partial \zeta} \right) x_i & \sum_i \left( \frac{\partial N_i}{\partial \zeta} \right) y_i & \sum_i \left( \frac{\partial N_i}{\partial \zeta} \right) z_i \end{bmatrix} \quad (55)$$

To summarize, suppose it is required to transform the integral

$$I = \int_{Vol} F'(L_1, L_2, L_3, \zeta) \, dx dy dz \quad (56)$$

to an integral entirely in terms of local coordinates. The determinant of equation (55) is then utilized to give

$$I = \int_{-1}^1 \int_0^1 \int_0^{1-L_1} F'(L_1, L_2, L_3, \zeta) \det[J(L_1, L_2, L_3, \zeta)] dL_1 dL_2 d\zeta \quad (57)$$

Equation (56) is now in a form suitable for numerical integration.

The development of coordinate transformation to this point has been under the general assumption that all the sides of the element are curved. In this work, the only



curved side is in the plane of the triangle. The sides of the element along the prism axis are straight. As such, a modified Jacobian matrix can be employed, thereby reducing the number of calculations to be performed. This modified Jacobian is the 2x2 matrix defined by

$$[J^*] = \begin{bmatrix} \sum_i \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) x_i & \sum_i \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) y_i \\ \sum_i \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right) x_i & \sum_i \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) y_i \end{bmatrix} \quad (58)$$

Along the prism axis or  $\zeta$  direction, it can be shown that

$$dz = \frac{h}{2} d\zeta \quad (59)$$

where  $h$  = height of the element. The volume relationship is then given by

$$dxdydz = \frac{h}{2} \det[J^*] dL_1 dL_2 d\zeta \quad (60)$$

The integral of equation (56) then assumes the form of

$$I = \frac{h}{2} \int_{-1}^1 \int_0^1 \int_0^{1-L_1} F'(L_1, L_2, L_3, \zeta) \det[J^*(L_1, L_2, L_3, \zeta)] dL_1 dL_2 d\zeta \quad (61)$$

Equation (61) is the basis of numerical integration applied in this work.

#### F. CONSTRUCTION OF ELEMENT MATRICES

The system matrix operators can be constructed through the use of the global basis functions  $G_j$  or through the





application of the element shape functions. Although the global basis functions were used to demonstrate the method of Galerkin, the construction of the system matrices in this work was achieved through element considerations. The solution of the unknown variable  $\psi$  within the element domain was approximated by

$$\tilde{\psi}^e = \sum_{i=1} N_i \psi_i^e(t) , \quad i=1,2,\dots,n^e \quad (62)$$

where  $n^e$  is the number of element nodal points,  $N_i$  are the element shape functions, and  $\psi_i^e(t)$  are the time-dependent nodal magnitudes of  $\tilde{\psi}^e$ . The element contribution to the system matrix operators is defined by Galerkin's orthogonality condition expressed by

$$\int_{Vol} N_j R^e dVol = 0 , \quad j=1,2,\dots,n^e \quad (63)$$

where  $R^e$  is defined by replacing  $\tilde{\psi}$  with  $\tilde{\psi}^e$  in equations (37) and (38), and the integration is over the element volume. Using equation (62), the element contribution is portrayed for the core by

$$\begin{aligned} & \left[ \int_{Vol} N_j N_i dVol \right] \{ \dot{\psi}_i^e(t) \} - \left[ \int_{Vol} N_j \nabla^2 N_i dVol \right] \{ (c1 \cdot \psi^e)_i \} + \left[ \int_{Vol} N_j N_i dVol \right] \{ (c2 \cdot \psi^e)_i \} \\ & + \left[ \int_{Vol} N_j N_i \ln \left( 1 + \frac{K}{\gamma T_0} \sum_k N_k \psi_k^e(t) \right) dVol \right] \{ (c4 \cdot \psi^e)_i \} \\ & + \left[ \int_{Vol} N_j N_i dVol \right] \left\{ \int_0^t e^{-\bar{\lambda}(t-t')} (c5 \cdot \psi(t'))_i dt' \right\} = 0 \end{aligned} \quad (64)$$



where  $i, j, k = 1, 2, \dots, 15$ , and  $c_1, c_2$ , etc., are constants at node  $i$ . The bracketed expressions represent square matrices of size  $15 \times 15$ , and the braced expressions represent column vectors of size  $15 \times 1$ . For the reflector, the last two terms of equation (64) are zero. Before any operation can be performed on equation (64), the nonlinear terms (last two terms) must be "linearized" and the term with the  $\nabla^2$  operator must be integrated by parts.

The nonlinear feedback term,  $\lambda_n [1 + \frac{K}{\gamma T_0} \sum_{k=1}^{15} N_k \psi_k^e(t)]$ , was linearized by using predicted values of the unknown function at time  $t$ . These predicted values come from the time integration scheme by Franke [7]. Basically, the integration scheme utilizes a predictor-corrector method which predicts values of the unknown function at the next time by using the derivatives of the function. Adopting the predicted values  $\psi_k^P$  enabled integration over space. The term in equation (64) involving the nonlinear feedback can be written as

$$\left[ \int_{Vol} N_j N_i \lambda_n \left( 1 + \frac{K}{\gamma T_0} (N_1 \psi_1^P + N_2 \psi_2^P + \dots + N_{15} \psi_{15}^P) \right) dVol \right] \{ (c_4 \cdot \psi^e)_i \}$$

The last term of equation (64) describes the delayed neutron contribution. In general,

$$\begin{aligned} e^{-\bar{\lambda} t_n} \int_0^t e^{\bar{\lambda} t'} \psi_i^e(t') dt' &= e^{-\bar{\lambda} (t_n - t_{n-1})} \left[ e^{-\bar{\lambda} t_{n-1}} \int_0^{t_{n-1}} e^{\bar{\lambda} t'} \psi_i^e(t') dt' \right] \\ &+ e^{-\bar{\lambda} t_n} \int_{t_{n-1}}^t e^{\bar{\lambda} t'} \psi_i^e(t') dt' \end{aligned} \quad (65)$$

where  $n$  is number of time steps,  $t_n$  is the current time,



and  $t_{n-1}$  is the previous time. To approximate the integrals in equation (65), the predicted values  $\psi_i^P$  were employed in a simple trapezoidal rule. The trapezoidal rule was believed to be sufficient since very small time steps were utilized in this work. For example, at time  $t_2$ ,

$$(SUM_i) = (\text{previous } SUM_i) e^{-\bar{\lambda}H} + \frac{H}{2}(e^{\bar{\lambda}H} \psi_i^e(t_1) + \psi_i^P)$$

where

$$H = \text{current time step} = t_2 - t_1$$

$$(SUM_i) = e^{-\bar{\lambda}t_2} \int_0^{t_2} e^{\bar{\lambda}t'} \psi_i^e(t') dt'$$

$$(\text{previous } SUM_i) = e^{-\bar{\lambda}t_1} \int_0^{t_1} e^{\bar{\lambda}t'} \psi_i^e(t') dt'$$

The previous sum is formed by accumulating the trapezoidal integration of each time step. In general then, for time  $t_n$

$$(SUM_i) = (\text{previous } SUM_i) e^{-\bar{\lambda}H} + \frac{H}{2}(e^{\bar{\lambda}H} \psi_i^e(t_{n-1}) + \psi_i^P) \quad (66)$$

The delayed neutron term in equation (64) can be expressed as

$$\left[ \int_{Vol} N_j N_i dVol \right] \{ (c5 \cdot SUM)_i \}$$



In order to bring equation (64) to final form, the  $\nabla^2$  operator was integrated by parts. Since the flux at the surface of the reactor is zero, integration by parts yields

$$\int_{Vol} N_j \nabla^2 N_i dx dy dz = - \int_{Vol} \left( \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} + \frac{\partial N_i}{\partial y} \frac{\partial N_j}{\partial y} + \frac{\partial N_i}{\partial z} \frac{\partial N_j}{\partial z} \right) dx dy dz \quad (67)$$

In vector notation, equation (67) is expressed as

$$\int_{Vol} N_j \nabla^2 N_i dx dy dz = - \int_{Vol} \left\langle \frac{\partial N_i}{\partial x}, \frac{\partial N_i}{\partial y}, \frac{\partial N_i}{\partial z} \right\rangle \begin{Bmatrix} \frac{\partial N_j}{\partial x} \\ \frac{\partial N_j}{\partial y} \\ \frac{\partial N_j}{\partial z} \end{Bmatrix} dx dy dz \quad (68)$$

Using the chain rule produces

$$\begin{Bmatrix} \frac{\partial N_i}{\partial x} \\ \frac{\partial N_i}{\partial y} \\ \frac{\partial N_i}{\partial z} \end{Bmatrix} = \begin{bmatrix} \frac{\partial L_1}{\partial x}, \frac{\partial L_2}{\partial x}, 0 \\ \frac{\partial L_1}{\partial y}, \frac{\partial L_2}{\partial y}, 0 \\ 0, 0, \frac{\partial \zeta}{\partial z} \end{bmatrix} \begin{Bmatrix} \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right) \\ \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right) \\ \left( \frac{\partial N_i}{\partial \zeta} \right) \end{Bmatrix} \quad (69)$$

Letting the 3x3 matrix above be  $[B']$ , the  $\nabla^2$  term becomes

$$\begin{aligned} \int_{Vol} N_j \nabla^2 N_i dx dy dz = & - \int_{Vol} \left\langle \left( \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3} \right), \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right), \frac{\partial N_i}{\partial \zeta} \right\rangle [B']^T [B'] \begin{Bmatrix} \left( \frac{\partial N_j}{\partial L_1} - \frac{\partial N_j}{\partial L_3} \right) \\ \left( \frac{\partial N_j}{\partial L_2} - \frac{\partial N_j}{\partial L_3} \right) \\ \frac{\partial N_j}{\partial \zeta} \end{Bmatrix} dx dy dz \end{aligned} \quad (70)$$





where  $[B']^T$  is the transpose of  $[B']$ . By applying the chain rule in equation (47) and using equation (59), it can be shown that for this work

$$[B'] = \begin{bmatrix} \frac{1}{(\frac{\partial N_1}{\partial L_1} x_1 + \dots + \frac{\partial N_{15}}{\partial L_1} x_{15})}, \frac{1}{(\frac{\partial N_1}{\partial L_2} x_1 + \dots + \frac{\partial N_{15}}{\partial L_2} x_{15})}, 0 \\ \frac{1}{(\frac{\partial N_1}{\partial L_1} y_1 + \dots + \frac{\partial N_{15}}{\partial L_1} y_{15})}, \frac{1}{(\frac{\partial N_1}{\partial L_2} y_1 + \dots + \frac{\partial N_{15}}{\partial L_2} y_{15})}, 0 \\ 0, 0, \frac{2}{h} \end{bmatrix} \quad (71)$$

$[B']^T$  can be derived from equation (71).

Note from equation (64) that there are three basic element matrices which are defined as follows after applying equation (60):

$$[G_{ji}] = \frac{h}{2} \int_{-1}^1 \int_0^1 \int_0^{1-L_1} N_j N_i \det[J^*] dL_1 dL_2 d\zeta \quad (72)$$

$$[GG_{ji}] = \frac{h}{2} \int_{-1}^1 \int_0^1 \int_0^{1-L_1} \left\langle \frac{\partial N_i}{\partial L_1} - \frac{\partial N_i}{\partial L_3}, \left( \frac{\partial N_i}{\partial L_2} - \frac{\partial N_i}{\partial L_3} \right), \left( \frac{\partial N_i}{\partial \zeta} \right) \right\rangle [B']^T$$

$$[B'] \left\{ \begin{array}{c} \left( \frac{\partial N_j}{\partial L_1} - \frac{\partial N_j}{\partial L_3} \right) \\ \left( \frac{\partial N_j}{\partial L_2} - \frac{\partial N_j}{\partial L_3} \right) \\ \frac{\partial N_j}{\partial \zeta} \end{array} \right\} \det[J^*] dL_1 dL_2 d\zeta \quad (73)$$

$$[GGG_{ji}] = \frac{h}{2} \int_{-1}^1 \int_0^1 \int_0^{1-L_1} N_j N_i \ln \left( 1 + \frac{K}{\gamma T_0} (N_1 \psi_1^P + \dots + N_{15} \psi_{15}^P) \right) \det[J^*] dL_1 dL_2 d\zeta \quad (74)$$



Element matrices  $[G_{ji}]$  and  $[GG_{ji}]$  are independent of time. However,  $[GGG_{ji}]$  is time dependent due to the utilization of the predicted values  $\psi_i^p$  which changes with time. In terms of these three basic element matrices, equation (64) is

$$[G_{ji}]\{\dot{\psi}_i^e\} + [GG_{ji}]\{c1 \cdot \psi_i^e\} + [G_{ji}]\{(c2 \cdot \psi_i^e)\} \\ + [GGG_{ji}]\{(c4 \cdot \psi_i^e)\} + [G_{ji}]\{(c5 \cdot \text{SUM})_i\} = 0 \quad (75)$$

The last two terms of equation (75) are zero for the reflector.

#### G. CONSTRUCTION OF THE SYSTEM MATRICES

The 15x15 coefficient element matrices were calculated according to equations (72), (73), and (74); and the results were collected element by element into the corresponding system coefficient matrices. The system coefficient matrix  $[BIGG]$  is developed from  $[G_{ji}]$ ,  $[BIGGG]$  from  $[GG_{ji}]$ , and  $[BIGH]$  from  $[GGG_{ji}]$ .  $[BIGG]$  and  $[BIGGG]$  are independent of time and can be constructed once and for all from geometry considerations.  $[BIGH]$  is dependent on both geometry and time due to the time dependence of the predicted flux utilized in the feedback term. Thus,  $[BIGH]$  is recalculated at each time increment.

Non-zero contributions to a global nodal point I come only from adjacent elements sharing that same nodal point I. Thus, the system matrices are sparse and banded. The process of assembling contributions from element matrices



requires the identification of a local nodal point ( $i=1,2,\dots,15$ ) with a global nodal point ( $I=1,2,\dots,\text{NUMNP}$ , where NUMNP is the total number of system nodes). This correspondence between element and global nodes is accomplished via the connectivity matrix.

The formal treatment of the field equations in terms of the system coefficient matrices is described by the equation

$$[\text{BIGG}] \{\dot{\psi}_I\} + [\text{BIGGG}] \{(c1 \cdot \psi)_I\} + [\text{BIGG}] \{(c2 \cdot \psi)_I\} + [\text{BIGH}] \{(c4 \cdot \psi)_I\} + [\text{BIGG}] \{(c5 \cdot \text{SUM})_I\} = 0 \quad (76)$$

where the system matrices are NUMNP x NUMNP and the column vectors are of length NUMNP x 1. However, the direct application of equation (76) requires a large amount of computer storage. To take advantage of the sparsity of the system matrices, an optimum compacting scheme described by Ref. 8 was employed.

The concept behind OCS is simply to store only the non-zero terms of a coefficient matrix. OCS requires two integer arrays, say JB and NAME, and a vector of non-zero coefficients of the square system matrix. For purposes of illustration, the square system matrix is called B, and the vector of non-zero coefficients of B is called BB. The  $i^{\text{th}}$  integer entry in the NUMNP x 1 JB vector is the number  $q_i$ . This number is defined by

$$q_i = 1 + \sum_{j=1}^{i=1} p_j, \quad i=1,2,\dots,\text{NUMNP} \quad (77)$$

where  $p_j$  is the number of terms in the  $i^{\text{th}}$  equation. In



other words,  $p_j$  is the number of nodes that the  $i^{\text{th}}$  node "sees". JB is therefore a pointer vector of length NUMNP+1 whose  $i^{\text{th}}$  term locates the initial position in the BB vector of the contributing coefficients to the  $i^{\text{th}}$  equation. The NAME vector of length Mx1, where  $M = \sum_{i=1}^{\text{NUMNP}} p_i$ , consists of NUMNP successive vector blocks of variable length  $p_i$ ,  $i=1,2,\dots,\text{NUMNP}$ . The  $p_i$  integer numbers in the  $i^{\text{th}}$  block of NAME list the  $p_i$  contributors to the  $i^{\text{th}}$  equation. The Mx1 BB vector contains the real non-zero coefficients of the NUMNPxNUMNP B matrix, arranged in the same contiguous block arrangement as the NAME vector. The  $j^{\text{th}}$  term in the  $i^{\text{th}}$  block or  $\text{BB}(\text{JB}(I) + J - 1)$  is  $B(I,K)$ , where  $K = \text{NAME}(\text{JB}(I) + J - 1)$ . To illustrate, consider the grid shown in figure 12. The two array vectors and the coefficient vector matrix of non-zero terms are:

$$\text{JB} = \langle 1, 5, 10, 13, 18, 25, 30, 33, 38, 42 \rangle$$

$$\text{NAME} = \langle 1, 2, 4, 5 | 2, 3, 1, 6, 5 | \text{---} | 9, 8, 6, 5 \rangle$$

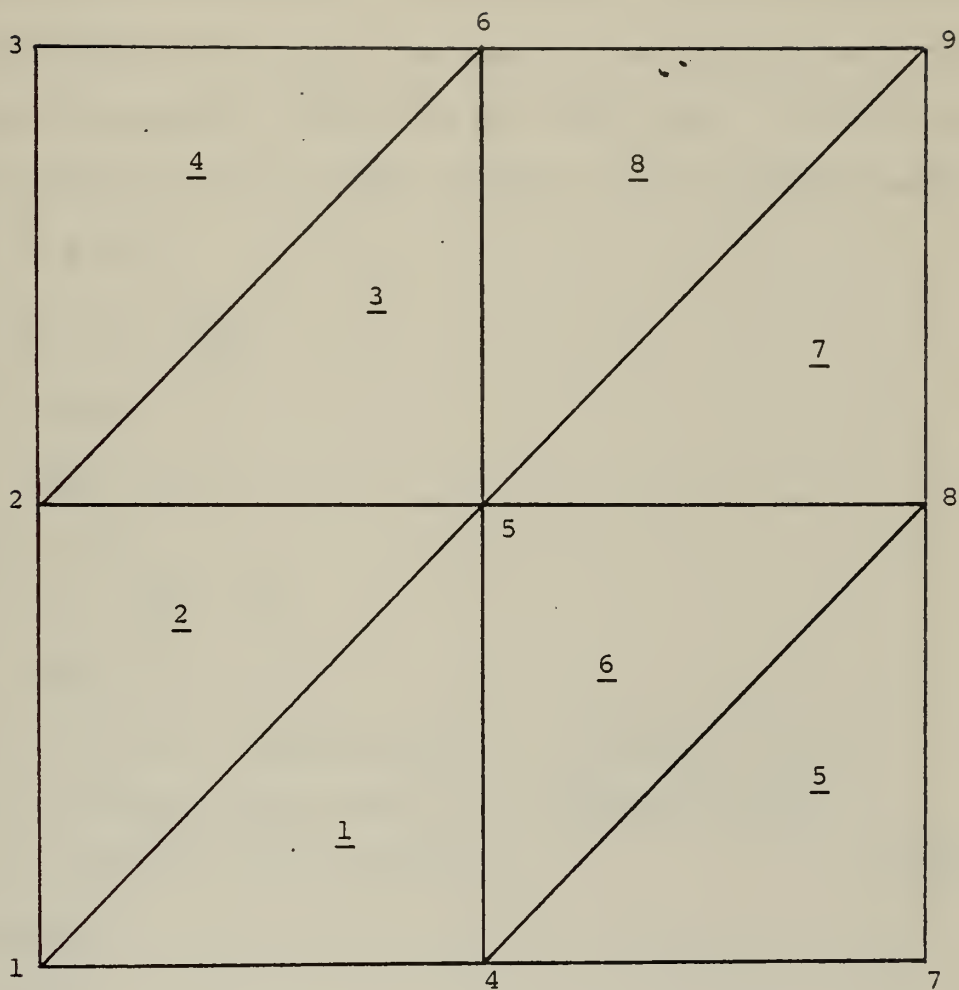
$$\text{BB} = \langle B_{11}, B_{12}, B_{14}, B_{15} | B_{22}, B_{23}, B_{21}, B_{26}, B_{25} | \text{---} | B_{99}, B_{98}, B_{96}, B_{95} \rangle$$

In this illustration, NUMNP = 9 and M = 41 [8].

In this work, a judicious method of numbering the system nodal points was adopted to further reduce computer storage requirements. Since the surface boundary nodes of the reactor represent zero neutron fluxes, the contributions of these nodes to interior or non-zero nodes can be discarded. Thus, only the interior nodal points need to be considered. These non-zero nodes were numbered first in the finite element mesh







Number - element number

Figure 12. Sample grid used for illustrating OCS



used so that in the OCS, the number of non-zero nodes (NNZ) replaces NUMNP.

The vectors of non-zero coefficients will be designated BIGG, BIGGG and BIGH since the square coefficient matrices described in equation (76) were not utilized. To illustrate the application of OCS in the system, the following sample program is given:

```
DO 450 I=1, NNZ
JBB = JB(I)
JE = JB(I+1)-1
DY(I) = 0.0
DO 500 J=JBB, JE
LL = NAME(J)
DY(I) = DY(I) + BIGG(J)* $\dot{\psi}$ (LL) + BIGGG(J)*c1(LL)*
 $\psi$ (LL) + BIGG(J)*c2(LL)* $\psi$ (LL) + BIGH(J)*c4(LL)*
 $\psi$ (LL) + BIGG(J)*c5(LL)*SUM(LL) (78)
500 CONTINUE
450 CONTINUE
```

where

LL = nodal point "seen" by node I

JE-JBB = total number of nodal points that node I "sees"

DY(I) = summation of all contributions to node I and  
should sum to zero as stated by the field equation

The assumption of a homogeneous reflector was relaxed in the application of equation (78). Interface nodes were assigned properties of the core. Therefore, if LL is a reflector



node not on the core reflector interface, the last two terms of equation (78) are nonexistent.



#### IV. NUMERICAL INTEGRATION

##### A. LINE AND AREA INTEGRATION

The solution of the element matrix equations was achieved through numerical integration since an exact closed form solution cannot be established. The volume integration was accomplished by using a line integration in the  $\zeta$ -direction and an area integration in the plane of the triangle. The line integral is described by the Gaussian quadrature formula [12]

$$\int_{-1}^1 f(\zeta) d\zeta \approx \sum_{k=1}^n H_k f(a_k) \quad (79)$$

where  $n$  is the number of Gauss integration points

$H_k$  = weighting coefficients

$f(a_k)$  = the function  $f(\zeta)$  evaluated at Gauss point  $a_k$

Table III lists  $a_k$ ,  $H_k$  and  $n$  [11].

The area integration was achieved by the equation

$$\int_0^1 \int_0^{1-L_1} f(L_1, L_2, L_3) dL_1 dL_2 \approx \sum_{m=1}^{\bar{m}} w_m f(L_1^m, L_2^m, L_3^m) \quad (80)$$

where  $\bar{m}$  is the number of area integration points and  $w_m$  are the weights. The numerical integration points for the area integration are given in Table IV which was extracted





TABLE III

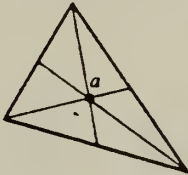
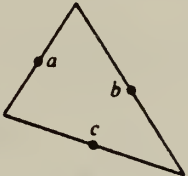
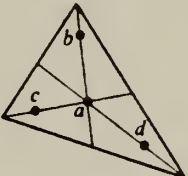
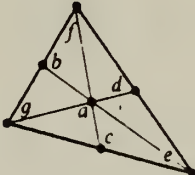
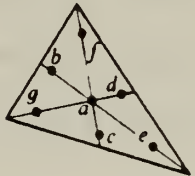
Abcissae and Weight Coefficients  
of the Gaussian Quadrature Formula

$$\int_{-1}^1 f(\zeta) d\zeta = \sum_{k=1}^n H_k f(a_k)$$

$\pm a$			$H$		
$n = 2$					
0.57735	02691	89626	1.00000	00000	00000
$n = 3$					
0.77459	66692	41483	0.55555	55555	55556
0.00000	00000	00000	0.88888	88888	88889
$n = 4$					
0.86113	63115	94053	0.34785	48451	37454
0.33998	10435	84856	0.65214	51548	62546
$n = 5$					
0.90617	98459	38664	0.23692	68850	56189
0.53846	93101	05683	0.47862	86704	99366
0.00000	00000	00000	0.56888	88888	88889
$n = 6$					
0.93246	95142	03152	0.17132	44923	79170
0.66120	93864	66265	0.36076	15730	48139
0.23861	91860	83197	0.46791	39345	72691
$n = 7$					
0.94910	79123	42759	0.12948	49661	68870
0.74153	11855	99394	0.27970	53914	89277
0.40584	51513	77397	0.38183	00505	05119
0.00000	00000	00000	0.41795	91836	73469
$n = 8$					
0.96028	98564	97536	0.10122	85362	90376
0.79666	64774	13627	0.22238	10344	53374
0.52553	24099	16329	0.31370	66458	77887
0.18343	46424	95650	0.36268	37833	78362
$n = 9$					
0.96816	02395	07626	0.08127	43883	61574
0.83603	11073	26636	0.18064	81606	54857
0.61337	14327	00590	0.26061	06964	02935
0.32425	34234	03809	0.31234	70770	40003
0.00000	00000	00000	0.33023	93550	01260
$n = 10$					
0.97390	65285	17172	0.06667	13443	08688
0.86506	33666	88985	0.14945	13491	50581
0.67940	95682	99024	0.21908	63625	15982
0.43339	53941	29247	0.26926	67193	09996
0.14887	43389	81631	0.29552	42247	14753



TABLE IV  
Numerical Formulas for Triangles

Order	Fig.	Error	Points	Triangular Co-ordinates	Weights $2W_i$
Linear		$R = O(h^2)$	$a$	$\frac{1}{3}, \frac{1}{3}, \frac{1}{3}$	1
Quadratic		$R = O(h^3)$	$a$ $b$ $c$	$\frac{1}{2}, \frac{1}{2}, 0$ $0, \frac{1}{2}, \frac{1}{2}$ $\frac{1}{2}, 0, \frac{1}{2}$	$\frac{1}{3}$ $\frac{1}{3}$ $\frac{1}{3}$
Cubic		$R = O(h^4)$	$a$ $b$ $c$ $d$	$\frac{1}{3}, \frac{1}{3}, \frac{1}{3}$ $\frac{1}{3}, \frac{2}{3}, \frac{1}{3}$ $\frac{2}{3}, \frac{1}{3}, \frac{1}{3}$ $\frac{1}{3}, \frac{1}{3}, \frac{1}{3}$	$-\frac{27}{248}$ $\frac{27}{248}$
Cubic		$R = O(h^4)$	$a$ $b$ $c$ $d$ $e$ $f$ $g$	$\frac{1}{3}, \frac{1}{3}, \frac{1}{3}$ $\frac{1}{2}, \frac{1}{2}, 0$ $0, \frac{1}{2}, \frac{1}{2}$ $\frac{1}{2}, 0, \frac{1}{2}$ $1, 0, 0$ $0, 1, 0$ $0, 0, 1$	$\frac{27}{80}$ $\frac{8}{80}$ $\frac{3}{80}$
Quintic		$R = O(h^6)$	$a$ $b$ $c$ $d$ $e$ $f$ $g$	$\frac{1}{3}, \frac{1}{3}, \frac{1}{3}$ $\alpha_1, \beta_1, \beta_1$ $\beta_1, \alpha_1, \beta_1$ $\beta_1, \beta_1, \alpha_1$ $\alpha_2, \beta_2, \beta_2$ $\beta_2, \alpha_2, \beta_2$ $\beta_2, \beta_2, \alpha_2$	0.225 0.13239415 0.12593918
with $\alpha_1 = 0.05971587$ $\beta_1 = 0.47014206$ $\alpha_2 = 0.79742699$ $\beta_2 = 0.10128651$					



from Ref. 11. The volume integration of the function  $f(L_1, L_2, L_3, \zeta)$  using numerical integration is therefore

$$\int_{-1}^1 \int_0^1 \int_0^{1-L_1} f(L_1, L_2, L_3, \zeta) dL_1 dL_2 d\zeta \approx \sum_{k=1}^{\bar{n}} w_k \sum_{m=1}^{\bar{m}} w_m f(L_1^m, L_2^m, L_3^m, \zeta^k) \quad (81)$$

where  $w_k = H_k$ . In the manner of equation (81), the element matrices become

$$[G_{ji}] = \frac{h}{2} \sum_{k=1}^{\bar{n}} w_k \sum_{m=1}^{\bar{m}} w_m N_j(L_1^m, L_2^m, L_3^m, \zeta^k) N_i \det[J^*] \quad (82)$$

$$[GG_{ji}] = \frac{h}{2} \sum_{k=1}^{\bar{n}} w_k \sum_{m=1}^{\bar{m}} w_m F(L_1^m, L_2^m, L_3^m, \zeta^k) \det[J^*] \quad (83)$$

$$[GGG_{ji}] = \frac{h}{2} \sum_{k=1}^{\bar{n}} w_k \sum_{m=1}^{\bar{m}} w_m T(L_1^m, L_2^m, L_3^m, \zeta^k) \det[J^*] \quad (84)$$

where  $F$  and  $T$  are easily derived from equations (73) and (74), respectively. Note that  $N_i$  and  $\det[J^*]$  are also evaluated at each integration point. They are not shown as such merely for the sake of convenience in writing the equations.

## B. NUMBER OF INTEGRATION POINTS

It is difficult to estimate the number of integration points required for good accuracy due to the complexity of the functions involved. The basic rule that the best number of integration points is found by trial and experience was adopted. For the  $[G_{ji}]$  element matrix, the numbers involved



can be approximated. This was done by letting the determinant of the Jacobian equal to twice the area of the curved triangle (checking  $\det[J^*]$  at each integration point showed that this assumption was not too unreasonable). With  $\det[J^*]$  outside the integration process,  $[G_{ji}]$  can be solved in closed form by integrating out  $\zeta$  from -1 to +1 and then applying the closed form equation [12]

$$\int_{\text{Area}} L_1^{m_1} L_2^{m_2} L_3^{m_3} d(\text{Area}) = 2A \frac{m_1! m_2! m_3!}{(m_1 + m_2 + m_3 + 2)!} \quad (85)$$

where  $m_1, m_2, m_3$  are positive integer exponents and  $A$  is the area of the triangle.

Five test points within the 15x15 element matrix were selected, as listed in Table V. The values obtained from the application of equation (85) are also given in Table V. Three sets of area integration points, each with a different number of  $\zeta$  Gauss points, were used. These three sets of area integration points, given in Table IV, are:

1. cubic order (4 points)
2. cubic order (7 points)
3. quintic order (7 points)

Using each of the three integration points above with different  $\zeta$  Gauss points in equation (82) produced the results obtained in Table V. From these results, the quintic order area integration points were selected for the element matrix  $[G_{ji}]$  with three  $\zeta$  Gauss points in the prism axes. This set of integration points was also used for  $[GGG_{ji}]$ .





TABLE V.

Selection of Integration Points for  $[G_{ji}]$ 

$\zeta$ points	Area points	$G(1,1)$	$G(2,2)$	$G(1,9)$	$G(6,9)$	$G(9,9)$
13	cubic (4 pts)	1415	5278	-2756	5072	8536
5	"	1817	5278	-3170	5072	10240
7	"	1817	5278	-3170	5072	10240
3	cubic (7 pts)	3248	8247	-3114	4960	10243
5	"	3249	8247	-3114	4960	10240
7	"	3249	8247	-3114	4960	10240
3	quintic (7 pts)	2464	6600	-3144	5020	10240
5	"	2464	6600	-3144	5020	10240
7	"	2464	6600	-3144	5020	10240
Approximated values		2500	6700	-3142	5026	10053



The numbers for the  $[GG_{ji}]$  element matrix could not be approximated. The best that could be done was to obtain an idea of the order of magnitude of this element matrix. Towards this end, a linear triangular element was assumed. Using linear approximation, the order of magnitude was found to be about  $10^3$ . Using the three different area integration points mentioned above with varying  $\zeta$  Gauss points in equation (83) yielded the results given in Table VI. Further checks on the  $[GG_{ji}]$  element matrix showed that the cubic order with four area integration points yielded the desired order of magnitude. Thus, the fourth order cubic with five  $\zeta$  Gauss points was employed for the  $[GG_{ji}]$  element matrix. A note should be mentioned here in regards to the vast difference in results obtained for  $[GG_{ji}]$  using different area integration points. Most likely, it was due to the  $[B']$  and  $[B']^T$  matrices which required the inversion of  $\frac{\partial x}{\partial L_1}$ ,  $\frac{\partial x}{\partial L_2}$ , etc. Or perhaps it was caused by the nature of the hybrid element used in this work. In any case, further investigation is warranted in this area.



TABLE VI

Selection of Integration Points for  $[GG_{ji}]$ 

$\zeta$ points	Area points	GG(1,9)	GG(3,3)	GG(14,9)	GG(8,8)	GG(15,15)
3	cubic (4 pts)	18.4	170.1	-165	165.9	106.1
5	"	22.0	177.6	-186	195.9	106.1
7	"	22.0	177.6	-186	195.9	106.1
3	cubic (7 pts)	$-1.03 \times 10^4$	$2.11 \times 10^7$	$1.05 \times 10^7$	$8.42 \times 10^7$	2056
5	"	$-1.03 \times 10^4$	$2.11 \times 10^7$	$1.05 \times 10^7$	$8.42 \times 10^7$	2056
7	"	$-1.03 \times 10^4$	$2.11 \times 10^7$	$1.05 \times 10^7$	$8.41 \times 10^7$	2056
3	quintic (7 pts)	-52.6	$1.27 \times 10^4$	$-6.64 \times 10^4$	$6.84 \times 10^4$	$1.51 \times 10^5$
5	"	-52.6	$1.27 \times 10^4$	$-6.64 \times 10^4$	$6.84 \times 10^4$	$1.51 \times 10^5$
7	"	-52.6	$1.27 \times 10^4$	$-6.64 \times 10^4$	$6.84 \times 10^4$	$1.51 \times 10^5$



## V. TEST PROBLEMS AND RESULTS

The reactor was subjected to uniform and local perturbations in the form of a ramp input described by

$$\Sigma_f = \Sigma_f^* + \alpha t \quad (86)$$

where  $\alpha$  is the change in  $\Sigma_f$  per unit time and  $\Sigma_f^*$  is the critical fission cross-section.  $\Sigma_f^*$  must first be obtained before the perturbations can be applied. This was accomplished by trial and error until a stationary solution was reached. For mesh I,  $\Sigma_f^*$  was found to be 0.0057360 per cm. No attempt was made to find  $\Sigma_f^*$  for mesh II due to time limitations. As such, the test problems outlined below were applied to mesh I. Future work is planned to apply the test problems on mesh II.

The following perturbations were applied:

a) Uniform perturbation of 10 dollar of reactivity per second:

$$\Sigma_f(\bar{r}, t) = \Sigma_f^* + \alpha t, \quad \text{in the core}$$

where  $\alpha = 0.005893/\text{cm-sec}$

b) Local perturbation at the core center of 100 dollar of reactivity per second:

$$\Sigma_f(\bar{r}, t) = \Sigma_f^* + \alpha t \delta(\bar{r}_0), \quad \text{in the core}$$





where  $\bar{r}_0$  is (0,0,0) and  $\alpha = 0.015407/\text{cm-sec}$

c) Local, off-center perturbation:

$$\Sigma_f(\bar{r},t) = \Sigma_f^* + \alpha_i t \delta(\bar{r}_1) , \quad i = 1,2,3, \text{ in the core}$$

where

$$\bar{r}_1 = (0,60,40)$$

$$\alpha_1 = 0.015407/\text{cm-sec} = 100 \text{ dollar per second}$$

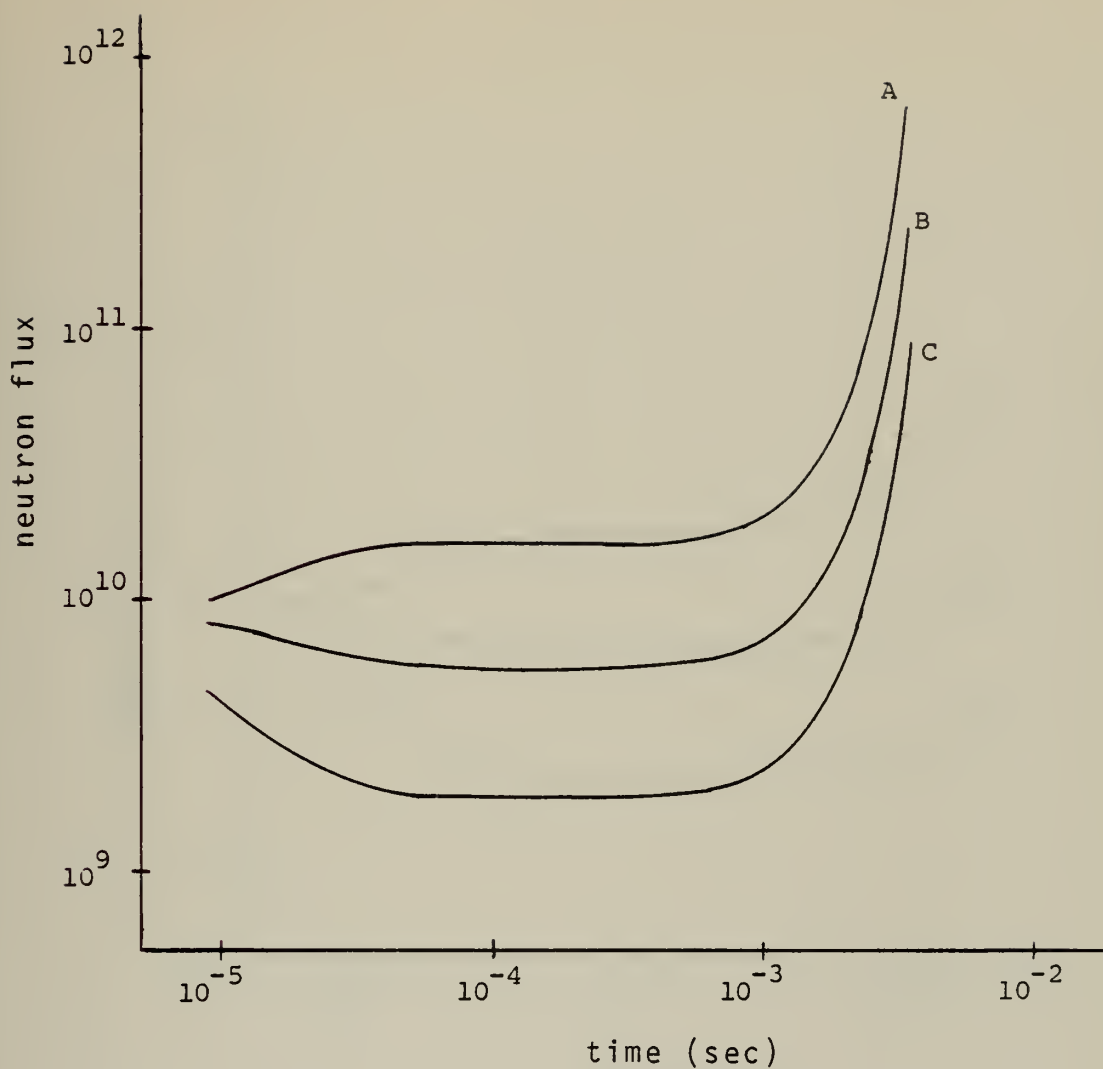
$$\alpha_2 = 0.008123/\text{cm-sec} = 50 \text{ dollar per second}$$

$$\alpha_3 = 0.005894/\text{cm-sec} = 10 \text{ dollar per second}$$

Three test points, (0,0,0), (60,0,0), and (-60,0,80), were selected to trace the neutron time history. For cases a) and b), the neutron flux was plotted at each test point during transience. This is shown in figures 13 and 14. Case c) involved three ramp inputs and were conducted for both the linear and nonlinear reactor equations. The linear and nonlinear responses were compared at each test point for each ramp input and are illustrated in figures 15 through 23. The radial and axial flux distributions at time  $t = 0.0123$  second were plotted for the steady state and the 100 dollar perturbation case. These are embodied in figures 24 and 25. Finally, a neutron flux early time history between mesh I and mesh II was plotted to check the effect of using a finer element mesh. The result is portrayed in figure 26.

Figures 13 and 14 revealed a clear space dependence of the neutron flux during transience as expected. Figures 15 through 23 demonstrated the effects of temperature feedback

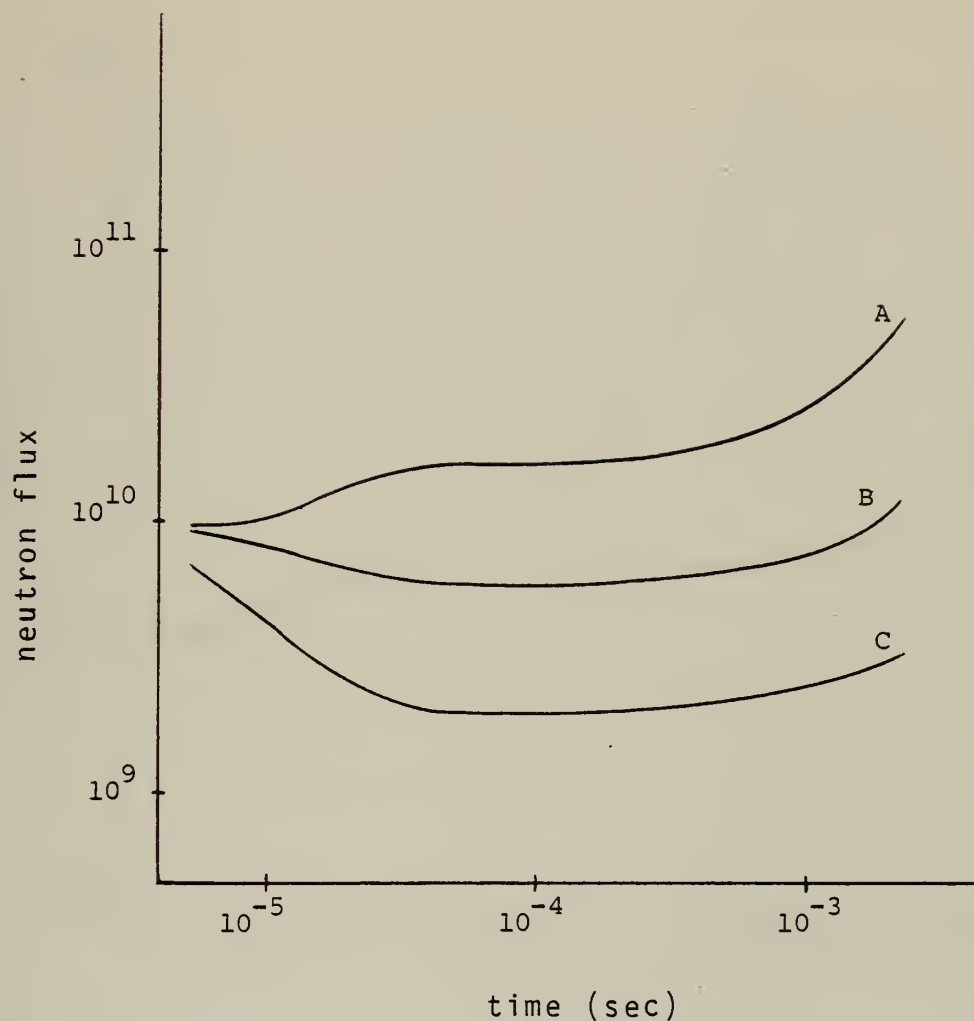




- A - neutron flux at test point (0,0,0)
- B - neutron flux at test point (60,0,0)
- C - neutron flux at test point (-60,0,80)

Figure 13. Neutron flux transient history at three test points with a uniform perturbation of 10 dollar of reactivity per second.

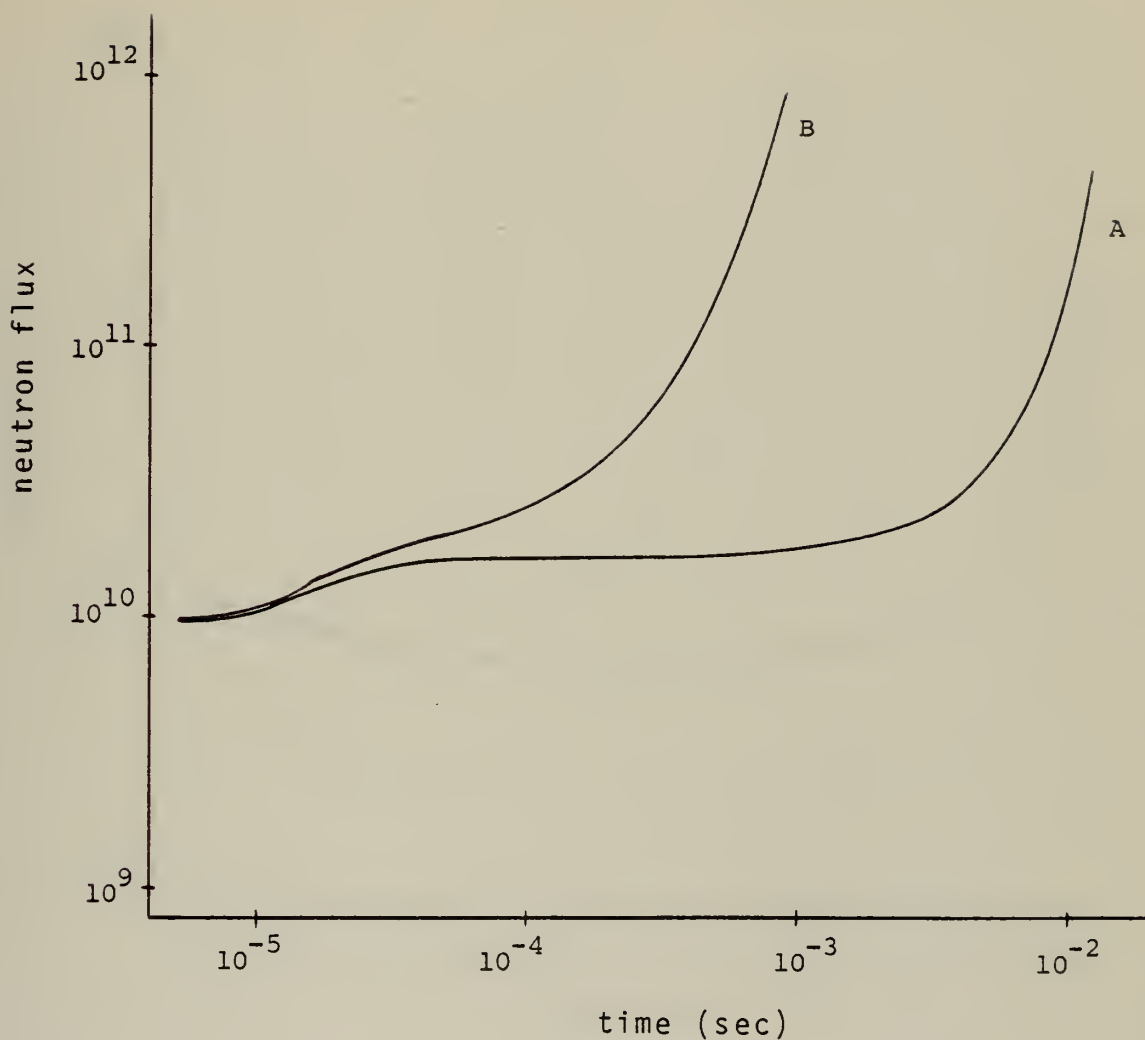




A - neutron flux at test point (0,0,0)  
 B - neutron flux at test point (60,0,0)  
 C - neutron flux at test point (-60,0,80)

Figure 14. Neutron flux transient history at various test points for a local central perturbation of 100 dollar/sec of reactivity.



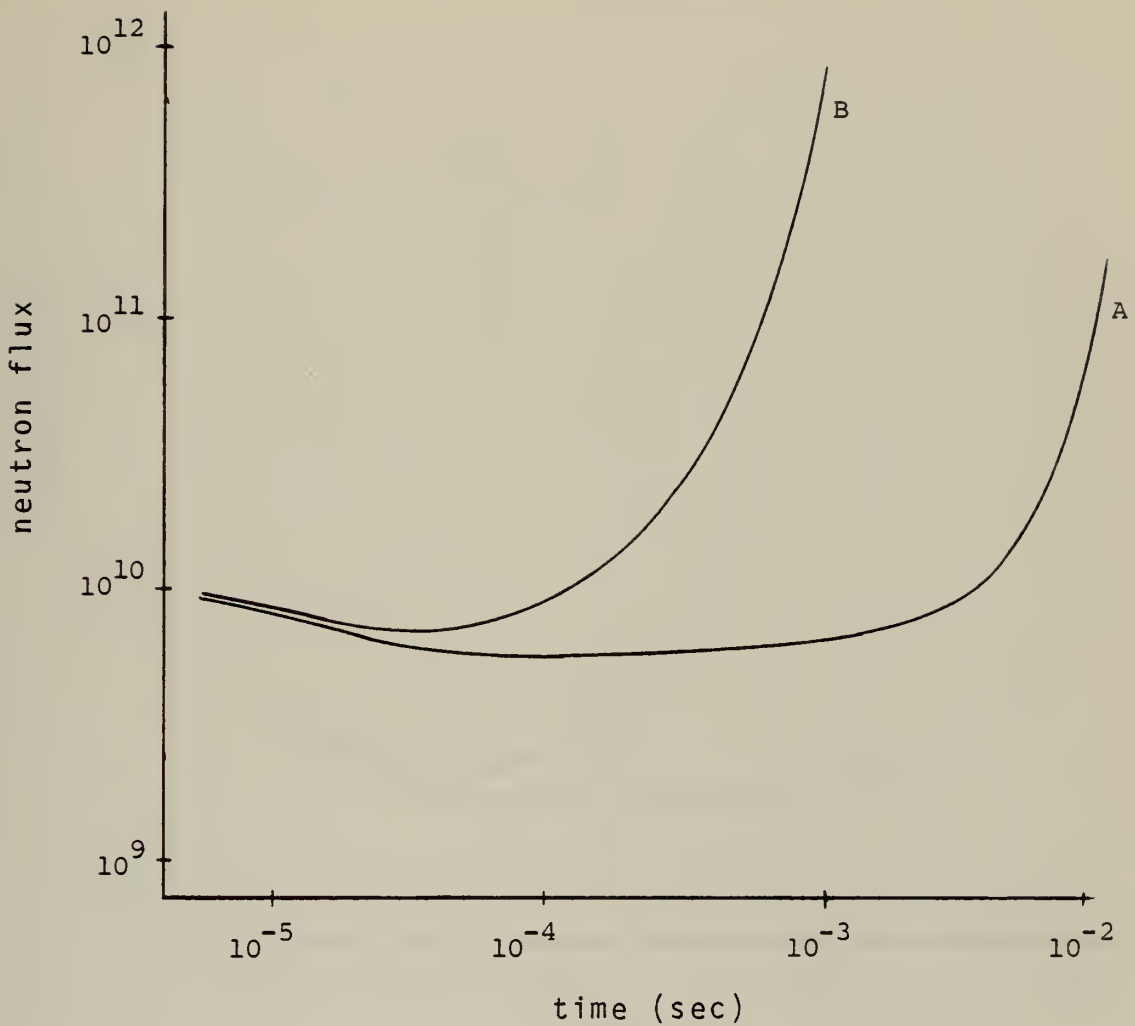


A - nonlinear neutron response  
B - linear neutron response

Figure 15. Linear and nonlinear fluxes at (0,0,0) due to a local perturbation of 100 dollar/sec of reactivity at (0,60,40).







A - nonlinear neutron response  
B - linear neutron response

Figure 16. Linear and nonlinear fluxes at (60,0,0) due to a local perturbation of 100 dollar/sec at (0,60,40).



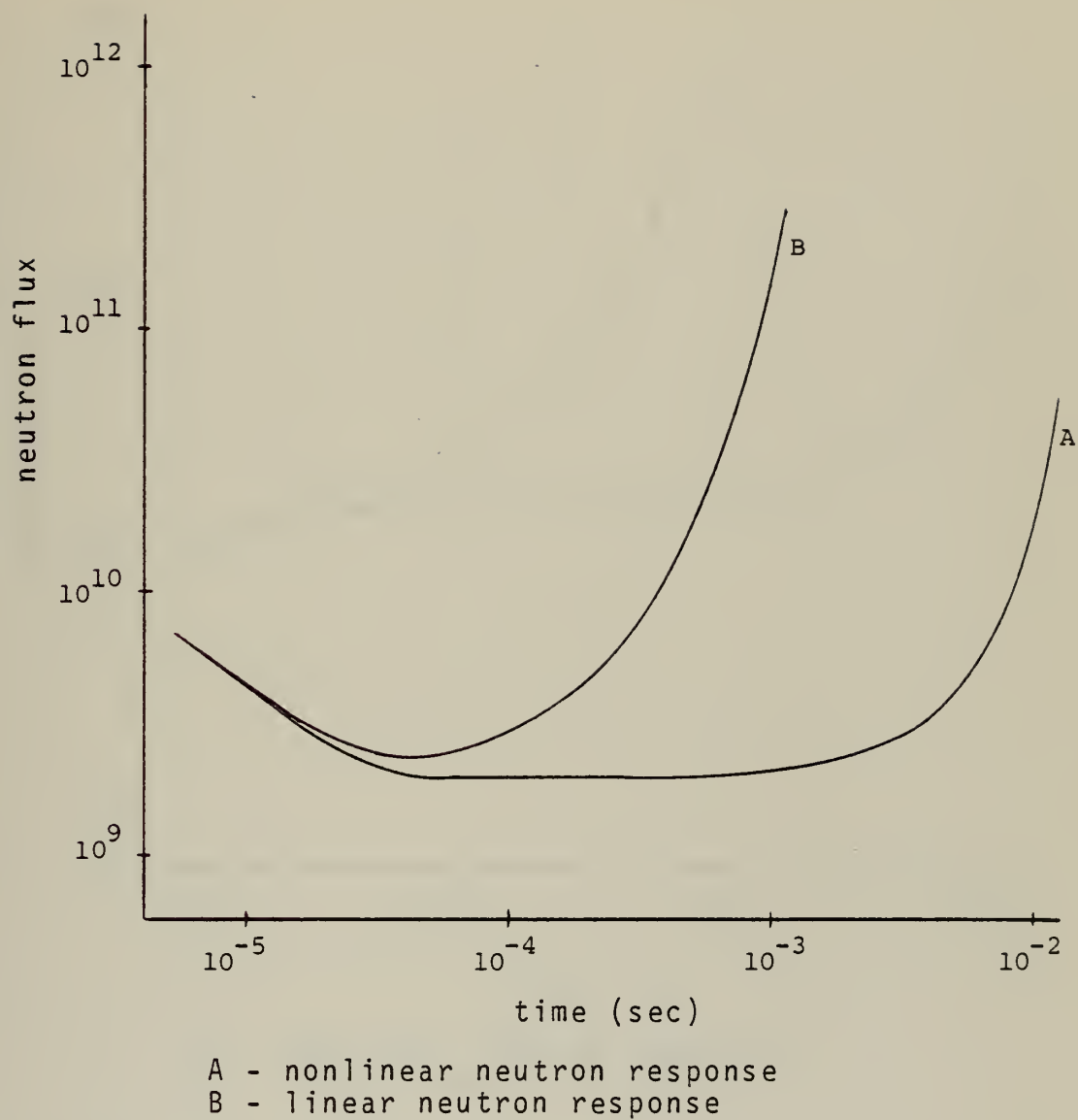
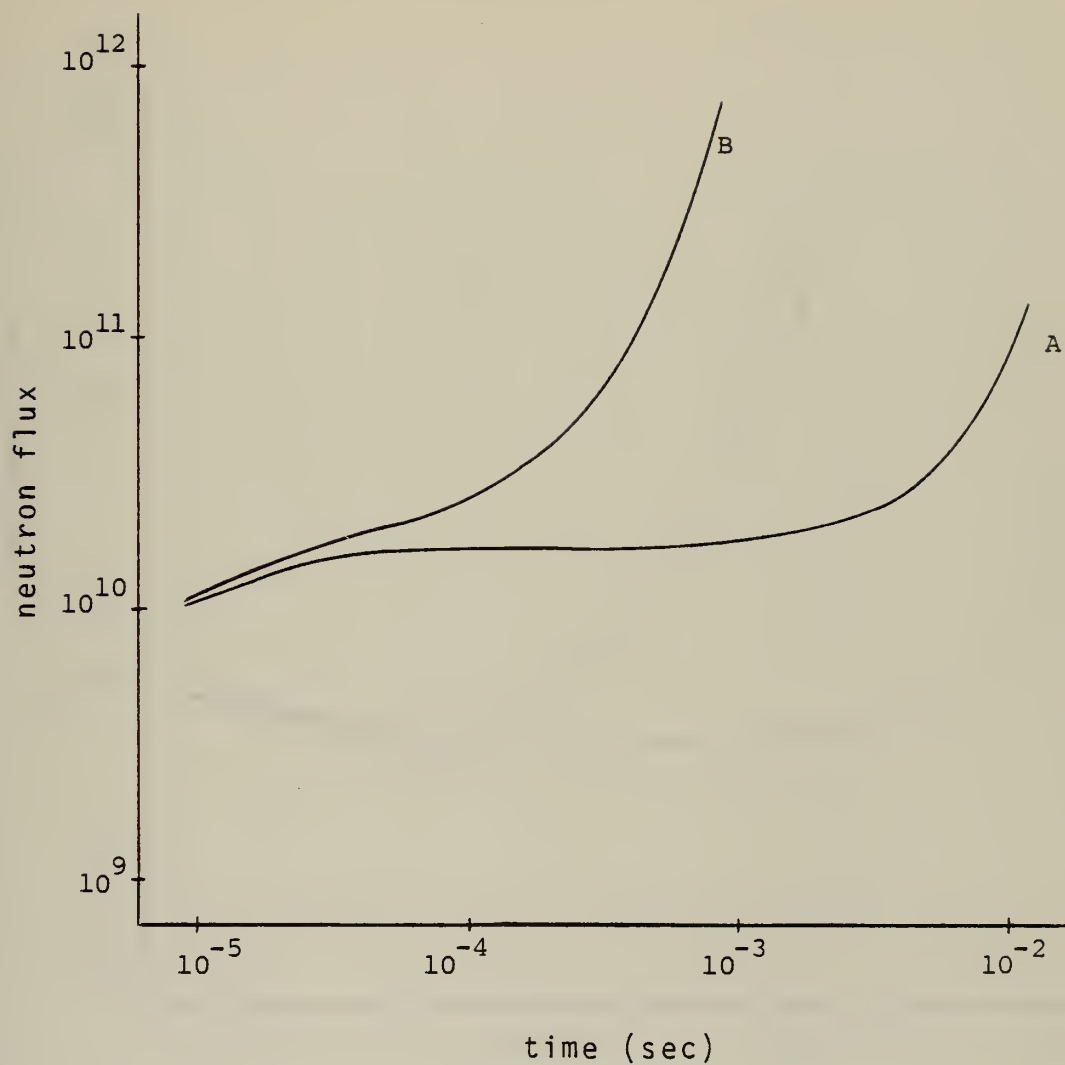


Figure 17. Linear and nonlinear fluxes at  $(-60, 0, 80)$  due to a local perturbation of 100 dollar/sec at  $(0, 60, 40)$ .





A - nonlinear neutron response  
B - linear neutron response

Figure 18. Linear and nonlinear fluxes at (0,0,0) due to a 50 dollar/sec local perturbation at (0,60,40).



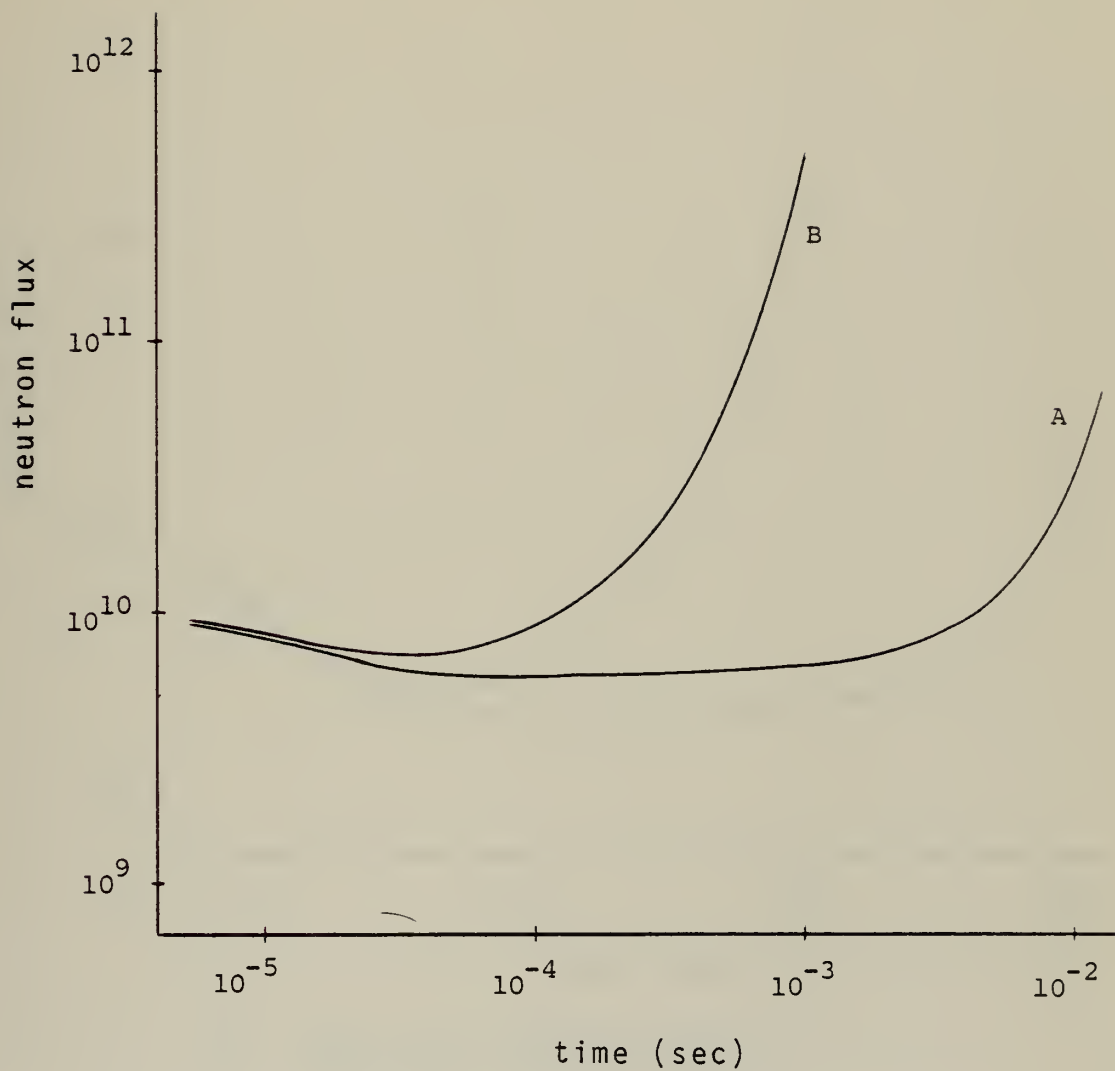


Figure 19. Linear and nonlinear fluxes at (60,0,0) due to a 50 dollar/sec local perturbation at (0,60,40).





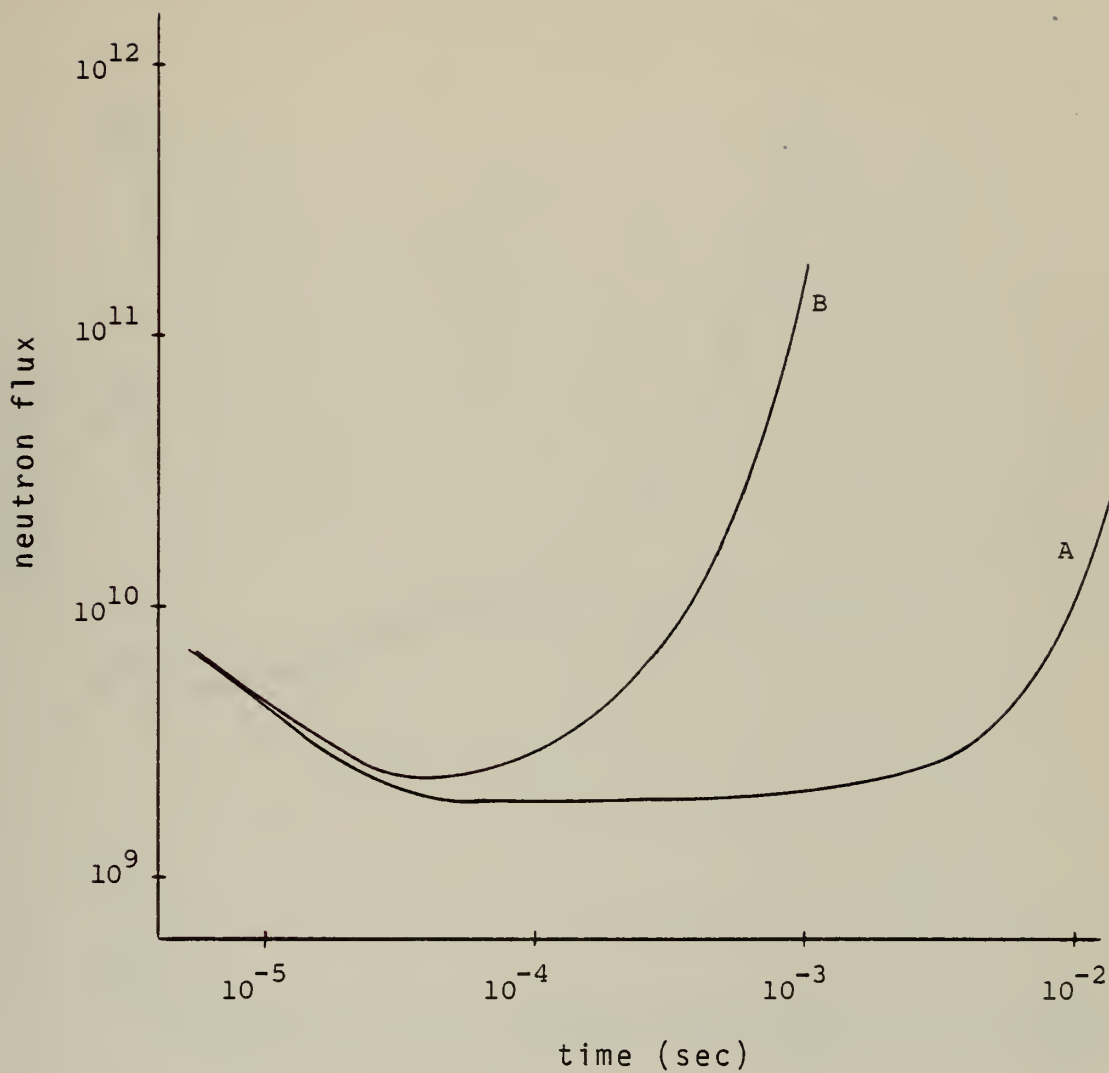
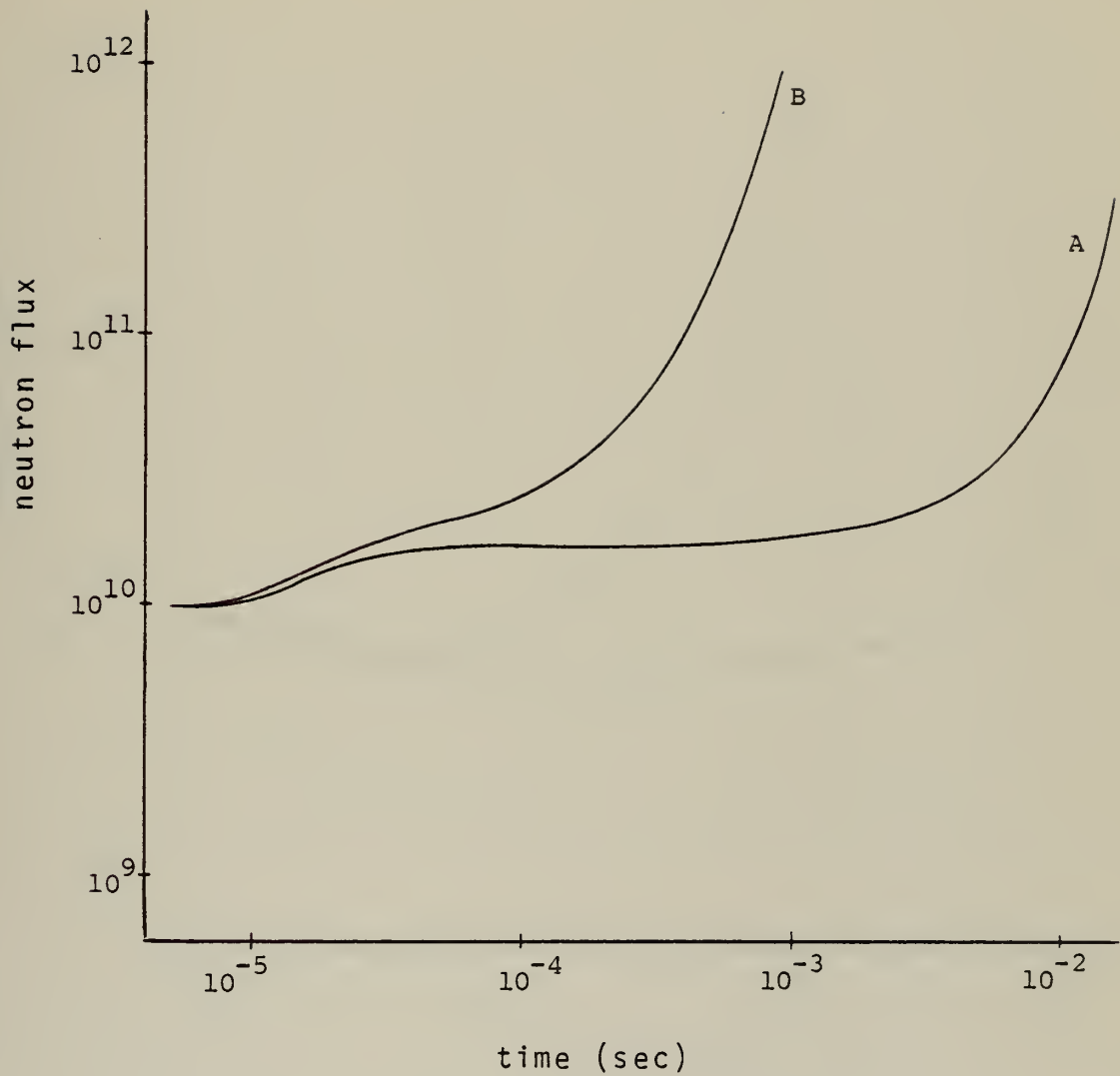


Figure 20. Linear and nonlinear fluxes at  $(-60,0,80)$  due to a 50 dollar/sec local perturbation at  $(0,60,40)$ .

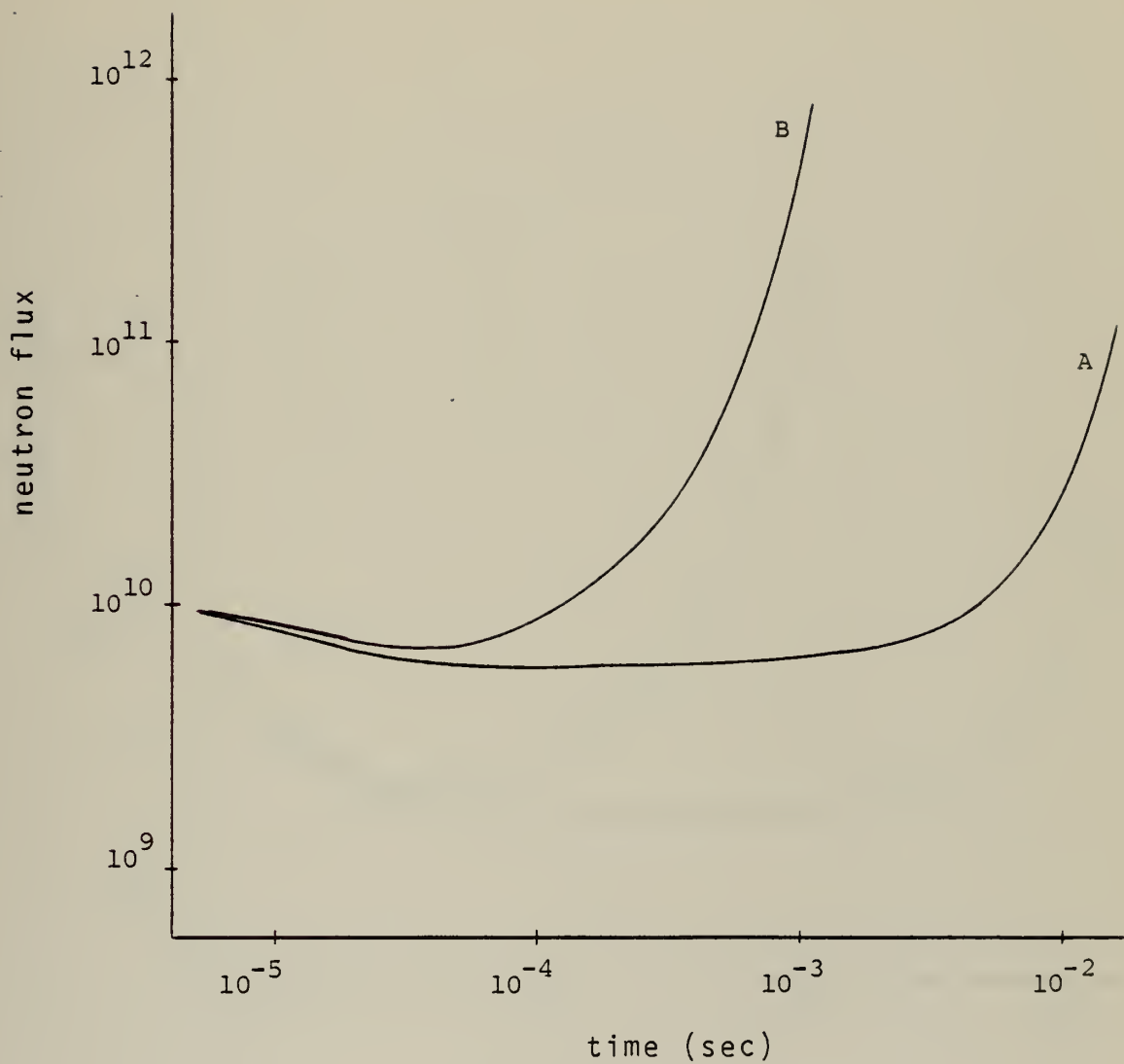




A - nonlinear neutron response  
B - linear neutron response

Figure 21. Linear and nonlinear fluxes at (0,0,0) due to a 10 dollar/sec local perturbation at (0,60,40).





A - nonlinear neutron response  
B - linear neutron response

Figure 22. Linear and nonlinear fluxes at (60,0,0) due to a 10 dollar/sec local perturbation at (0,60,40).



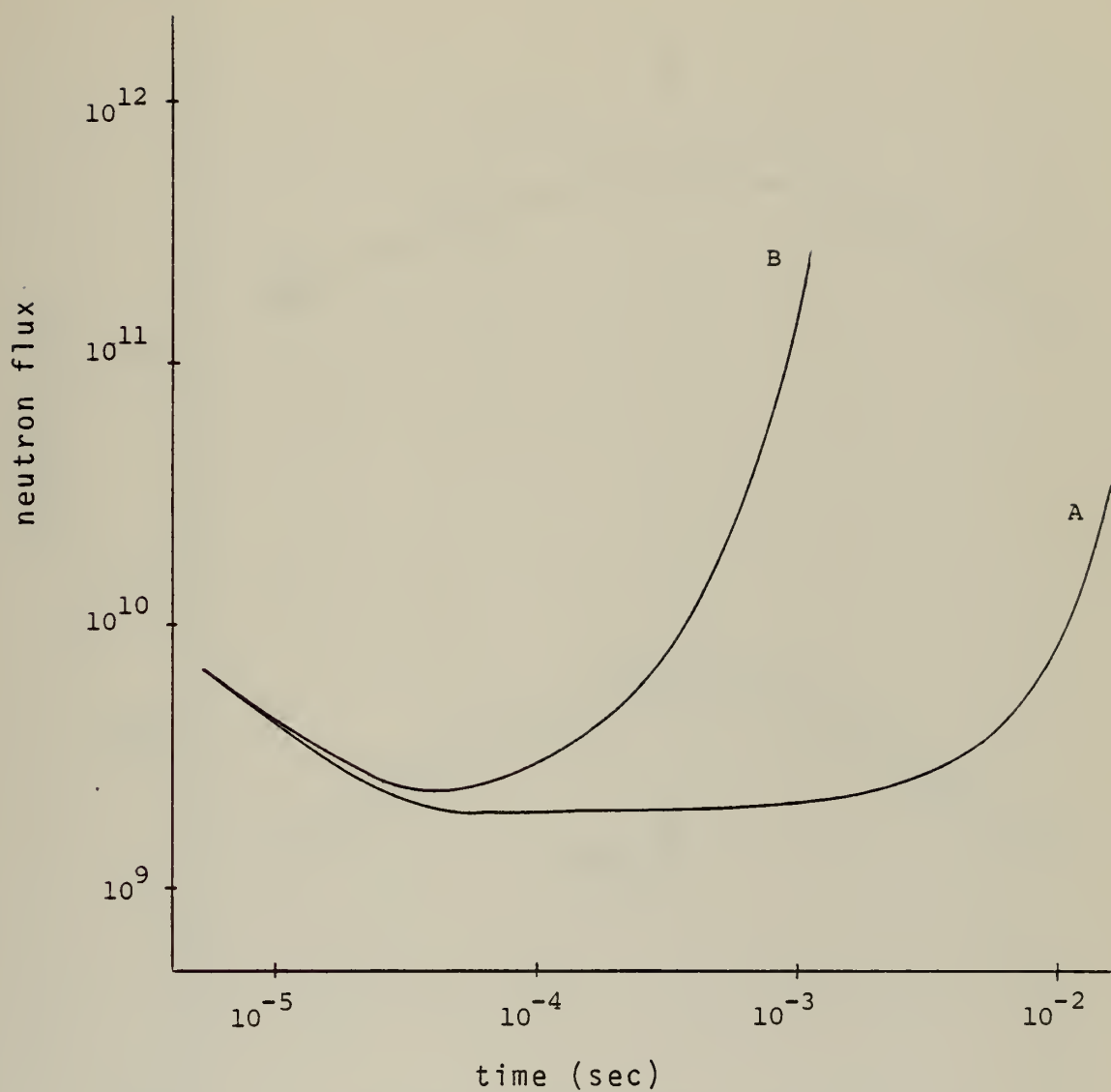
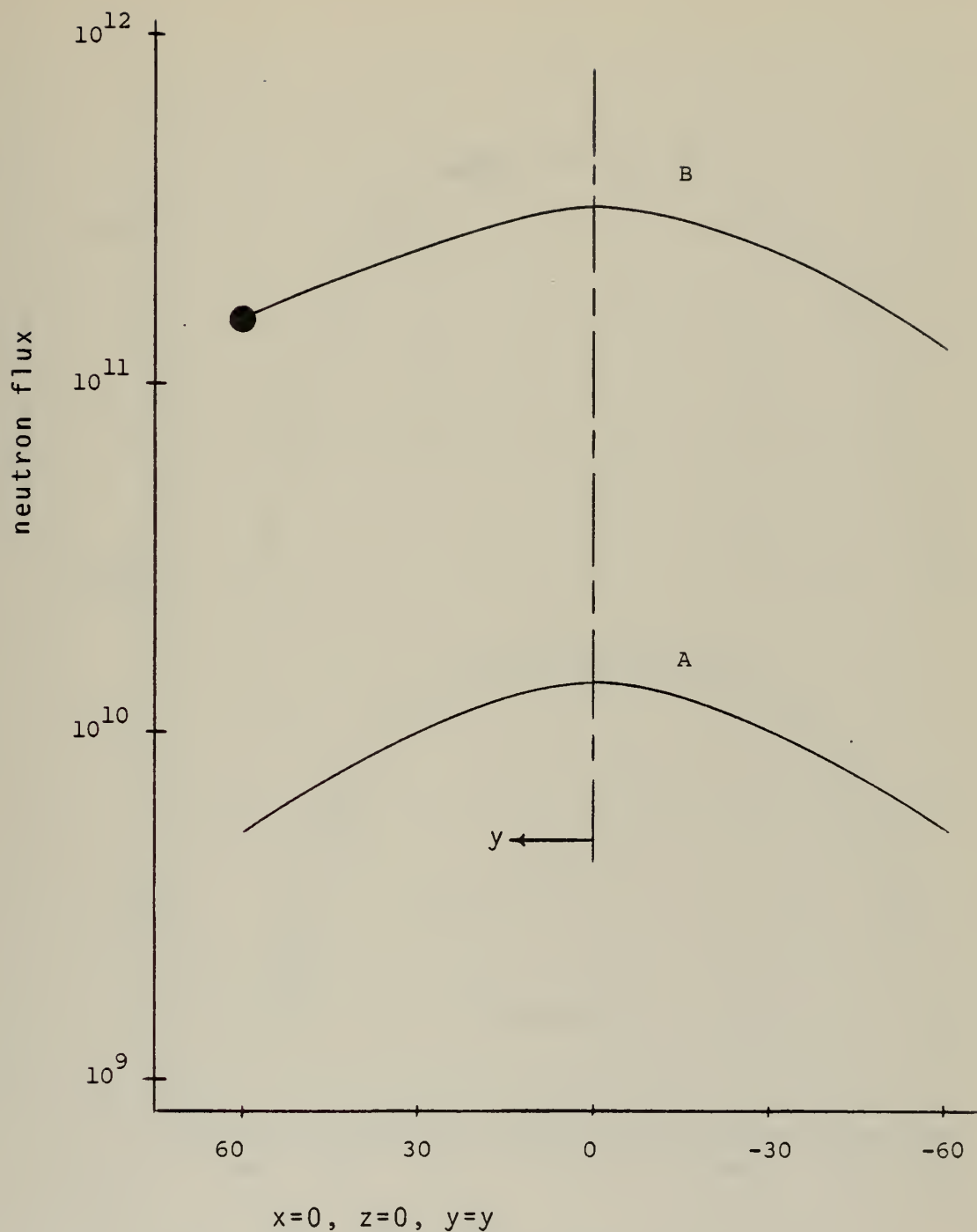


Figure 23. Linear and nonlinear fluxes at  $(-60, 0, 80)$  due to a 10 dollar/sec local perturbation at  $(0, 60, 40)$ .



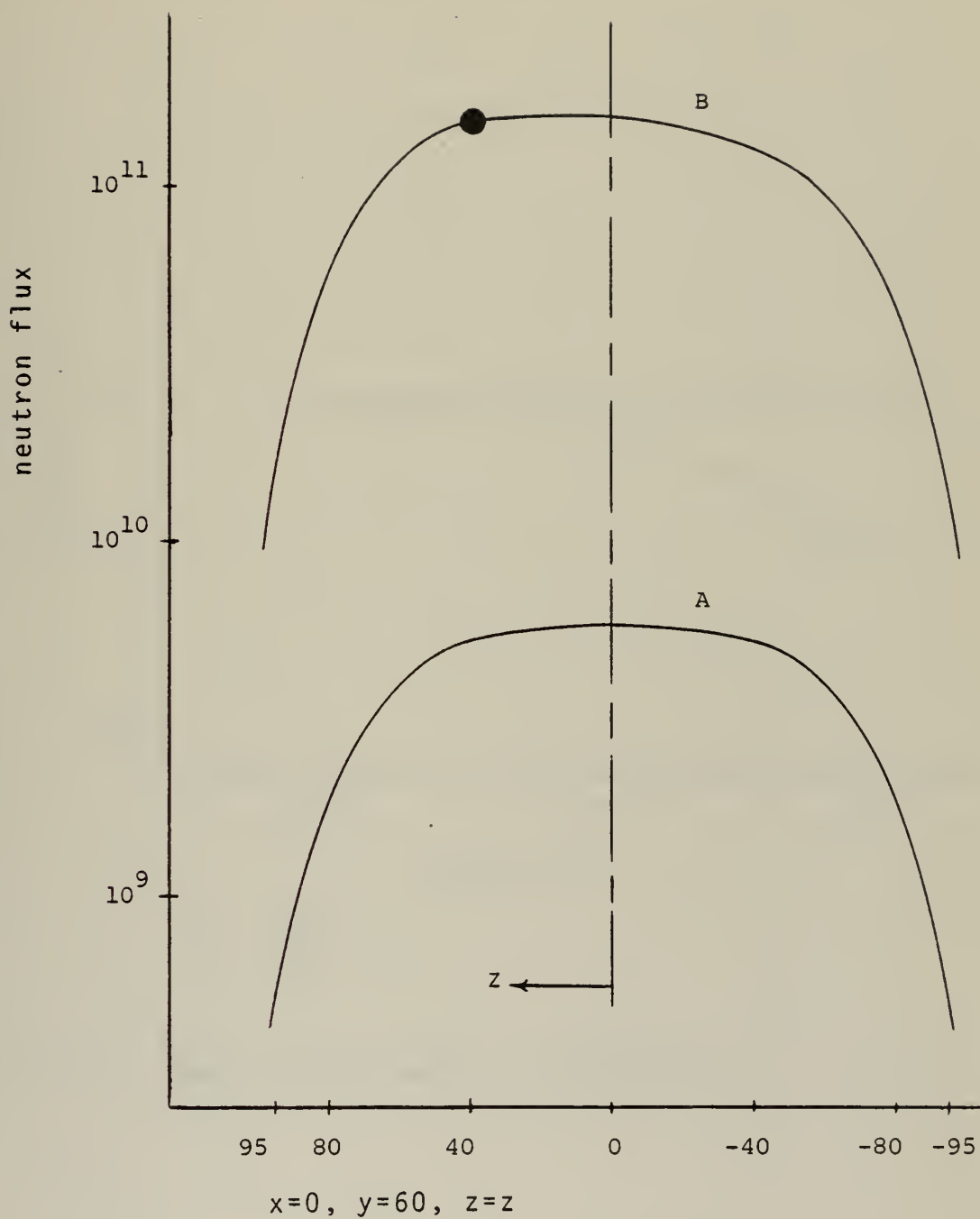




- - point of local perturbation
- A - neutron flux at steady-state
- B - neutron flux under local perturbation

Figure 24. Radial flux distribution for the steady state and 100 dollar/sec local perturbation.





- - point of local perturbation
- A - axial flux distribution during steady state
- B - axial flux distribution during perturbation

Figure 25. Axial flux distribution for the steady state and 100 dollar/sec local perturbation.



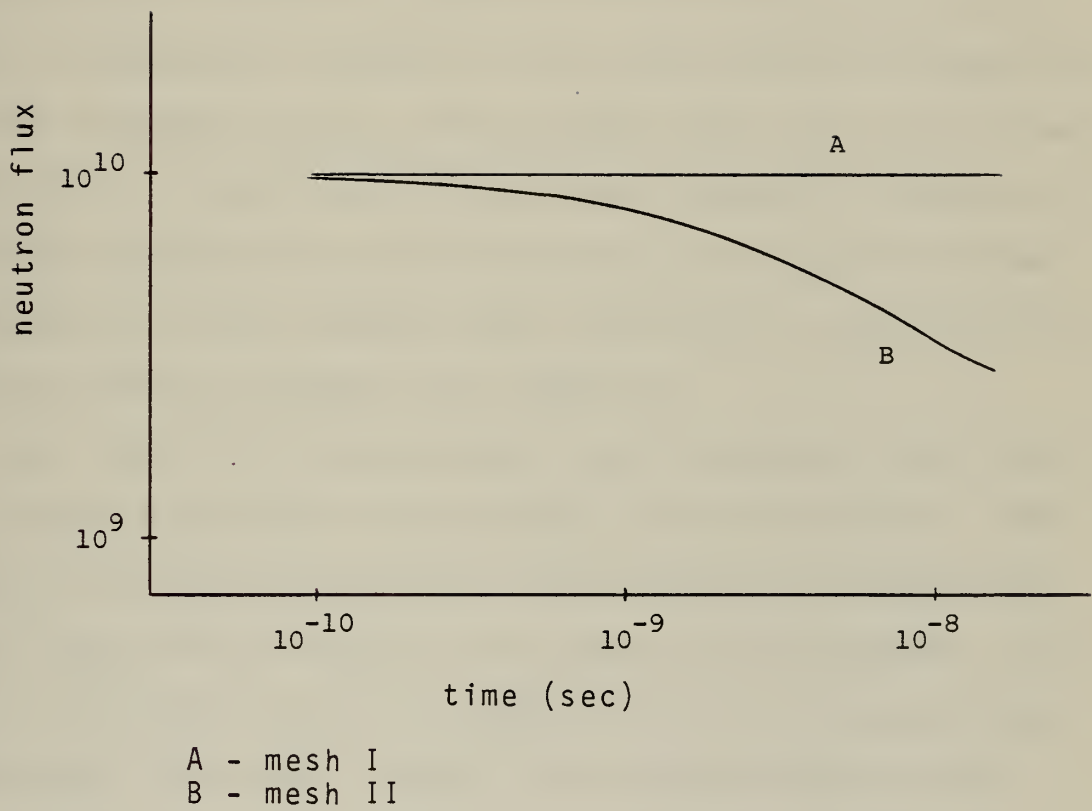


Figure 26. Early time history of the neutron flux at core center using mesh I and mesh II.



and delayed neutrons on the flux. To obtain an idea of the difference in magnitude between the linear and nonlinear flux for the 100 dollar local perturbation, at time  $t = 10^{-3}$  second, the linear case predicted a neutron flux at the core center of  $2.024 \times 10^{12}$  neutrons per  $\text{cm}^2$  per sec. The nonlinear reactor predicted  $1.85 \times 10^{10}$  neutrons per  $\text{cm}^2$  per sec. Although the numbers are not sufficiently accurate due to the crudeness of the finite element mesh, the significant difference in the order of magnitude between the linear and nonlinear neutron flux leads to the belief that the three-dimensional finite element model utilized is predicting the correct trend of neutron flux behavior.

The radial flux distribution for the steady state shown in figure 24 portrays the expected flux distribution. The radial flux distributions for the local perturbation case show a skew distribution at the point of perturbation. The axial flux distribution at steady state is very much symmetric about the center. Again, the expected skew at the point of perturbation is there, as shown in figure 25. Figure 26 gives an indication that the  $\Sigma_f^*$  for finer meshes is higher. However, curve B of figure 26 should be extended to longer times to verify this hypothesis. Curve B, as plotted in figure 26, used four hours of computer time employing the H-compiler of the IBM 360/67.





## VI. CONCLUSIONS AND RECOMMENDATIONS

The finite element mesh employed here is crude, and, thus, results that were obtained should be considered as positive indicators rather than numerically conclusive facts. The trend of neutron flux behavior is the major thrust of this work. From the results obtained, it is concluded that the expected patterns of neutron behavior as predicted by the three-dimensional finite element model used do occur. These patterns were best demonstrated by the differences between the linear and nonlinear flux responses and by the spatial flux distribution at steady state and during local perturbation. The three-dimensional quadratic finite element model utilized should produce better results by resorting to a finer mesh. The draw-back to a finer mesh, of course, is the significant increase in computer time and storage requirements. Once accurate results are obtained through finer element meshes, comparisons between three- and two-dimensional models can be attempted.

A mesh generator was not developed for this problem. It is recommended that this be done to ease the transition from one mesh to another and to minimize human error. In addition, a similar calculation using a three-dimensional linear element should be performed to corroborate the results obtained here. It should be particularly noted that this type of problem is highly sensitive to the fission cross-section.



In the search for  $\Sigma_f^*$ , it was necessary to adjust  $\Sigma_f$  to the sixth decimal place. There is no exact method of deriving  $\Sigma_f^*$  due to the highly nonlinear aspect of the problem; therefore, trial and error must be used. Finally, the Gaussian quadrature used to determine the element matrices should be investigated. The cause of the differences in values obtained by using different number of integration points should be established. Was it due to the integration points, the coordinate transformation, or the element shape functions? This is an important question upon which future numerical results could be based.



# APPENDIX A

## MESH I CONNECTIVITY AND COORDINATES

The connectivity matrix for the 128-element mesh is as follows:

ELEMENT NUMBER	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15
1	1	3	12	11	10	2	26	28	27	35	37	46	45	44	36
2	1	4	14	13	12	3	26	29	28	35	38	48	47	46	37
3	1	5	15	14	13	4	26	30	29	35	39	50	49	48	38
4	1	6	16	15	14	5	26	31	30	35	40	52	51	50	39
5	1	7	17	16	15	6	26	32	31	35	41	54	53	52	40
6	1	8	18	17	16	7	26	33	32	35	42	56	55	54	41
7	1	9	19	18	17	8	26	34	33	35	43	58	57	56	42
8	1	10	20	19	18	9	26	35	34	35	36	44	43	42	43
9	1	11	21	20	19	10	26	36	35	35	37	80	79	78	70
10	3	37	45	44	43	36	60	62	61	69	72	82	81	80	71
11	3	38	46	45	44	37	60	63	62	69	73	84	83	82	72
12	3	39	47	46	45	38	60	64	63	69	74	86	85	84	73
13	3	40	48	47	46	39	60	65	64	69	75	88	87	86	74
14	3	41	49	48	47	40	60	66	65	69	76	90	89	88	75
15	3	42	50	49	48	41	60	67	66	69	77	92	91	90	76
16	3	43	51	50	49	42	60	68	67	69	78	93	92	91	77
17	3	44	52	51	50	43	60	69	68	69	79	94	93	92	78
18	3	45	53	52	51	44	60	70	69	69	80	95	94	93	79
19	3	46	54	53	52	45	60	71	70	69	81	96	95	94	80
20	3	47	55	54	53	46	60	72	71	69	82	97	96	95	81
21	3	48	56	55	54	47	60	73	72	69	83	98	97	96	82
22	3	49	57	56	55	48	60	74	73	69	84	99	98	97	83
23	3	50	58	57	56	49	60	75	74	69	85	100	99	98	84
24	3	51	59	58	57	50	60	76	75	69	86	101	100	99	85
25	3	52	60	59	58	51	60	77	76	69	87	102	101	100	86
26	3	53	61	60	59	52	60	78	77	69	88	103	102	101	87
27	3	54	62	61	60	53	60	79	78	69	89	104	103	102	88
28	3	55	63	62	61	54	60	80	79	69	90	105	104	103	89
29	3	56	64	63	62	55	60	81	80	69	91	106	105	104	90
30	3	57	65	64	63	56	60	82	81	69	92	107	106	105	91
31	3	58	66	65	64	57	60	83	82	69	93	108	107	106	92
32	3	59	67	66	65	58	60	84	83	69	94	109	108	107	93
33	3	60	68	67	66	59	60	85	84	69	95	110	109	108	94
34	3	61	69	68	67	60	60	86	85	69	96	111	110	109	95
35	3	62	70	69	68	61	60	87	86	69	97	112	111	110	96
36	3	63	71	70	69	62	60	88	87	69	98	113	112	111	97
37	3	64	72	71	70	63	60	89	88	69	99	114	113	112	98
38	3	65	73	72	71	64	60	90	89	69	100	115	114	113	99
39	3	66	74	73	72	65	60	91	90	69	101	116	115	114	100
40	3	67	75	74	73	66	60	92	91	69	102	117	116	115	101
41	3	68	76	75	74	67	60	93	92	69	103	118	117	116	102
42	3	69	77	76	75	68	60	94	93	69	104	119	118	117	103
43	3	70	78	77	76	69	60	95	94	69	105	120	119	118	104
44	3	71	79	78	77	70	60	96	95	69	106	121	120	119	105
45	3	72	80	79	78	71	60	97	96	69	107	122	121	120	106
46	3	73	81	80	79	72	60	98	97	69	108	123	122	121	107
47	3	74	82	81	80	73	60	99	98	69	109	124	123	122	108
48	3	75	83	82	81	74	60	100	99	69	110	125	124	123	109
49	3	76	84	83	82	75	60	101	100	69	111	126	125	124	110
50	3	77	85	84	83	76	60	102	101	69	112	127	126	125	111
51	3	78	86	85	84	77	60	103	102	69	113	128	127	126	112
52	3	79	87	86	85	78	60	104	103	69	114	129	128	127	113
53	3	80	88	87	86	79	60	105	104	69	115	130	129	128	114
54	3	81	89	88	87	80	60	106	105	69	116	131	130	129	115
55	3	82	90	89	88	81	60	107	106	69	117	132	131	130	116
56	3	83	91	90	89	82	60	108	107	69	118	133	132	131	117
57	3	84	92	91	90	83	60	109	108	69	119	134	133	132	118
58	3	85	93	92	91	84	60	110	109	69	120	135	134	133	119
59	3	86	94	93	92	85	60	111	110	69	121	136	135	134	120
60	3	87	95	94	93	86	60	112	111	69	122	137	136	135	121
61	3	88	96	95	94	87	60	113	112	69	123	138	137	136	122
62	3	89	97	96	95	88	60	114	113	69	124	139	138	137	123
63	3	90	98	97	96	89	60	115	114	69	125	140	139	138	124
64	3	91	99	98	97	90	60	116	115	69	126	141	140	139	125
65	3	92	100	99	98	91	60	117	116	69	127	142	141	140	126
66	3	93	101	100	99	92	60	118	117	69	128	143	142	141	127
67	3	94	102	101	100	93	60	119	118	69	129	144	143	142	128
68	3	95	103	102	101	94	60	120	119	69	130	145	144	143	129
69	3	96	104	103	102	95	60	121	120	69	131	146	145	144	130
70	3	97	105	104	103	96	60	122	121	69	132	147	146	145	131
71	3	98	106	105	104	97	60	123	122	69	133	148	147	146	132
72	3	99	107	106	105	98	60	124	123	69	134	149	148	147	133
73	3	100	108	107	106	99	60	125	124	69	135	150	149	148	134
74	3	101	109	108	107	100	60	126	125	69	136	151	150	149	135
75	3	102	110	109	108	101	60	127	126	69	137	152	151	150	136
76	3	103	111	110	109	102	60	128	127	69	138	153	152	151	137
77	3	104	112	111	110	103	60	129	128	69	139	154	153	152	138
78	3	105	113	112	111	104	60	130	129	69	140	155	154	153	139
79	3	106	114	113	112	105	60	131	130	69	141	156	155	154	140
80	3	107	115	114	113	106	60	132	131	69	142	157	156	155	141
81	3	108	116	115	114	107	60	133	132	69	143	158	157	156	142
82	3	109	117	116	115	108	60	134	133	69	144	159	158	157	143
83	3	110	118	117	116	109	60	135	134	69	145	160	159	158	144
84	3	111	119	118	117	110	60	136	135	69	146	161	160	159	145
85	3	112	120	119	118	111	60	137	136	69	147	162	161	160	146
86	3	113	121	120	119	112	60	138	137	69	148	163	162	161	147
87	3	114	122	121	120	113	60	139	138	69	149	164	163	162	148
88	3	115	123	122	121	114	60	140	139	69	150	165	164	163	149
89	3	116	124	123	122	115	60	141	140	69	151	166	165	164	150
90	3	117	125	124	123	116	60	142	141	69	152	167	166	165	151
91	3	118	126	125	124	117	60	143	142	69	153	168	167	166	152
92	3	119	127	126	125	118	60	144	143	69	154	169	168	167	153
93	3	120	128	127	126	119	60	145	144	69	155	170	169	168	154
94	3	121	129	128	127	120	60	146	145	69	156	171	170	169	155
95	3	122	130	129	128	121	60	147	146	69	157	172	171	170	156
96	3	123	131	130	129	122	60	148	147	69	158	173	172	171	157
97	3	124	132	131	130	123	60	149	148	69	159	174	173	172	158
98	3	125	133	132	131	124	60	150	149	69	160	175	174	173	159
99	3	126	134	133	132	125	60								





[illegible]





[illegible]



The coordinates of the 128-element mesh are as follows:

NODE NUMBER	X	Y	Z
1	0.0	0.0	80.00
2	0.0	30.00	80.00
3	21.21	21.21	80.00
4	30.00	0.0	80.00
5	21.21	-21.21	80.00
6	0.0	-30.00	80.00
7	-21.21	-21.21	80.00
8	-30.00	0.0	80.00
9	-21.21	21.21	80.00
10	0.0	60.00	80.00
11	22.96	55.43	80.00
12	42.43	42.43	80.00
13	55.43	22.96	80.00
14	60.00	0.0	80.00
15	55.43	-22.96	80.00
16	42.43	-42.43	80.00
17	22.96	-55.43	80.00
18	0.0	-60.00	80.00
19	-22.96	-55.43	80.00
20	-42.43	-42.43	80.00
21	-55.43	-22.96	80.00
22	-60.00	0.0	80.00
23	-55.43	22.96	80.00
24	-42.43	42.43	80.00
25	-22.96	55.43	80.00
26	0.0	60.00	40.00
27	0.0	42.43	40.00
28	42.43	0.0	40.00
29	60.00	0.0	40.00
30	42.43	-42.43	40.00
31	0.0	-60.00	40.00
32	-42.43	-42.43	40.00
33	-60.00	0.0	40.00
34	-55.43	22.96	40.00
35	-42.43	42.43	40.00
36	-22.96	55.43	40.00
37	0.0	60.00	0.0
38	21.21	21.21	0.0
39	30.00	0.0	0.0
40	21.21	-21.21	0.0
41	0.0	-30.00	0.0
42	-21.21	-21.21	0.0
43	-30.00	0.0	0.0
44	-21.21	21.21	0.0



45	22.43	55.43	00.00	00.00
46	22.43	42.96	40.00	00.00
47	55.00	22.96	40.00	00.00
48	55.43	-42.43	40.00	00.00
49	55.43	-55.43	40.00	00.00
50	42.96	-55.00	40.00	00.00
51	22.96	-42.96	40.00	00.00
52	22.43	-55.43	40.00	00.00
53	22.43	-42.43	40.00	00.00
54	55.43	-55.43	40.00	00.00
55	55.00	-42.96	40.00	00.00
56	42.96	22.96	40.00	00.00
57	55.43	42.43	40.00	00.00
58	55.43	55.43	40.00	00.00
59	22.96	60.00	40.00	00.00
60	00.00	42.43	40.00	00.00
61	42.43	40.00	40.00	00.00
62	60.00	40.43	40.00	00.00
63	42.00	-42.43	40.00	00.00
64	42.00	-60.00	40.00	00.00
65	42.43	-42.43	40.00	00.00
66	60.00	42.43	40.00	00.00
67	42.00	40.00	40.00	00.00
68	42.43	30.00	80.00	00.00
69	00.00	21.21	80.00	00.00
70	21.21	21.00	80.00	00.00
71	30.00	21.21	80.00	00.00
72	21.00	-21.21	80.00	00.00
73	21.21	-30.00	80.00	00.00
74	00.00	-21.21	80.00	00.00
75	21.21	-30.00	80.00	00.00
76	00.00	-21.21	80.00	00.00
77	21.21	21.00	80.00	00.00
78	00.00	55.43	80.00	00.00
79	22.96	42.43	80.00	00.00
80	42.43	55.43	80.00	00.00
81	55.00	60.00	80.00	00.00
82	55.43	42.96	80.00	00.00
83	42.96	-42.96	80.00	00.00
84	22.43	-55.43	80.00	00.00
85	22.43	-55.00	80.00	00.00
86	00.00	-42.96	80.00	00.00
87	22.96	-55.43	80.00	00.00
88	42.43	-42.43	80.00	00.00
89	55.43	-55.43	80.00	00.00
90	60.00	-60.00	80.00	00.00
91	42.43	22.96	80.00	00.00
92	42.43	22.43	80.00	00.00



93	-22.96	55.43	-80.00	00
94	16.80	75.00	-80.00	00
95	38.10	71.60	-80.00	00
96	53.03	62.90	-80.00	00
97	62.80	53.03	-80.00	00
98	71.30	38.60	-80.00	00
99	75.00	17.00	-80.00	00
100	71.30	17.00	-80.00	00
101	62.80	-38.60	-80.00	00
102	53.03	-53.03	-80.00	00
103	38.10	-62.90	-80.00	00
104	17.10	-71.80	-80.00	00
105	17.10	-75.00	-80.00	00
106	0.90	-71.80	-80.00	00
107	-16.50	-62.80	-80.00	00
108	-38.50	-53.03	-80.00	00
109	-53.03	-38.60	-80.00	00
110	-62.70	-17.10	-80.00	00
111	-71.30	-17.10	-80.00	00
112	-75.00	17.00	-80.00	00
113	-71.30	17.10	-80.00	00
114	-62.90	38.30	-80.00	00
115	-53.03	53.03	-80.00	00
116	-38.60	62.80	-80.00	00
117	-17.30	71.30	-80.00	00
118	-17.30	71.30	-95.00	00
119	0.00	60.00	-95.00	00
120	42.43	42.43	-95.00	00
121	50.00	0.00	-95.00	00
122	42.43	42.43	-95.00	00
123	42.43	42.43	-95.00	00
124	42.43	42.43	-95.00	00
125	42.43	42.43	-95.00	00
126	42.43	42.43	-95.00	00
127	42.43	42.43	-95.00	00
128	16.80	75.00	0.00	00
129	38.10	71.60	0.00	00
130	53.03	62.90	0.00	00
131	62.80	53.03	0.00	00
132	71.30	38.60	0.00	00
133	75.00	17.00	0.00	00
134	71.30	17.00	0.00	00
135	62.80	-38.60	0.00	00
136	53.03	-53.03	0.00	00
137	38.10	-62.90	0.00	00
138	17.10	-71.80	0.00	00
139	17.10	-75.00	0.00	00
140	-16.90	-71.80	0.00	00





141	-38.50	-62.80	0.0
142	-53.70	-53.03	0.0
143	-62.70	-38.60	0.0
144	-71.30	-17.10	0.0
145	-75.00	17.00	0.0
146	-71.30	17.10	0.0
147	-62.90	38.03	0.0
148	-53.90	53.03	0.0
149	-38.60	62.80	0.0
150	-17.30	71.30	0.0
151	16.80	75.00	80.00
152	38.10	71.60	80.00
153	53.03	62.90	80.00
154	62.80	53.03	80.00
155	71.30	38.60	80.00
156	75.00	17.00	80.00
157	71.30	17.00	80.00
158	62.80	-17.00	80.00
159	53.90	-38.60	80.00
160	38.10	-53.03	80.00
161	17.00	-62.90	80.00
162	16.90	-71.80	80.00
163	-38.50	-75.00	80.00
164	-53.03	-71.80	80.00
165	-62.70	-62.03	80.00
166	-71.30	-38.60	80.00
167	-75.00	-17.10	80.00
168	-71.30	17.00	80.00
169	-62.90	17.10	80.00
170	-53.90	38.03	80.00
171	-38.60	53.03	80.00
172	-17.30	62.80	80.00
173	16.80	71.30	80.00
174	38.10	75.00	95.00
175	53.03	60.00	95.00
176	62.80	42.43	95.00
177	71.30	42.00	95.00
178	75.00	42.43	95.00
179	71.30	42.00	95.00
180	62.90	42.43	95.00
181	53.90	42.00	95.00
182	38.60	42.43	95.00
183	17.30	42.00	95.00
184	16.80	42.43	95.00
185	38.10	30.00	110.00
186	53.03	21.21	110.00
187	62.80	21.00	110.00
188	71.30	-21.21	110.00



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## MESH II CONNECTIVITY AND COORDINATES

The connectivity matrix for the 192-element mesh is as follows:

[illegible]



304	305	306	307	308	223	4	5	6	7	8	9	0	180	182	182	183	185	186	188	189	191	192	194	195	197	197	198	200	201	203	204	206	207	209	210	212	213	215																																																																																																																																																																																																																																																																																																																																																																																				
315	317	319	321	323	10	12	14	16	18	20	22	24	589	591	593	594	595	597	599	601	603	605	607	609	621	613	615	617	619	654	616	635	645	648	649	650	651	652	653	655																																																																																																																																																																																																																																																																																																																																																																																		
316	318	320	322	324	11	13	15	17	19	21	23	25	590	591	592	594	596	598	599	602	604	606	608	610	612	614	616	618	620	622	624	626	628	630	632	634	636	638	640	642	644	646	648	650	652	654	656	658	660	662	664	666	668	670	672	674	676	678	680	682	684	686	688	690	692	694	696	698	700	702	704	706	708	710	712	714	716	718	720	722	724	726	728	730	732	734	736	738	740	742	744	746	748	750	752	754	756	758	760	762	764	766	768	770	772	774	776	778	780	782	784	786	788	790	792	794	796	798	800	802	804	806	808	810	812	814	816	818	820	822	824	826	828	830	832	834	836	838	840	842	844	846	848	850	852	854	856	858	860	862	864	866	868	870	872	874	876	878	880	882	884	886	888	890	892	894	896	898	900	902	904	906	908	910	912	914	916	918	920	922	924	926	928	930	932	934	936	938	940	942	944	946	948	950	952	954	956	958	960	962	964	966	968	970	972	974	976	978	980	982	984	986	988	990	992	994	996	998	1000	1002	1004	1006	1008	1010	1012	1014	1016	1018	1020	1022	1024	1026	1028	1030	1032	1034	1036	1038	1040	1042	1044	1046	1048	1050	1052	1054	1056	1058	1060	1062	1064	1066	1068	1070	1072	1074	1076	1078	1080	1082	1084	1086	1088	1090	1092	1094	1096	1098	1100	1102	1104	1106	1108	1110	1112	1114	1116	1118	1120	1122	1124	1126	1128	1130	1132	1134	1136	1138	1140	1142	1144	1146	1148	1150	1152	1154	1156	1158	1160	1162	1164	1166	1168	1170	1172	1174	1176	1178	1180	1182	1184	1186	1188	1190	1192	1194	1196	1198	1200	1202	1204	1206	1208	1210	1212	1214	1216	1218	1220	1222	1224	1226	1228	1230	1232	1234	1236	1238	1240	1242	1244	1246	1248	1250	1252	1254	1256	1258	1260	1262	1264	1266	1268	1270	1272	1274	1276	1278	1280	1282	1284	1286	1288	1290	1292	1294	1296	1298	1300	1302	1304	1306	1308	1310	1312	1314	1316	1318	1320	1322	1324	1326	1328	1330	1332	1334	1336	1338	1340	1342	1344	1346	1348	1350	1352	1354	1356	1358	1360	1362	1364	1366	1368	1370	1372	1374	1376	1378	1380	1





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180	154	288	413	412	411	287	167	389	388	317	337	365	364	337
181	155	289	414	414	413	288	167	390	389	317	338	365	366	338
182	156	290	415	415	414	290	390	167	389	317	339	365	368	339
183	157	291	416	416	415	291	168	391	390	319	340	367	368	340
184	158	292	417	417	416	292	168	392	391	319	341	369	370	342
185	159	293	418	418	417	293	169	393	392	321	342	371	372	343
186	160	294	419	419	418	294	169	394	393	321	343	371	374	344
187	161	295	420	420	419	295	170	395	394	321	344	373	374	345
188	162	296	421	421	420	296	170	396	395	321	345	373	376	345
189	163	297	422	422	421	297	170	396	395	321	346	375	378	346
190	164	298	423	423	422	298	170	396	395	321	347	377	378	347
191	165	299	424	424	423	299	170	396	395	321	347	377	378	348
192	166	300	425	425	424	300	170	396	395	321	347	377	378	348

The coordinates of the 192-element mesh are as follows:

NODE NUMBER	X	Y	Z
1	0.0	0.0	80.00
2	0.0	30.21	80.00
3	21.00	21.00	80.00
4	21.00	0.0	80.00
5	0.0	-21.00	80.00
6	0.0	-30.21	80.00
7	-21.00	-21.00	80.00
8	-21.00	0.0	80.00
9	-30.21	21.00	80.00
10	-30.21	60.43	80.00
11	0.0	55.43	80.00
12	22.96	42.96	80.00
13	22.96	22.96	80.00
14	55.43	0.0	80.00
15	55.43	-22.96	80.00
16	42.96	-42.96	80.00
17	42.96	-55.43	80.00
18	0.0	-60.43	80.00
19	0.0	-55.43	80.00
20	22.96	-42.96	80.00
21	22.96	-22.96	80.00
22	55.43	0.0	80.00
23	55.43	22.96	80.00
24	42.96	42.96	80.00
25	42.96	55.43	80.00
26	0.0	60.43	60.00



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75	-21.21	0.00	0.00
76	-30.00	0.00	0.00
77	-21.21	0.00	0.00
78	0.00	0.00	0.00
79	22.96	0.00	0.00
80	42.43	0.00	0.00
81	55.00	0.00	0.00
82	55.43	0.00	0.00
83	42.43	0.00	0.00
84	2.00	0.00	0.00
85	-22.96	0.00	0.00
86	-42.43	0.00	0.00
87	-55.00	0.00	0.00
88	-55.43	0.00	0.00
89	-42.43	0.00	0.00
90	-2.00	0.00	0.00
91	22.96	0.00	0.00
92	42.43	0.00	0.00
93	55.00	0.00	0.00
94	55.43	0.00	0.00
95	42.43	0.00	0.00
96	2.00	0.00	0.00
97	-22.96	0.00	0.00
98	-42.43	0.00	0.00
99	-55.00	0.00	0.00
100	-55.43	0.00	0.00
101	-42.43	0.00	0.00
102	-2.00	0.00	0.00
103	22.96	0.00	0.00
104	42.43	0.00	0.00
105	55.00	0.00	0.00
106	55.43	0.00	0.00
107	42.43	0.00	0.00
108	2.00	0.00	0.00
109	-22.96	0.00	0.00
110	-42.43	0.00	0.00
111	-55.00	0.00	0.00
112	-55.43	0.00	0.00
113	-42.43	0.00	0.00
114	-2.00	0.00	0.00
115	22.96	0.00	0.00
116	42.43	0.00	0.00
117	55.00	0.00	0.00
118	55.43	0.00	0.00
119	42.43	0.00	0.00
120	2.00	0.00	0.00
121	-22.96	0.00	0.00
122	-42.43	0.00	0.00



123	-5.43	-2.96	-40.00
124	-50.43	20.00	-40.00
125	-55.43	22.96	-40.00
126	-42.43	42.43	-40.00
127	-22.96	55.43	-40.00
128	0.00	0.00	-60.00
129	42.43	60.00	-60.00
130	60.00	42.43	-60.00
131	42.43	0.00	-60.00
132	0.00	-42.43	-60.00
133	-42.43	-60.00	-60.00
134	-60.00	-42.43	-60.00
135	-42.43	0.00	-60.00
136	0.00	42.43	-60.00
137	42.43	30.00	-80.00
138	30.00	21.21	-80.00
139	21.21	0.00	-80.00
140	0.00	-21.21	-80.00
141	-21.21	-30.00	-80.00
142	-30.00	-21.21	-80.00
143	-21.21	0.00	-80.00
144	0.00	21.21	-80.00
145	21.21	30.00	-80.00
146	30.00	42.43	-80.00
147	42.43	55.43	-80.00
148	55.43	60.00	-80.00
149	60.00	42.43	-80.00
150	42.43	22.96	-80.00
151	22.96	0.00	-80.00
152	0.00	-22.96	-80.00
153	-22.96	-42.43	-80.00
154	-42.43	-55.43	-80.00
155	-55.43	-60.00	-80.00
156	-60.00	-42.43	-80.00
157	-42.43	-22.96	-80.00
158	-22.96	0.00	-80.00
159	0.00	22.96	-80.00
160	22.96	42.43	-80.00
161	42.43	55.43	-95.00
162	55.43	60.00	-95.00
163	60.00	42.43	-95.00
164	42.43	0.00	-95.00
165	0.00	-42.43	-95.00
166	-42.43	-60.00	-95.00
167	-60.00	-42.43	-95.00
168	-42.43	0.00	-95.00
169	0.00	42.43	-95.00
170	42.43	42.43	-95.00





171	0.0	0.0	95.00
172	2.43	60.00	95.00
173	60.00	42.03	95.00
174	42.43	0.0	95.00
175	42.43	-42.43	95.00
176	0.0	-60.00	95.00
177	-42.42	-42.03	95.00
178	-42.43	0.0	95.00
179	0.0	42.43	95.00
180	16.80	75.00	80.00
181	38.10	71.90	80.00
182	53.03	62.90	80.00
183	62.80	53.03	80.00
184	71.30	38.60	80.00
185	75.00	17.00	80.00
186	71.30	0.0	80.00
187	52.80	-17.00	80.00
188	53.03	-38.60	80.00
189	38.10	-53.03	80.00
190	17.0	-62.90	80.00
191	0.0	-71.80	80.00
192	-16.90	-75.00	80.00
193	-38.50	-71.90	80.00
194	-53.03	-62.80	80.00
195	-62.70	-53.03	80.00
196	-71.30	-38.60	80.00
197	-75.00	-17.10	80.00
198	-62.30	0.0	80.00
199	-53.90	17.10	80.00
200	-38.03	38.30	80.00
201	-38.60	53.03	80.00
202	-17.30	62.80	80.00
203	0.80	71.30	40.00
204	16.80	75.00	40.00
205	38.10	71.60	40.00
206	53.03	62.90	40.00
207	62.80	53.03	40.00
208	71.30	38.60	40.00
209	75.00	17.00	40.00
210	71.30	0.0	40.00
211	52.80	-17.00	40.00
212	53.03	-38.60	40.00
213	38.10	-53.03	40.00
214	17.10	-62.90	40.00
215	0.90	-71.80	40.00
216	-16.90	-75.00	40.00
217	-38.50	-71.90	40.00
218		-62.80	40.00



219	53.03	-53.03	40.00
220	-62.70	-38.60	40.00
221	-71.30	-17.10	40.00
222	-75.00	0.00	40.00
223	-71.30	17.10	40.00
224	-62.90	38.30	40.00
225	-53.03	53.03	40.00
226	-38.60	62.80	40.00
227	-17.30	71.30	40.00
228	-17.30	75.00	0.00
229	16.80	71.60	0.00
230	38.10	62.90	0.00
231	53.03	53.03	0.00
232	62.80	38.60	0.00
233	71.30	17.00	0.00
234	75.00	0.00	0.00
235	71.30	17.00	0.00
236	53.03	-17.60	0.00
237	38.10	-53.03	0.00
238	17.10	-62.90	0.00
239	17.10	-71.80	0.00
240	0.00	-71.80	0.00
241	-16.90	-71.80	0.00
242	-38.50	-62.80	0.00
243	-53.03	-53.03	0.00
244	-52.70	-38.60	0.00
245	-71.30	-17.10	0.00
246	-75.00	0.00	0.00
247	-71.30	17.10	0.00
248	-62.90	38.30	0.00
249	-53.03	53.03	0.00
250	-38.60	62.80	0.00
251	-17.30	71.30	0.00
252	16.80	75.00	0.00
253	38.10	71.60	-40.00
254	53.03	62.90	-40.00
255	62.80	53.03	-40.00
256	71.30	38.60	-40.00
257	75.00	17.00	-40.00
258	71.30	0.00	-40.00
259	62.90	17.00	-40.00
260	53.03	-38.60	-40.00
261	38.10	-53.03	-40.00
262	17.10	-62.90	-40.00
263	17.10	-71.80	-40.00
264	0.00	-71.80	-40.00
265	-16.90	-75.00	-40.00
266	-38.50	-62.80	-40.00



267	-53.03	-53.00	-40.00
268	-62.70	-38.60	-40.00
269	-71.30	-17.10	-40.00
270	-75.00	17.10	-40.00
271	-71.30	38.30	-40.00
272	-62.90	53.03	-40.00
273	-53.03	62.80	-40.00
274	-38.60	71.30	-40.00
275	-17.30	75.00	-40.00
276	-10.00	71.60	-80.00
277	16.80	62.90	-80.00
278	38.10	53.03	-80.00
279	53.03	38.60	-80.00
280	62.80	17.00	-80.00
281	71.30	17.00	-80.00
282	75.00	17.00	-80.00
283	71.30	-17.00	-80.00
284	62.80	-38.60	-80.00
285	53.03	-53.03	-80.00
286	38.10	-62.90	-80.00
287	17.00	-71.80	-80.00
288	16.90	-75.00	-80.00
289	-38.50	-71.80	-80.00
290	-53.03	-62.80	-80.00
291	-62.70	-53.03	-80.00
292	-71.30	-38.60	-80.00
293	-75.00	-17.10	-80.00
294	-71.30	17.00	-80.00
295	-62.90	17.10	-80.00
296	-53.03	38.30	-80.00
297	-38.60	53.03	-80.00
298	-17.30	62.80	-80.00
299	-10.00	71.30	-80.00
300	0.00	30.00	-110.00
301	21.21	21.21	-110.00
302	30.00	21.21	-110.00
303	21.21	21.00	-110.00
304	20.00	-21.21	-110.00
305	-21.21	-21.21	-110.00
306	-21.21	-21.00	-110.00
307	-30.00	21.21	-110.00
308	-21.21	21.00	-110.00
309	22.96	60.43	-110.00
310	42.43	55.43	-110.00
311	55.43	42.96	-110.00
312	60.43	22.00	-110.00
313	55.43	0.96	-110.00
314	60.43	-22.96	-110.00













411	34.44	-83.15	-80.00
412	17.56	-88.27	-80.00
413	17.00	-90.07	-80.00
414	-17.56	-88.15	-80.00
415	-34.44	-74.83	-80.00
416	-50.00	-63.64	-80.00
417	-63.64	-50.00	-80.00
418	-74.83	-34.44	-80.00
419	-83.15	-17.56	-80.00
420	-88.27	17.00	-80.00
421	-90.00	17.56	-80.00
422	-88.27	34.44	-80.00
423	-83.15	50.00	-80.00
424	-74.83	63.64	-80.00
425	-63.64	74.83	-80.00
426	-50.00	83.15	-80.00
427	-34.44	88.27	-80.00
428	-17.56	90.00	-60.00
429	-17.00	83.15	-60.00
430	34.44	63.64	-60.00
431	63.64	34.44	-60.00
432	83.15	0.00	-60.00
433	90.00	-34.44	-60.00
434	83.15	-63.64	-60.00
435	63.64	-90.00	-60.00
436	34.44	-83.15	-60.00
437	0.00	-63.64	-60.00
438	-34.44	-34.00	-60.00
439	-63.64	34.44	-60.00
440	-83.15	63.64	-60.00
441	-90.00	83.15	-60.00
442	-83.15	90.00	-40.00
443	-63.64	88.27	-40.00
444	-34.44	83.15	-40.00
445	0.00	74.83	-40.00
446	17.56	50.00	-40.00
447	34.44	34.44	-40.00
448	50.00	17.56	-40.00
449	63.64	17.00	-40.00
450	74.83	-17.56	-40.00
451	83.15	-34.44	-40.00
452	88.27	-50.00	-40.00
453	90.00	-63.64	-40.00
454	88.27	-74.83	-40.00
455	83.15	-83.15	-40.00
456	74.83	-90.00	-40.00
457	63.64	-83.15	-40.00
458	50.00	-74.83	-40.00



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[illegible]



507	34.44	-83.15	0.0
508	17.56	-88.27	0.0
509	17.56	-90.27	0.0
510	-17.56	-88.27	0.0
511	-34.44	-83.15	0.0
512	-50.00	-74.83	0.0
513	-63.64	-63.64	0.0
514	-74.83	-50.00	0.0
515	-83.15	-34.44	0.0
516	-88.27	-17.56	0.0
517	-90.00	17.56	0.0
518	-88.27	34.44	0.0
519	-74.83	50.00	0.0
520	-63.64	63.64	0.0
521	-50.00	74.83	0.0
522	-34.44	83.15	0.0
523	-17.56	88.27	0.0
524	17.56	90.00	0.0
525	34.44	83.15	20.00
526	63.64	63.64	20.00
527	83.15	34.44	20.00
528	90.00	0.0	20.00
529	83.15	-34.44	20.00
530	63.64	-63.64	20.00
531	34.44	-83.15	20.00
532	0.0	-90.00	20.00
533	-34.44	-83.15	20.00
534	-63.64	-63.64	20.00
535	-83.15	-34.44	20.00
536	-90.00	17.56	20.00
537	-88.27	34.44	20.00
538	-74.83	50.00	20.00
539	-63.64	63.64	20.00
540	-50.00	74.83	40.00
541	-34.44	83.15	40.00
542	17.56	88.27	40.00
543	34.44	90.00	40.00
544	50.00	83.15	40.00
545	63.64	74.83	40.00
546	74.83	63.64	40.00
547	83.15	50.00	40.00
548	88.27	34.44	40.00
549	90.00	17.56	40.00
550	88.27	-17.56	40.00
551	74.83	-34.44	40.00
552	63.64	-50.00	40.00
553	50.00	-63.64	40.00
554	34.44	-74.83	40.00









603	34.44	-83.15	80.00
604	17.56	-88.27	80.00
605	17.00	-90.00	80.00
606	-17.56	-88.27	80.00
607	-34.44	-83.15	80.00
608	-50.00	-74.83	80.00
609	-63.64	-63.64	80.00
610	-74.83	-50.00	80.00
611	-88.27	-34.44	80.00
612	-90.00	-17.56	80.00
613	-88.27	17.56	80.00
614	-83.15	34.44	80.00
615	-74.83	50.00	80.00
616	-63.64	63.64	80.00
617	-50.00	74.83	80.00
618	-34.44	88.27	80.00
619	-17.56	90.00	80.00
620	17.56	83.15	95.00
621	34.44	63.64	95.00
622	63.64	34.44	95.00
623	83.15	0.00	95.00
624	90.00	-34.44	95.00
625	83.15	-63.64	95.00
626	63.64	-83.15	95.00
627	34.44	-90.00	95.00
628	0.00	-83.15	95.00
629	-34.44	-63.64	95.00
630	-63.64	-34.44	95.00
631	-83.15	34.44	95.00
632	-90.00	63.64	95.00
633	-83.15	83.15	95.00
634	-63.64	90.00	95.00
635	-34.44	88.27	95.00
636	17.56	83.15	95.00
637	34.44	74.83	110.00
638	50.00	63.64	110.00
639	63.64	50.00	110.00
640	74.83	34.44	110.00
641	83.15	17.56	110.00
642	90.00	0.00	110.00
643	88.27	-17.56	110.00
644	83.15	-34.44	110.00
645	74.83	-50.00	110.00
646	63.64	-63.64	110.00
647	50.00	-74.83	110.00
648	34.44	-83.15	110.00
649	17.56	-90.00	110.00
650	0.00	-83.15	110.00







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42.43  
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-50.00  
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-22.96  
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KK=CONN(I,K)
IF(KK.GT.183) GO TO 80
JA(JJ)=JA(JJ)+I
JAA=JA(JJ)
DO 70 L=2,JAA
JJL=MAME(JJ,L)
IF(JJL.EQ.KK) JA(JJ)=JA(JJ)-1
IF(JJL.EQ.KK) GO TO 80
IF(JJL.EQ.0) MAME(JJ,JAA)=KK
CONTINUE
70 CONTINUE
80 CONTINUE
90 CONTINUE
100 CONTINUE
C
JB(1)=1
JC=0
DO 200 I=1,183
JN=JA(I)
JB(I+1)=JB(I)+JA(I)
JC=JC+JA(I)
CONTINUE
200 WRITE(6,215) JC
DO 250 I=1,183
JAA=JA(I)
JBL=JB(I)
DO 240 J=1,JAA
JJ=JBL+J-1
NAME(JJ)=MAME(I,J)
CONTINUE
240 CONTINUE
250 WRITE(6,204)
WRITE(6,205) (JA(I),I=1,183)
WRITE(6,207)
WRITE(6,205) (JB(I),I=1,183)
WRITE(6,208)
WRITE(6,205) (NAME(I),I=1,JC)
STOP
END

```



```

C-----
C  THIS IS THE MAIN PROGRAM WHICH CALCULATES THE BIGG AND BIGGG
C  SYSTEM MATRICES BY USING SUBROUTINES FLOWIE AND TANYA. THE SYSTEM
C  MATRICES ARE THEN PUT ON TAPE NUMBER NPS182.
C-----
      REAL*8 G,GG,BIGG,BIGGG
      REAL*8 WL,WS,S,CL1,CL2,CL3,FN,DL1,DL2,DL3,DS,DETJ
      INTEGER*2 NAME,JA,JB,NNZ,JC,NUMNP,NUMEL,NELDOF,CONN
      COMMON/GTRY2/NAME(4615),JA(183),JB(183),NNZ,JC,NUMNP,NUMEL,NELDOF,
1     CONN(128,15)
      COMMON/GTRY1/X(505),P(505),Z(505)
      COMMON/XOCAL/XX(15),YY(15),ZZ(15),PPSI(15)
      COMMON/GMAT/G(15,15)
      COMMON/GGMAT/GG(15,15)
      COMMON/SYSMT1/BIGG(4615)
      COMMON/SYSMT2/BIGGG(4615)
      COMMON/GAJSS/WL(7),WS(5),S(5),CL1(7),CL2(7),CL3(7),FN(15,21),DL1(1
15,21),DL2(15,21),DL3(15,21),DS(15,21),DETJ(21)
      FORMAT(16I5)
150  FORMAT(25I5)
160  FORMAT(8E16.8)
211  FORMAT(12F10.4)
609  FORMAT(20I6)
702  FORMAT(10F8.2)
800  FORMAT(20I4)
952  NNZ=183
      NUMEL=128
      NUMNP=505
      JC=4615
      NELDOF=15
455  DO 455 I=1,NUMEL
      READ(5,150)I,(CONN(I,J),J=1,NELDOF)
      CONTINUE
      READ(5,952)(JA(I),I=1,NNZ)
      READ(5,952)(JB(I),I=1,NNZ)
      READ(5,952)(NAME(I),I=1,JC)
      READ(5,800)(X(I),I=1,NUMNP)
      READ(5,800)(P(I),I=1,NUMNP)
      READ(5,800)(Z(I),I=1,NUMNP)
      DO 465 I=1,JC
      BIGG(I)=0.0
      BIGGG(I)=0.0
465  CONTINUE
      CALL TANYA
      CALL FLOWIE
      STOP

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M-010  
M-015

M-060

M-100

M-120  
M-125  
M-130

M-160  
M-140  
M-145

M-355  
M-360  
M-365  
M-370  
M-405  
M-415  
M-615



END

M-620

THIS SUBROUTINE EVALUATES THE GG ELEMENT MATRIX USING 20 POINTS OF  
INTEGRATION. THE GG ELEMENT MATRIX IS THEN INSERTED INTO THE BIGGG  
SYSTEM MATRIX.

SUBROUTINE TANYA

REAL\*8 GG,BIGGG

REAL\*8 WL,WS,S,CL1,CL2,CL3,FN,DL1,DL2,DL3,DS,DETJ

INTEGER\*2 NAME,JA,JB,NNZ,JC,NJMNP,NUMEL,NELDOF,CONN

COMMON/GTRY2/NAME(8199),JA(299),JB(299),NNZ,JC,NUMNP,NUMEL,NELDOF,

1CONN(192,15)

COMMON/GTRY1/X(717),P(717),Z(717)

COMMON/XOCAL/KX(15),YY(15),ZZ(15),PPSI(15)

COMMON/GGMAT/GG(15,15)

COMMON/SYSMT2/BIGGG(8199)

COMMON/GAUSS/WL(7),WS(5),S(5),CL1(7),CL2(7),CL3(7),FN(15,21),DL1(1

15,21),DL2(15,21),DL3(15,21),DS(15,21),DETJ(21)

WS(1)=0.236926885056189

WS(2)=0.478628670499366

WS(3)=0.56888888888889

WS(4)=0.478628670499366

WS(5)=0.236926885056189

S(1)=0.906179845938664

S(2)=0.538469310105683

S(3)=0.0

S(4)=0.538469310105683

S(5)=0.906179845938664

WL(1)=0.281250

WL(2)=0.260416666666667

WL(3)=0.260416666666667

WL(4)=0.260416666666667

CL1(1)=0.333333333333333

CL1(2)=0.60

CL1(3)=0.20

CL1(4)=0.20

CL2(1)=0.333333333333333

CL2(2)=0.20

CL2(3)=0.60

CL2(4)=0.20

CL3(1)=0.333333333333333

CL3(2)=0.20

CL3(3)=0.60

CL3(4)=0.20

DO 80 K=1,5

DO 90 M=1,4

T-015





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IF(K.EQ.1) GO TO 10
IF(K.EQ.2) GO TO 20
IF(K.EQ.3) GO TO 30
IF(K.EQ.4) GO TO 40
IF(K.EQ.5) GO TO 50
10 N=M
GO TO 60
20 N=M+4
GO TO 60
30 N=M+8
GO TO 60
40 N=M+12
GO TO 60
50 N=M+16
GO TO 60
60 CONTINUE
FN(1,V)=0.50*CL1(M)*((2.0*CL1(M)-1.0)*(1.0+S(K))-(1.0-S(K)**2))
FN(2,V)=0.50*CL1(M)*CL2(M)*(1.0+S(K))
FN(3,V)=0.50*CL2(M)*((2.0*CL2(M)-1.0)*(1.0+S(K))-(1.0-S(K)**2))
FN(4,V)=0.50*CL2(M)*CL3(M)*(1.0+S(K))
FN(5,V)=0.50*CL3(M)*((2.0*CL3(M)-1.0)*(1.0+S(K))-(1.0-S(K)**2))
FN(6,V)=0.50*CL3(M)*CL3(M)*(1.0+S(K))
FN(7,V)=CL1(M)*((1.0-S(K)**2))
FN(8,V)=CL2(M)*((1.0-S(K)**2))
FN(9,V)=CL3(M)*((1.0-S(K)**2))
FN(10,V)=0.50*CL1(M)*((2.0*CL1(M)-1.0)*(1.0-S(K))-(1.0-S(K)**2))
FN(11,V)=0.50*CL1(M)*CL2(M)*(1.0-S(K))
FN(12,V)=0.50*CL2(M)*((2.0*CL2(M)-1.0)*(1.0-S(K))-(1.0-S(K)**2))
FN(13,V)=0.50*CL2(M)*CL3(M)*(1.0-S(K))
FN(14,V)=0.50*CL3(M)*((2.0*CL3(M)-1.0)*(1.0-S(K))-(1.0-S(K)**2))
FN(15,V)=0.50*CL3(M)*CL3(M)*(1.0-S(K))
DL1(1,V)=0.50*CL1(M)*CL1(M)-1.0
DL1(2,V)=2.0*CL2(M)*(1.0+S(K))
DL1(3,V)=0.0
DL1(4,V)=0.0
DL1(5,V)=0.0
DL1(6,V)=2.0*CL3(M)*(1.0+S(K))
DL1(7,V)=0.0
DL1(8,V)=0.0
DL1(9,V)=0.0
DL1(10,V)=0.50*CL1(M)*((4.0*CL1(M)-1.0)-(1.0-S(K)**2))
DL1(11,V)=2.0*CL2(M)*(1.0-S(K))
DL1(12,V)=0.0
DL1(13,V)=0.0
DL1(14,V)=0.0
DL1(15,V)=2.0*CL3(M)*(1.0-S(K))
DL2(1,V)=0.0
DL2(2,V)=2.0*CL1(M)*(1.0+S(K))
DL2(3,V)=0.50*CL1(M)*((4.0*CL2(M)-1.0)-(1.0-S(K)**2))

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DL2(4,N)=2.0*CL3(M)*(1.0+S(K))
DL2(5,N)=0.0
DL2(6,N)=0.0
DL2(7,N)=0.0
DL2(8,N)=(1.0-S(K)**2)
DL2(9,N)=0.0
DL2(10,N)=0.0
DL2(11,N)=2.0*CL1(M)*(1.0-S(K))
DL2(12,N)=0.50*((1.0-S(K))*(4.0*CL2(M)-1.0)-(1.0-S(K)**2))
DL2(13,N)=2.0*CL3(M)*(1.0-S(K))
DL2(14,N)=0.0
DL2(15,N)=0.0
DL3(1,N)=0.0
DL3(2,N)=0.0
DL3(3,N)=0.0
DL3(4,N)=2.0*CL2(M)*(1.0+S(K))
DL3(5,N)=0.50*((1.0+S(K))*(4.0*CL3(M)-1.0)-(1.0-S(K)**2))
DL3(6,N)=2.0*CL1(M)*(1.0+S(K))
DL3(7,N)=0.0
DL3(8,N)=0.0
DL3(9,N)=(1.0-S(K)**2)
DL3(10,N)=0.0
DL3(11,N)=0.0
DL3(12,N)=0.0
DL3(13,N)=2.0*CL2(M)*(1.0-S(K))
DL3(14,N)=0.50*((1.0-S(K))*(4.0*CL3(M)-1.0)-(1.0-S(K)**2))
DL3(15,N)=2.0*CL1(M)*(1.0-S(K))
DS(1,V)=0.50*(2.0*CL1(M)**2-CL1(M))+CL1(M)*S(K)
DS(2,N)=2.0*CL1(M)*CL2(M)
DS(3,N)=0.50*(2.0*CL2(M)**2-CL2(M))+CL2(M)*S(K)
DS(4,V)=2.0*CL2(M)*CL3(M)
DS(5,N)=0.50*(2.0*CL3(M)**2-CL3(M))+CL3(M)*S(K)
DS(6,V)=2.0*CL1(M)*CL3(M)
DS(7,N)=-2.0*CL1(M)*S(K)
DS(8,V)=-2.0*CL2(M)*S(K)
DS(9,N)=-2.0*CL3(M)*S(K)
DS(10,N)=-0.50*(2.0*CL1(M)**2-CL1(M))+CL1(M)*S(K)
DS(11,N)=-0.50*CL1(M)*CL2(M)
DS(12,N)=-0.50*(2.0*CL2(M)**2-CL2(M))+CL2(M)*S(K)
DS(13,V)=-0.50*CL2(M)*CL3(M)
DS(14,N)=-0.50*(2.0*CL3(M)**2-CL3(M))+CL3(M)*S(K)
DS(15,V)=-2.0*CL1(M)*CL3(M)
90 CONTINUE
80
DO 510 L=1,NUMEL
DO 475 I=1,NELDOF
DO 470 J=1,NELDOF
GG(I,J)=0.0

```



```

470 CONTINUE
475 CONTINUE
DO 480 JJ=1,NELDOF
NN=CONN(L,JJ)
XX(JJ)=X(NN)
YY(JJ)=P(NN)
ZZ(JJ)=Z(NN)
480 CONTINUE
A5=2.0/(ZZ(1)-ZZ(15))
DETJJ=0.0
DO 200 K=1,5
DETJ1=0.0
DO 210 M=1,4
IF(K.EQ.1) GO TO 11
IF(K.EQ.2) GO TO 21
IF(K.EQ.3) GO TO 31
IF(K.EQ.4) GO TO 41
IF(K.EQ.5) GO TO 51
11 N=M
GO TO 61
21 N=M+4
GO TO 61
31 N=M+8
GO TO 61
41 N=M+12
GO TO 61
51 N=M+16
GO TO 61
61 CONTINUE
T1=DL1(1,N)*XX(1)+DL1(2,N)*XX(2)-DL3(4,N)*XX(4)-DL3(5,N)*XX(5)
T2=DL1(7,N)*XX(7)-DL3(9,N)*XX(9)+DL1(10,N)*XX(10)+DL1(11,N)*XX(11)
T3=(DL1(6,N))*XX(6)-DL3(13,N)*XX(13)-DL3(14,N)*XX(14)
T4=(DL1(15,N))-DL3(15,N)*XX(15)
DJ11=T1+T2+T3+T4
D1=DL1(1,N)*YY(1)+DL1(2,N)*YY(2)-DL3(4,N)*YY(4)-DL3(5,N)*YY(5)
D2=DL1(7,N)*YY(7)-DL3(9,N)*YY(9)+DL1(10,N)*YY(10)+DL1(11,N)*YY(11)
D3=(DL1(6,N))*YY(6)-DL3(13,N)*YY(13)-DL3(14,N)*YY(14)
D4=(DL1(15,N))-DL3(15,N)*YY(15)
DJ12=D1+D2+D3+D4
T5=DL2(2,N)*XX(2)+DL2(3,N)*XX(3)-DL3(5,N)*XX(5)-DL3(6,N)*XX(6)
T6=DL2(8,N)*XX(8)-DL3(9,N)*XX(9)+DL2(4,N)*XX(4)
T7=DL2(11,N)*XX(11)+DL2(12,N)*XX(12)+(DL2(13,N)-DL3(13,N))*XX(13)
T8=-DL3(14,N)*XX(14)-DL3(15,N)*XX(15)
DJ21=T5+T6+T7+T8
D5=DL2(2,N)*YY(2)+DL2(3,N)*YY(3)-DL3(5,N)*YY(5)-DL3(6,N)*YY(6)
D6=DL2(8,N)*YY(8)-DL3(9,N)*YY(9)+DL2(4,N)*YY(4)
D7=DL2(11,N)*YY(11)+DL2(12,N)*YY(12)+(DL2(13,N)-DL3(13,N))*YY(13)
D8=-DL3(14,N)*YY(14)-DL3(15,N)*YY(15)
DJ22=D5+D6+D7+D8

```



```

DETJ(N)=DJ11*DJ22-DJ12*DJ21
DETJ1=DETJ1+D=J(N)*WL(M)
CONTINUE
210 DETJJ=DETJJ+WS(K)*DETJ1
CONTINUE
200 DO 600 J=1,15
DO 610 I=1,15
FEG=0.0
DO 660 K=1,5
FEG=0.0
DO 670 M=1,4
IF(K.EQ.1) GO TO 12
IF(K.EQ.2) GO TO 22
IF(K.EQ.3) GO TO 32
IF(K.EQ.4) GO TO 42
IF(K.EQ.5) GO TO 52
12 N=M
GO TO 62
22 N=M+4
GO TO 62
32 N=M+8
GO TO 62
42 N=M+12
GO TO 62
52 N=M+16
GO TO 62
62 CONTINUE
DXDL5=DL1(1,N)*XX(1)+(DL1(2,N)-DL2(2,N))*XX(2)-DL2(3,N)*XX(3)
DXDL5=-{(DL2(4,N)+DL3(4,N))*XX(4)-DL3(5,N)*XX(5)}
DXDL7=(DL1(6,N)-DL3(6,N))*XX(6)+DL1(7,N)*XX(7)-DL2(8,N)*XX(8)
DXDL9=-DL3(9,N)*XX(9)+DL1(10,N)*XX(10)+(DL1(11,N)-DL2(11,N))*XX(11)
1)
DXDL9=-DL2(12,N)*XX(12)-(DL2(13,N)+DL3(13,N))*XX(13)-DL3(14,N)*XX(14)
1)
DXDL1=(DL1(15,N)-DL3(15,N))*XX(15)
DXDL1=DXDL5+DXDL6+DXDL7+DXDL8+DXDL9
DYDL5=DL1(1,N)*YY(1)+(DL1(2,N)-DL2(2,N))*YY(2)-DL2(3,N)*YY(3)
DYDL6=-{(DL2(4,N)+DL3(4,N))*YY(4)-DL3(5,N)*YY(5)}
DYDL7=(DL1(6,N)-DL3(6,N))*YY(6)+DL1(7,N)*YY(7)-DL2(8,N)*YY(8)
DYDL8=-DL3(9,N)*YY(9)+DL1(10,N)*YY(10)+(DL1(11,N)-DL2(11,N))*YY(11)
1)
DYDL9=-DL2(12,N)*YY(12)-(DL2(13,N)+DL3(13,N))*YY(13)-DL3(14,N)*YY(14)
1)
DYDL1=(DL1(15,N)-DL3(15,N))*YY(15)
DYDL1=DYDL5+DYDL6+DYDL7+DYDL8+DYDL9
DXDL55=-DL1(1,N)*XX(1)+(DL2(2,N)-DL1(2,N))*XX(2)+DL2(3,N)*XX(3)
DXDL56=(DL2(4,N)-DL3(4,N))*XX(4)-DL3(5,N)*XX(5)-DL1(7,N)*XX(7)
DXDL77=-{(DL1(5,N)+DL3(6,N))*XX(6)+DL2(8,N)*XX(8)-DL3(9,N)*XX(9)}
DXDL88=-DL1(10,N)*XX(10)+(DL2(11,N)-DL1(11,N))*XX(11)
DXDL99=DL2(12,N)*XX(12)+(DL2(13,N)-DL3(13,N))*XX(13)-DL3(14,N)*XX(14)
1)
DXDL99=-DL1(15,N)+DL3(15,N))*XX(15)

```





```

DXDL2=DXDL55+DXDL66+DXDL77+DXDL88+DXDL99
DYDL55=-DL1(1,N)*YY(1)+(DL2(2,N)-DL1(2,N))*YY(2)+DL2(3,N)*YY(3)
DYDL66=(DL2(4,N)-DL3(4,N))*YY(4)-DL3(5,N)*YY(5)-DL1(7,N)*YY(7)
DYDL77=-DL1(6,N)+DL3(6,N))*YY(6)+DL2(8,N)*YY(8)-DL3(9,N)*YY(9)
DYDL88=-DL1(10,N)*YY(10)+(DL2(11,N)-DL1(11,N))*YY(11)
DYDL99=DL2(12,N)*YY(12)+(DL2(13,N)-DL3(13,N))*YY(13)-DL3(14,N)*YY(
114)-(DL1(15,N)+DL3(15,N))*YY(15)
DYDL2=DYDL55+DYDL66+DYDL77+DYDL88+DYDL99
DL1DY=1.0/DYDL1
DL2DX=1.0/DXDL2
DL1DX=1.0/DXDL1
DL2DY=1.0/DYDL2
BB11=(DL1DX**2)+(DL1DY**2)
BB12=DL1DX*DL2DX+DL1DY*DL2DY
BB22=(DL2DX**2)+(DL2DY**2)
H1=BB11*((DL1(J,N)-DL3(J,N))*DL1(1,N)-DL3(1,N))
H2=BB12*((DL2(J,N)-DL3(J,N))*DL1(1,N)-DL3(1,N))
H3=BB12*((DL1(J,N)-DL3(J,N))*DL2(1,N)-DL3(1,N))
H4=BB22*((DL2(J,N)-DL3(J,N))*DL2(1,N)-DL3(1,N))
H5=(A5**2)*DS(I,N)*DS(I,N)
H6=(H1+H2+H3+H4+H5)*DETJ(N)
FFG=FFG+WL(M)*H6
670 CONTINUE
FFG=WS(K)*FFG+FFG
660 CONTINUE
GG(J,I)=FFG/A5
610 CONTINUE
600 CONTINUE
*** INSERT ELEMENT MATRIX INTO SYSTEM MATRIX *****
DO 505 K=1,NELDOF
KK=CONN(L,K)
IF(KK.GT.NNZ) GO TO 505
KKK=JA(KK)
LLL=JB(KK)-1
DO 500 I=1,NELDOF
II=CONN(L,I)
IF(II.GT.NNZ) GO TO 500
DO 490 M=1,KKK
MM=LLL+M
KKM=NAME(MM)
IF(II.EQ.KKM) GO TO 495
490 CONTINUE
495 CONTINUE
BIGGG(MM)=BIGGG(MM)+GG(K,I)
500 CONTINUE
505 CONTINUE
510 RETURN

```



END

THIS SUBROUTINE EVALUATES THE G ELEMENT MATRIX USING 21 POINTS OF  
INTEGRATION. THE G ELEMENT MATRIX IS THEN INSERTED INTO THE BIGG  
SYSTEM MATRIX.

SUBROUTINE FLOWIE

```

REAL*8 G,BIGG
REAL*8 WL,WS,S,CL1,CL2,CL3,FN,DL1,DL2,DL3,DS,DETJ
INTEGER*2 NAME,JA,JB,NNZ,JC,NJMNP,NUMEL,NELDOF,CONN
COMMON/GTRY2/NAME(8199),JA(299),JB(299),NNZ,JC,NUMNP,NUMEL,NELDOF,
1 CONN(192,15)
COMMON/GIRYL/X(717),P(717),Z(717)
COMMON/XOCAL/XX(15),YY(15),ZZ(15),PPSI(15)
COMMON/GMAT/G(15,15)
COMMON/SYSTL/BIGG(8199)
COMMON/GAUSS/WL(7),WS(5),S(5),CL1(7),CL2(7),CL3(7),FN(15,21),DL1(1
1 5,21),DL2(15,21),DL3(15,21),DS(15,21),DETJ(21)
WS(1)=0.5555555555555556
WS(2)=0.8888888888888889
WS(3)=0.5555555555555556
S(1)=0.774596669241483
S(2)=0.0
S(3)=-.774596669241483
WL(1)=0.112500
WL(2)=0.06619707500
WL(3)=0.06619707500
WL(4)=0.06619707500
WL(5)=0.0629695900
WL(6)=0.0629695900
WL(7)=0.0629695900
CL1(1)=0.333333333333333
CL1(2)=0.059615870
CL1(3)=0.470142060
CL1(4)=0.470142060
CL1(5)=0.797426990
CL1(6)=0.101286510
CL1(7)=0.101286510
CL2(1)=0.333333333333333
CL2(2)=0.470142060
CL2(3)=0.059615870
CL2(4)=0.470142060
CL2(5)=0.101286510
CL2(6)=0.797426990
CL2(7)=0.101286510
CL3(1)=0.333333333333333

```



```

CL3(2)=0.470142060
CL3(3)=0.470142060
CL3(4)=0.059615870
CL3(5)=0.101286510
CL3(6)=0.101286510
CL3(7)=0.797426990
DO 10 K=1,3
DO 20 M=1,7
C THE 3 X 7 MATRICES ARE BEING CONVERTED TO 1 X 21 VECTORS
30 N=M
IF(K-2)30,40,50
40 N=M+7
50 N=M+14
60 CONTINUE
FN(1,N)=0.50*CL1(M)*((2.0*CL1(M)-1.0)*(1.0+S(K))-(1.0-S(K)**2))
FN(2,N)=2.0*CL1(M)*CL2(M)*(1.0+S(K))
FN(3,N)=0.50*CL2(M)*((2.0*CL2(M)-1.0)*(1.0+S(K))-(1.0-S(K)**2))
FN(4,N)=2.0*CL2(M)*CL3(M)*(1.0+S(K))
FN(5,N)=0.50*CL3(M)*((2.0*CL3(M)-1.0)*(1.0+S(K))-(1.0-S(K)**2))
FN(6,N)=2.0*CL1(M)*CL3(M)*(1.0+S(K))
FN(7,N)=CL1(M)*(1.0-S(K)**2)
FN(8,N)=CL2(M)*(1.0-S(K)**2)
FN(9,N)=CL3(M)*((2.0*CL1(M)-1.0)*(1.0-S(K))-(1.0-S(K)**2))
FN(10,N)=0.50*CL1(M)*CL2(M)*(1.0-S(K))
FN(11,N)=2.0*CL1(M)*CL2(M)*(1.0-S(K))
FN(12,N)=0.50*CL2(M)*((2.0*CL2(M)-1.0)*(1.0-S(K))-(1.0-S(K)**2))
FN(13,N)=2.0*CL2(M)*CL3(M)*(1.0-S(K))
FN(14,N)=0.50*CL3(M)*((2.0*CL3(M)-1.0)*(1.0-S(K))-(1.0-S(K)**2))
FN(15,N)=2.0*CL1(M)*CL3(M)*(1.0-S(K))
DL1(1,N)=0.50*((1.0+S(K))*(4.0*CL1(M)-1.0)-(1.0-S(K)**2))
DL1(2,N)=2.0*CL2(M)*(1.0+S(K))
DL1(3,N)=0.0
DL1(4,N)=0.0
DL1(5,N)=0.0
DL1(6,N)=2.0*CL3(M)*(1.0+S(K))
DL1(7,N)=(1.0-S(K)**2)
DL1(8,N)=0.0
DL1(9,N)=0.0
DL1(10,N)=0.50*((1.0-S(K))*(4.0*CL1(M)-1.0)-(1.0-S(K)**2))
DL1(11,N)=2.0*CL2(M)*(1.0-S(K))
DL1(12,N)=0.0
DL1(13,N)=0.0
DL1(14,N)=0.0
DL1(15,N)=2.0*CL3(M)*(1.0-S(K))
DL2(1,N)=0.0
DL2(2,N)=2.0*CL1(M)*(1.0+S(K))

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```

DL2(3,N)= 0.50*((1.0+S(K))*(4.0*CL2(M)-1.0)-(1.0-S(K)**2))
DL2(4,N)= 2.0*CL3(M)*(1.0+S(K))
DL2(5,N)= 0.0
DL2(6,N)= 0.0
DL2(7,N)= 0.0
DL2(8,N)= (1.0-S(K)**2)
DL2(9,N)= 0.0
DL2(10,N)= 0.0
DL2(11,N)= 2.0*CL1(M)*(1.0-S(K))
DL2(12,N)= 0.50*((1.0-S(K))*(4.0*CL2(M)-1.0)-(1.0-S(K)**2))
DL2(13,N)= 2.0*CL3(M)*(1.0-S(K))
DL2(14,N)= 0.0
DL2(15,N)= 0.0
DL3(1,N)= 0.0
DL3(2,N)= 0.0
DL3(3,N)= 0.0
DL3(4,N)= 2.0*CL2(M)*(1.0+S(K))
DL3(5,N)= 0.50*((1.0+S(K))*(4.0*CL3(M)-1.0)-(1.0-S(K)**2))
DL3(6,N)= 2.0*CL1(M)*(1.0+S(K))
DL3(7,N)= 0.0
DL3(8,N)= 0.0
DL3(9,N)= (1.0-S(K)**2)
DL3(10,N)= 0.0
DL3(11,N)= 0.0
DL3(12,N)= 0.0
DL3(13,N)= 2.0*CL2(M)*(1.0-S(K))
DL3(14,N)= 0.50*((1.0-S(K))*(4.0*CL3(M)-1.0)-(1.0-S(K)**2))
DL3(15,N)= 2.0*CL1(M)*(1.0-S(K))
DS(1,N)= 0.50*(2.0*CL1(M)**2-CL1(M))+CL1(M)*S(K)
DS(2,N)= 2.0*CL1(M)*CL2(M)
DS(3,N)= 0.50*(2.0*CL2(M)**2-CL2(M))+CL2(M)*S(K)
DS(4,N)= 2.0*(2.0*CL2(M)*CL3(M)
DS(5,N)= 0.50*(2.0*CL3(M)**2-CL3(M))+CL3(M)*S(K)
DS(6,N)= 2.0*CL1(M)*CL3(M)
DS(7,N)= -2.0*CL1(M)*S(K)
DS(8,N)= -2.0*CL2(M)*S(K)
DS(9,N)= -2.0*CL3(M)*S(K)
DS(10,N)= -50*(2.0*CL1(M)**2-CL1(M))+CL1(M)*S(K)
DS(11,N)= -2.0*CL1(M)*CL2(M)
DS(12,N)= -50*(2.0*CL2(M)**2-CL2(M))+CL2(M)*S(K)
DS(13,N)= -2.0*CL2(M)*CL3(M)
DS(14,N)= -50*(2.0*CL3(M)**2-CL3(M))+CL3(M)*S(K)
DS(15,N)= -2.0*CL1(M)*CL3(M)

```

```

20 CONTINUE
10

```

```

DO 510 L=1,NUMEL
DO 475 J=1,NELDOF
DO 470 J=1,NELDOF

```





```

470 G(I,J)=0.0
475 CONTINUE
DO 480 JJ=1,NELDOF
NN=CONN(L,JJ)
XX(JJ)=X(NN)
YY(JJ)=P(NN)
ZZ(JJ)=Z(NN)
480 CONTINUE
A5=2.0/(ZZ(1)-ZZ(15))
DETJJ=0.0
DO 200 K=1,3
DETJ1=0.0
DO 210 M=1,7
IF(K-2)31,41,51
31 N=M
GO TO 61
41 N=M+7
GO TO 61
51 N=M+14
61 CONTINUE
T1=DL1(1,N)*XX(1)+DL1(2,N)*XX(2)-DL3(4,N)*XX(4)-DL3(5,N)*XX(5)
T2=DL1(7,N)*XX(7)-DL3(9,N)*XX(9)+DL1(13,N)*XX(10)+DL1(11,N)*XX(11)
T3=(DL1(6,N)-DL3(6,N))*XX(6)-DL3(13,N)*XX(13)-DL3(14,N)*XX(14)
T4=(DL1(15,N)-DL3(15,N))*XX(15)
DJ1=T1+T2+T3+T4
D1=DL1(1,N)*YY(1)+DL1(2,N)*YY(2)-DL3(4,N)*YY(4)-DL3(5,N)*YY(5)
D2=DL1(7,N)*YY(7)-DL3(9,N)*YY(9)+DL1(10,N)*YY(10)+DL1(11,N)*YY(11)
D3=(DL1(6,N)-DL3(6,N))*YY(6)-DL3(13,N)*YY(13)-DL3(14,N)*YY(14)
D4=(DL1(15,N)-DL3(15,N))*YY(15)
DJ2=D1+D2+D3+D4
T5=DL2(2,N)*XX(2)+DL2(3,N)*XX(3)-DL3(5,N)*XX(5)-DL3(6,N)*XX(6)
T6=DL2(8,N)*XX(8)-DL3(9,N)*XX(9)+(DL2(4,N)-DL3(4,N))*XX(4)
T7=DL2(11,N)*XX(11)+DL2(12,N)*XX(12)+(DL2(13,N)-DL3(13,N))*XX(13)
T8=-DL3(14,N)*XX(14)-DL3(15,N)*XX(15)
DJ21=T5+T6+T7+T8
D5=DL2(2,N)*YY(2)+DL2(3,N)*YY(3)-DL3(5,N)*YY(5)-DL3(6,N)*YY(6)
D6=DL2(8,N)*YY(8)-DL3(9,N)*YY(9)+(DL2(4,N)-DL3(4,N))*YY(4)
D7=DL2(11,N)*YY(11)+DL2(12,N)*YY(12)+(DL2(13,N)-DL3(13,N))*YY(13)
D8=-DL3(14,N)*YY(14)-DL3(15,N)*YY(15)
DJ22=D5+D6+D7+D8
DETJ(N)=DJ1+DJ22-DJ12+DJ21
DETJ1=DETJ(N)*WL(M)+DETJ1
CONTINUE
210 DETJJ=DETJJ+WS(K)*DETJ1
200 CONTINUE
DO 600 J=1,15
DO 610 I=1,15

```



```

FG=0.0
DO 660 K=1,3
F=0.0
DO 670 M=1,7
IF(K-2)620,630,640
620 N=M
GO TO 650
630 N=M+7
GO TO 650
640 N=M+14
650 CONTINUE
F=F+WL(M)*FN(J,N)*FN(I,N)*DETJ(N)
670 CONTINUE
FG=FG+WS(K)*F
660 CONTINUE
G(J,I)=FG/A5
610 CONTINUE
600 ** INSERT ELEMENT MATRIX INTO SYSTEM MATRIX *****
DO 505 K=1,NELDOF
KK=CONN(L,K)
IF(KK.GT.NNZ) GO TO 505
KKK=JA(KK)
LLL=JB(KK)-1
DO 500 I=1,NELDOF
II=CONN(L,I)
IF(II.GT.NNZ) GO TO 500
DO 490 M=1,KK<
MM=LLL+M
KKM=NAME(MM)
IF(II.EQ.KKM) GO TO 495
490 CONTINUE
495 BIGG(MM)=BIGG(MM)+G(K,I)
500 CONTINUE
505 CONTINUE
510 CONTINUE
RETURN
END

```



-----  
 THIS IS THE MAIN PROGRAM FOR THE SOLUTION OF THE FIELD EQUATIONS.  
 THE DATA FROM THE TAPE IS READ OUT, NODAL CONSTANTS ARE EVALUATED,  
 THE INTEGRATION PACKAGE IS INITIATED, ETC.  
 -----

```

REAL*8 SIGF, SIGA, C1, C2, C4, C5
REAL*8 GGG, BIGG, BIGGG, BIGH
REAL*8 WL, WS, S, CL1, CL2, CL3, FN, DL1, DL2, DL3, DS, DETJ
INTEGER*2 NAME, JA, JB, NNZ, JC, NUMNP, NUMEL, NELDOF, CONN
COMMON/GTRY2/NAME(4615), JA(183), JB(184), NNZ, JC, NUMNP, NUMEL, NELDOF,
1CONN(128,15)
COMMON/GTRY1/X(505), P(505), Z(505)
COMMON/XOCAL/KX(15), YY(15), ZZ(15), PPSI(15)
DIMENSION W(700), Y(7,183)
COMMON/DELAY/SUM(183), PSI(93)
COMMON/TEMP/C3(183), C6(15)
DIMENSION SIGF(93), SIGA(183)
COMMON/COEF/C1(183), C2(183), C4(183), C5(183)
COMMON/GGGMAT/GGG(15,15)
COMMON/SYSMT1/BIGG(4615)
COMMON/SYSMT2/BIGGG(4615)
COMMON/SYSMT3/BIGH(4615)
COMMON/GAUSS/WL(7), WS(5), S(5), CL1(7), CL2(7), CL3(7), FN(15,21), DL1(1
15,21), DL2(15,21), DL3(15,21), DS(15,21), DETJ(21)
NNZ=183
NUMEL=128
NUMNP=505
NELDOF=15
JC=4615
FMAX=1.0E10
DO 771 I=1, NUMEL
  READ(4) (CONN(I, J), J=1, NELDOF)
771 CONTINUE
  (JA(I), I=1, NNZ)
  READ(4) (JB(I), I=1, NNZ)
  READ(4) (NAME(I), I=1, JC)
  READ(4) (X(I), I=1, NUMNP)
  READ(4) (P(I), I=1, NUMNP)
  READ(4) (Z(I), I=1, NUMNP)
  READ(4) (BIGG(I), I=1, JC)
  READ(4) (BIGGG(I), I=1, JC)
DO 772 I=1, 3
  READ(4) WS(I), S(I)
772 CONTINUE

```

M-010  
M-015  
M-020  
M-025  
M-055  
M-060

M-075  
M-080  
M-085  
M-095  
M-100

M-105



```

DO 773 I=1,7
READ(4) WL(I),CL1(I),CL2(I),CL3(I)
CONTINUE
DO 781 I=1,NELDOF
DO 782 J=1,21
READ(4) FN(I,J),DL1(I,J),DL2(I,J),DL3(I,J),DS(I,J)
CONTINUE
CONTINUE
DO 460 I=1,NNZ
Y(1,I)=FMAX
Y(2,I)=0.0
SUM(I)=0.0
IF(I.GT.93) GJ TO 460
PSI(I)=FMAX
CONTINUE
WRITE(6,3)
3 FORMAT(1,3)
SIGF(I)=0.0057450, SIGA(I)=0.014010,0.0080 .)
V=4.800E07
Q=0.4349710
BETA=0.006420
B=-0.00400
DO 445 I=1,NNZ
IF(I.GT.93) GJ TO 420
D=0.9130
SIGF(I)=0.0057450
SIGA(I)=0.014010
ZNU=2.54
C1(I)=V*D
C2(I)=V*SIGA(I)-V*(1.0-BETA)*SIGF(I)*ZNU
C3(I)=(2.0179325E-13)*SIGF(I)
C4(I)=-V*(1.0-BETA)*SIGA(I)*B
C5(I)=-V*BETA*Q*ZNU*SIGF(I)
GO TO 445
420 D=1.200
SIGA(I)=0.00800
C1(I)=V*D
C2(I)=V*SIGA(I)
C3(I)=0.0
C4(I)=0.0
C5(I)=0.0
CONTINUE
TEND=0.0010
NY=NNZ
NL=0
H=1.0E-19
T=0.0
HMIN=1.0E-20
HMAX=0.10

```

M-520  
M-535  
M-545  
M-540  
M-550

M-565  
M-570  
M-575  
M-580  
M-585  
M-590









```

SDESOL. JSKF > 0 INDICATES A CONTINUATION OF THE SDE
PREVIOUS CALL TO SDESOL. JSKF < -1 MAY HAVE RESULTS SDE
FROM THE USER NEGLECTING TO TEST FOR ERROR RETURNS SDE
FROM SDESOL. BECAUSE OF THIS POSSIBILITY, JSKF < -1 SDE
RESULTS IN TERMINATION OF THE RUN WITH THE SDE
APPROPRIATE COMMENT. SDE
ON OUTPUT, JSKF CONSISTS OF TWO DIGITS AND SIGN, SDE
+ OR - QP. Q IS THE ORDER OF THE FORMULA CURRENTLY SDE
BEING USED. P INDICATES THE TYPE OF RETURN, AS SDE
FOLLOWS. SDE
JSKF > 0, P = 1 IS THE NORMAL RETURN SDE
JSKF < 0, P IS AN ERROR RETURN, WITH THE FOLLOWING SDE
MEANINGS. SDE
P = 1 ERROR TEST FAILURE FOR H > HMIN SDE
P = 3 CORRECTOR FAILED TO CONVERGE FOR H > HMIN SDE
P = 4 ORDER METHOD FAILED TO CONVERGE FOR FIRST SDE
P = 5 ORDER METHOD FROM SUBROUTINE NUTSL SDE
P = 6 ERROR RETURN FROM SUBROUTINE DERVAL SDE
MAXDER - MAXIMUM ORDER DERIVATIVE THAT SHOULD BE USED IN SDE
METHOD. IT MUST BE NO GREATER THAN SIX. SDE
IPRT - INTERNAL PRINT CONTROL INDICATOR FOR LDASUB. SDE
IPRT = 0 NO PRINT SDE
IPRT > 0 PRINT COUNTERS, STEPSIZE, CURRENT TIMES SDE
AND VALUES OF DEPENDENT VARIABLES AT SDE
EACH STEP. SDE
H - CURRENT STEPSIZE. AN INITIAL VALUE MUST BE SUPPLIED SDE
3JT NEED NOT BE THE ONE WHICH MUST BE USED, SINCE THE SDE
SUBROUTINE WILL CHOOSE A SMALLER ONE IF NECESSARY TO SDE
KEEP THE ERROR PER STEP SMALLER THAN THE SPECIFIED SDE
VALUE. IT IS BETTER TO UNDERESTIMATE THE INITIAL SDE
STEPSIZE THAN TO OVERESTIMATE IT. THE STEPSIZE IS SDE
VARIABLELY NOT CHANGED BY THE USER. SDE
HMIN - MINIMUM STEPSIZE ALLOWED SDE
HMAX - MAXIMUM STEPSIZE ALLOWED SDE
RMSEPS - THE ERROR TEST CONSTANT. THE ROOT-MEAN-SQUARE OF SDE
THE SINGLE STEP ERROR ESTIMATES, ER(I), DIVIDED BY SDE
YMAX(I) = (MAXIMUM TO CURRENT TIME OF Y(I)) MUST BE SDE
LESS THAN EPS. THE STEPSIZE AND/OR THE ORDER SDE
ARE VARIED TO ACHIEVE THIS. SDE
W - SCATCH STORAGE ARRAY. MUST BE AT LEAST 13*NY + 5*NL SDE
LOCATIONS, PLUS THOSE REQUIRED FOR STORAGE OF THE SDE
MATRIX PW. (SEE DESCRIPTION OF SUBROUTINE JACMAT). SDE
THE STORAGE OF PW WILL NORMALLY REQUIRE NO MORE THAN SDE
N*2 + 2*N LOCATIONS, AND IF COMPACT STORAGE TECH- SDE
NIQUES ARE USED, CAN BE MUCH FEWER. SDE
-----SDE

```



```

C      DIMENSION Y(7,1), YL(1), W(1)
C      IF (JSKF.GT.0) GO TO 120
C      IF (JSKF.LT.-1) GO TO 140
C      N = NY+NL
C      IF (JSKF.LT.0) GO TO 110
C
C      IF THIS IS THE FIRST ENTRY, OBTAIN VALJES OF THE DERIVATIVES.
C      CALL Derval (Y,YL,T,N,NY,W,KRETR)
C      IF (KRETR.NE.0) GO TO 130
C
C      NOW SET UP STORAGE BLOCKS IN THE W ARRAY.  THIS NEEDS TO BE DONE
C      ONLY INITIALLY AND ON RESTARTS.
C
C      THE ARRAY   SAVE   STARTS AT LOCATION
C      THE ARRAY   YLSV   STARTS AT LOCATION
C      THE ARRAY   YMAX   STARTS AT LOCATION
C      THE ARRAY   ERSV   STARTS AT LOCATION
C      THE ARRAY   ESF    STARTS AT LOCATION
C      THE ARRAY   F1     STARTS AT LOCATION
C      THE ARRAY   DY     STARTS AT LOCATION
C      THE MATRIX   PW     STARTS AT LOCATION
C
C      110 NSVL = 7*NY+1
C          NYMAX = NSVL+NL
C          NER = NSVL+NY
C          NESV = NER+NY
C          NF1 = NESV+NY
C          NDY = NF1+N
C          NPW = NDY+N
C          JS = JSKF
C          CALL LDASUB (Y,YL,T,TEND,N,NY,M,JS,KF,MAXDER,IPRT,H,HMIN,HMAX,
C            IRMSEPS,W,W(NSVL),W(NYMAX),W(NER),W(NESV),W(NFY),W(NPW))
C
C      CODE JSKF ON RETURN FROM LDASUB
C
C      JSKF = ISIGN(JS*10+IABS(KF),KF)
C      RETURN
C      JSKF = -6
C      RETURN
C      PRINT 1, JSKF
C      STOP
C
C      1 FORMAT ('OIT IS AN ERROR TO ENTER SDESCL WITH JSKF = ',I10//
C        1, ' RUN HAS BEEN TERMINATED.')
C      END
C
C      SUBROUTINE LDASUB (Y,YL,T,TEND,N,NY,M,JSTART,KFLAG,MAXOR,IPRT,H,

```





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1HMIN,HMAX,RMSEPS,SAVE,YLSV,YMAX,ER,ESV,FI,DY,PW)
SUBROUTINE LDASUB IS A MODIFICATION OF SUBROUTINE DFASUB
WHICH IS DUE TO R. L. BROWN AND C. W. GEAR. DFASUB IS DOCUMENTED
IN THE REPORT
DOCUMENTATION FOR DFASUB--
BY R. L. BROWN AND C. W. GEAR
REPORT UIUCDCS-R-73-575, JULY 1973
UNIVERSITY OF ILLINOIS AT URBANA-CHAMPAIGN
URBANA, ILLINOIS 61801
THIS REPORT IS AVAILABLE FROM THE NATIONAL TECHNICAL INFORMATION
SERVICE OF THE U. S. DEPARTMENT OF COMMERCE UNDER ACCESSION NUMBER
C00-1459-225.

THE MODIFICATION HERE IS DOCUMENTED IN THE REPORT
A PROGRAM FOR THE NUMERICAL SOLUTION OF LARGE SPARSE SYSTEMS OF
ALGEBRAIC AND IMPLICITLY DEFINED STIFF DIFFERENTIAL EQUATIONS
BY RICHARD FRANK
REPORT NPS53FE76051, MAY 1976
NAVAL POSTGRADUATE SCHOOL
MONTEREY, CALIFORNIA 93940
-----

THE CALLING SEQUENCE FOR LDASUB IS

CALL LDASUB(Y,YL,T,TEND,V,NY,M,JSTART,KFLAG,MAXOR,IPRT,H,HMIN,
HMAX,RMSEPS,SAVE,YLSV,YMAX,ER,ESV,FI,DY,PW)

WHERE THE PARAMETERS ARE DEFINED AS FOLLOWS.
Y - ARRAY DIMENSIONED (7,NY). THIS ARRAY CONTAINS THE
DEPENDENT VARIABLES AND THEIR SCALED DERIVATIVES.
Y(J+1,I) CONTAINS THE J-TH DERIVATIVE OF THE I-TH VARIABLE
TABLE TIMES H**J/J-FACTIAL, WHERE H IS THE CURRENT
STEP SIZE. ON FIRST ENTRY THE CALLER SUPPLIES THE
INITIAL VALUES OF EACH VARIABLE IN Y(1,I) AND AN
ESTIMATE OF THE INITIAL VALUES OF THE DERIVATIVES
IN Y(2,I). ON SUBSEQUENT ENTRIES IT IS ASSUMED THAT
THE ARRAY HAS NOT BEEN CHANGED. TO INTERPOLATE TO
NON-MESH POINTS, THESE VALUES CAN BE USED AS FOLLOWS.
IF H IS THE CURRENT STEP SIZE AND VALUES AT TIME T+E
NEEDED, LET S = E/H AND THEN

      NQ      SUM Y(J+1,I)*S**J
      I-TH VARIABLE AT T+E IS
      J=0

      THE VALUE OF NQ IS OBTAINED IN THE CALLING PROGRAM

```





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3Y  NQ = JSTART.
-   ARRAY OF NL = N - NY VARIABLES WHICH APPEAR LINEARLY.
-   THE USER SUPPLIES INITIAL VALUES FOR THESE VARIABLES.
-   CURRENT VALUE OF THE INDEPENDENT VARIABLE (TIME)
-   END TIME
-   TOTAL NUMBER OF VARIABLES
-   NUMBER OF DIFFERENTIAL EQUATIONS AND NONLINEAR
-   VARIABLES.
-   NUMBER OF VARIABLES INCLUDED IN THE ERROR TEST.
-   THIS NUMBER CAN BE NO GREATER THAN NY. IF IT IS
-   GREATER THAN NY, NY VARIABLES ARE USED IN THE ERROR
-   TEST.
-   INPUT AND OUTPUT INDICATOR.
-   ON INPUT
-   ON INPUT <0
-   ON INPUT =0
-   ON INPUT >0
-   ON OUTPUT
-   ON OUTPUT, JSTART IS SET TO THE VALUE OF NQ, THE
-   ORDER OF THE FORMULA CURRENTLY BEING USED.
-   THE COMPLETION CODE INDICATOR, WITH THE FOLLOWING
-   MEANINGS
-   +1 THE INTEGRATION WAS SUCCESSFUL
-   -1 ERROR TEST FAILED FOR H > HMIN
-   -3 CORRECTOR FAILED TO CONVERGE FOR H > HMIN
-   -4 CORRECTOR FAILED TO CONVERGE FOR FIRST
-   ORDER METHOD
-   -5 ERROR RETURN FROM SUBROUTINE NUTSL
-   MAXOR ORDER DERIVATIVE THAT SHOULD BE USED IN THE
-   METHOD. IT MUST BE NO GREATER THAN SIX. IF IT IS
-   GREATER THAN SIX, THE MAXIMUM ORDER USED WILL BE SIX.

```

```

LDA 500
LDA 510
LDA 520
LDA 530
LDA 540
LDA 550
LDA 560
LDA 570
LDA 580
LDA 590
LDA 600
LDA 610
LDA 620
LDA 630
LDA 640
LDA 650
LDA 660
LDA 670
LDA 680
LDA 690
LDA 700
LDA 710
LDA 720
LDA 730
LDA 740
LDA 750
LDA 760
LDA 770
LDA 780
LDA 790
LDA 800
LDA 810
LDA 820
LDA 830
LDA 840
LDA 850
LDA 860
LDA 870
LDA 880
LDA 890
LDA 900
LDA 910
LDA 920
LDA 930
LDA 940
LDA 950
LDA 960
LDA 970

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CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC

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C      YMAX(J) = AMAX1(1.,ABS(Y(1,J)))
130  Y(2,J) = Y(2,J)*H
C
C      NQ = 1
C      BR = 1.
C      ASSIGN 190 TO IRET
C      SET COEFFICIENTS FOR THE ORDER CURRENTLY BEING USED.
C      E IS A TEST FOR ERRORS OF THE CURRENT ORDER NQ
C      EUP IS TO TEST FOR INCREASING THE ORDER, EDWN FOR DECREASING THE
C      ORDER.
C
140  K = NQ*(NQ-1)/2
C      CALL COPYZ (A(2),COF(K+1),NQ)
C      K = NQ+1
C      IDQUB = NQ
C      ENQ1 = .5/NQ
C      ENQ2 = .5/K
C      ENQ3 = .5/(NQ+2)
C      PEP SH = EPS**2
C      E = PERT(NQ,1)*PEP SH
C      EUP = PERT(NQ,2)*PEP SH
C      EDWN = PERT(NQ,3)*PEP SH
C      BND = (EPS*ENQ3)**2
C      IWEVAL = 1
C      GO TO IRET, (190,200,490,570)
150  IF (H.EQ.HNEW) GO TO 190
C
C      IF CALLER HAS CHANGED H, RESCALE DERIVATIVES TO REFLECT THAT HNEW
C      WAS USED ON THE LAST CALL.
C
C      R = H/HNEW
C      ASSIGN 190 TO IRET
C      GO TO 610
C
C      SET JSTART TO NQ, THE CURRENT ORDER OF THE METHOD, BEFORE EXIT,
C      AND SAVE THE CURRENT STEP SIZE IN HNEW.
C
160  JSTART = NQ
C      HNEW = H
C      RETURN
170  NS = NS+1
C      IF (IPRT.LE.0) GO TO 180
C
C      PRINT DATA IF DESIRED BY USER
C
C      PRINT 1, NS, NQ, H, T, (Y(1,I), I=1, NY)
C      IF (NL.GT.0) PRINT 2, (YL(I), I=1, NL)

```

```

LDA 1940
LDA 1950
LDA 1960
LDA 1970
LDA 1980
LDA 1990
LDA 2000
LDA 2010
LDA 2020
LDA 2030
LDA 2040
LDA 2050
LDA 2060
LDA 2070
LDA 2080
LDA 2090
LDA 2100
LDA 2110
LDA 2120
LDA 2130
LDA 2140
LDA 2150
LDA 2160
LDA 2170
LDA 2180
LDA 2190
LDA 2200
LDA 2210
LDA 2220
LDA 2230
LDA 2240
LDA 2250
LDA 2260
LDA 2270
LDA 2280
LDA 2290
LDA 2300
LDA 2310
LDA 2320
LDA 2330
LDA 2340
LDA 2350
LDA 2360
LDA 2370
LDA 2380
LDA 2390
LDA 2400
LDA 2410

```





```

180 CONTINUE
IF (KFLAG.LT.0) GO TO 160
IF (T.GE.TEND) GO TO 160
C-----
C TAKE ANOTHER STEP IF T < TEND
C-----
C JSTART = 1
C-----
C SAVE DATA FOR TRIAL WITH A SMALLER TIMESTEP IF THIS STEP FAILS
C-----
190 CALL COPYZ (SAVE,Y,LCOPYZ)
CALL COPYZ (YLSV,YL,LCOPYL)
RACUM = 1.
KFLAG = 1
HOLD = H
NQOLD = NQ
TOLD = T
T = T+H
HINV = 1./H
C-----
C COMPUTE PREDICTED VALUES BY EFFECTIVELY MULTIPLYING DERIVATIVE
C VECTOR BY PASCAL TRIANGLE MATRIX
C-----
DO 210 J=2,K
J3 = K+J-1
C
DO 210 J1=J,K
J2 = J3-J1
C
DO 210 I=1,NY
210 Y(J2,I) = Y(J2,I)+Y(J2+1,I)
C
DO 220 I=1,NY
220 ER(I) = 0.
C-----
C DO UP TO THREE CORRECTOR ITERATIONS. CONVERGENCE IS OBTAINED WHEN
C CHANGES ARE LESS THAN BND WHICH IS DEPENDENT ON THE ERROR TEST
C CONSTANT. THE SUM OF CORRECTIONS IS ACCUMULATED IN ER(I). IT IS
C EQUAL TO THE K-TH DERIVATIVE OF Y TIMES H**K/(K-FACTORIAL*A(K)),
C AND THUS IS PROPORTIONAL TO THE ACTUAL ERRORS TO THE LOWEST POWER
C OF H PRESENT, WHICH IS H**K.
C-----
DO 270 L=1,3
CALL DIFFUN (Y,YL,T,HINV,DY)
C-----

```



```

C C C C C C C
IF (IWEVAL.LT.1) GO TO 230
-----
IF THERE HAS BEEN A CHANGE OF ORDER OR THERE HAS BEEN TROUBLE
WITH CONVERGENCE, PW IS RE-EVALUATED PRIOR TO STARTING THE
CORRECTOR ITERATION. IWEVAL IS THEN SET TO -1 AS AN INDICATOR
THAT IT HAS BEEN DONE. NEWPW IS SET NONZERO TO INDICATE TO
SUBROUTINE NUTSL THAT A NEW PW HAS BEEN PROVIDED.
-----
LDA 2900
LDA 2910
LDA 2920
LDA 2930
LDA 2940
LDA 2950
LDA 2960
LDA 2970
-----
CALL JACMAT (Y, YL, T, HINV, A(2), N, NY, EPS, DY, F1, PW)
KFLAG = 1
IWEVAL = -1
NW = NW+1
NEWPW = 1
230 CALL NUTSL (PW, DY, F1, N, NY, EPS, YMAX, NEWPW, KRRET)
IF (KRRET.NE.0) GO TO 600
IF (NL.LE.0) GO TO 250
-----
DO 240 I=1,NL
240 YL(I) = YL(I)-F1(I+NY)
C
C
250 CONTINUE
DEL = 0.
C
DO 260 I=1,NY
Y(1,I) = Y(1,I)-F1(I)
Y(2,I) = Y(2,I)+A(2)*F1(I)
ER(I) = ER(I)+F1(I)
DEL = DEL+(F1(I)/AMAX1(YMAX(I),ABS(Y(1,I))))**2
260 CONTINUE
C
IF (L.GE.2) BR = AMAX1(.9*BR,DEL/DEL1)
DEL1 = DEL
IF (AMIN1(DEL,BR*DEL*2.).LE.BND) GO TO 330
270 CONTINUE
C
THE CORRECTOR ITERATION FAILED TO CONVERGE IN 3 TRIES. VARIOUS
POSSIBILITIES ARE CHECKED FOR. IF H IS ALREADY HMIN AND PW HAS
ALREADY BEEN RE-EVALUATED, A NO CONVERGENCE EXIT IS TAKEN.
OTHERWISE THE MATRIX PW IS RE-EVALUATED AND/OR (IN THAT ORDER) THE
STEP IS REDUCED TO TRY AND GET CONVERGENCE.
-----
T = TOLD
IF (IWEVAL) 280,300,290
280 IF (H.LE.HMIN*1.00001) GO TO 310
290 RACUM = RACUM*.25
300 CONTINUE
GO TO 560
-----
LDA 3250
LDA 3260
LDA 3270
LDA 3280
LDA 3290
LDA 3300
LDA 3310
LDA 3320
LDA 3330
LDA 3340
LDA 3350
LDA 3360
LDA 3370

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```

C      DO 390 I=1,NY
C      390 SAVE(2,I) = Y(2,I)
C      LDA 3850
C      400 KFLAG = KFLAG-2
C      IF (H.LE.HMIN) GO TO 550
C      T = TOLD
C      IF (KFLAG.LE.-5) GO TO 530
C      410 PR2 = (D/E)**NQ2*1.2
C      L = 0
C      IF (NQ.LE.1) GO TO 430
C      D = 0.
C      DO 420 J=1,M1
C      YM = AMAX1(ABS(Y(1,J)),YMAX(J))
C      420 D = D+((Y(K,J)/YM)**2)
C      LDA 4000
C      PR1 = (D/EDWN)**ENQ1*1.3
C      IF (PR1.GE.PR2) GO TO 430
C      PR2 = PR1
C      L = -1
C      430 IF (KFLAG.LT.0.JR.NQ.GE.MAXDER) GO TO 450
C      D = 0
C      DO 440 J=1,M1
C      YM = AMAX1(ABS(Y(1,J)),YMAX(J))
C      440 D = D+((ER(J)-ESV(J))/YM)**2
C      LDA 4010
C      PR1 = (D/EUP)**ENQ3*1.4
C      IF (PR1.GE.PR2) GO TO 450
C      PR2 = PR1
C      L = 1
C      450 R = 1./AMAX1(PR2,1.E-5)
C      IF (KFLAG.LT.0.JR.R.GE.1.1) GO TO 460
C      IDOUB = 9
C      GO TO 510
C      460 NEWQ = NQ+L
C      K = NEWQ+1
C      IF (NEWQ.LE.NQ) GO TO 480
C      R1 = A(NEWQ)/FLOAT(NEWQ)
C      DO 470 J=1,NY
C      470 Y(K,J) = ER(J)*R1
C      LDA 4250
C      480 CONTINUE
C      IF THE STEP WAS OKAY, SCALE THE Y VARIABLES IN ACCORDANCE
C      WITH THE NEW VALUE OF H. IF KFLAG < 0, HOWEVER, USE THE
C      SAVED VALUES (IN SAVE AND YLSV). IN EITHER CASE, IF THE ORDER

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```

C      HAS CHANGED IT IS NECESSARY TO FIX CERTAIN PARAMETERS BY CALLING
C      THE PROGRAM SEGMENT AT STATEMENT NUMBER 140.
C      -----
      IDOUB = NQ
      IF (NEWQ.EQ.NQ) GO TO 490
      NQ = NEWQ
      ASSIGN 490 TO IRET
      GO TO 140
490  IF (KFLAG.GT.J) GO TO 500
      RACUM = RACUM*R
      GO TO 560
500  R = AMAX1(AMIN1(HMAX/H,R),HMIN/H)
      H = H*R
      IWEVAL = 1
      ASSIGN 510 TO IRET
      GO TO 610
C
C      510 DO 520 I=1,M1
C      520 YMAX(I) = AMAX1(ABS(Y(1,I)),YMAX(I))
C      GO TO 170
C      -----
C      THE ERROR TEST HAS NOW FAILED THREE TIMES, SO THE DERIVATIVES ARE
C      IN BAD SHAPE. RETURN TO FIRST ORDER METHOD AND TRY AGAIN. OF
C      COURSE, IF NQ = 1 ALREADY, THEN THERE IS NO HOPE AND WE EXIT WITH
C      KFLAG = -4.
C      -----
530  IF (NQ.EQ.1) GO TO 540
      NQ = 1
      IDOUB = 1
      ASSIGN 570 TO IRET
      GO TO 140
540  NQOLD = 1
      KFLAG = -4
      GO TO 320
550  KFLAG = -1
      GO TO 170
C      -----
C      THIS SECTION RESTORES THE SAVED VALUES OF Y AND YL, SCALING THE
C      Y DERIVATIVES AS NECESSARY, AND THEN RETURNS TO THE PREDICTOR LOOP
C      -----
560  H = HOLD*RACUM
      H = AMAX1(HMIN,AMIN1(H,HMAX))
570  RACUM = H/HOLD
      R1 = 1.
C
C      DO 580 J=2,K
      R1 = R1*RACUM
C
C      580
C
C      590
C
C      600
C
C      610
C
C      620
C
C      630
C
C      640
C
C      650
C
C      660
C
C      670
C
C      680
C
C      690
C
C      700
C
C      710
C
C      720
C
C      730
C
C      740
C
C      750
C
C      760
C
C      770
C
C      780
C
C      790
C
C      800
C
C      810
C
C      820
C
C      830
C
C      840
C
C      850
C
C      860
C
C      870
C
C      880
C
C      890
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C      960
C
C      970
C
C      980
C
C      990
C
C      1000
C
C      1010
C
C      1020
C
C      1030
C
C      1040
C
C      1050
C
C      1060
C
C      1070
C
C      1080
C
C      1090
C
C      1100
C
C      1110
C
C      1120
C
C      1130
C
C      1140
C
C      1150
C
C      1160
C
C      1170
C
C      1180
C
C      1190
C
C      1200
C
C      1210
C
C      1220
C
C      1230
C
C      1240
C
C      1250
C
C      1260
C
C      1270
C
C      1280
C
C      1290
C
C      1300
C
C      1310
C
C      1320
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C      1330
C
C      1340
C
C      1350
C
C      1360
C
C      1370
C
C      1380
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C      DO 580 I=1,NY
580   Y(J,I) = SAVE(J,I)*R1
C
C      DO 590 I=1,NY
590   Y(1,I) = SAVE(1,I)
C      CALL COPYZ (YL,YLSV,LCOPYL)
IWEVAL=1
GO TO 200
600   KFLAG = 5
GO TO 160
C
C      THIS SECTION SCALES THE Y DERIVATIVES BY R**J
C
C      610 R1 = 1.
C
C      DO 620 J=2,K
R1 = R1*R
C
C      DO 620 I=1,NY
620   Y(J,I) = Y(J,I)*R1
C
C      GO TO IRET, (190,510)
C
C      THIS SECTION ALLOWS FOR RESTARTS AFTER SOLVING ANOTHER PROBLEM, OR
HAVING TERMINATED THE CURRENT COMPUTER RUN. SUBROUTINE LDASAV
SAVES THE NECESSARY VALUES WHICH ARE INTERNAL TO LDASUB. FOR
DOUBLE PRECISION, WITH COPYZ IN SINGLE PRECISION, THE NUMBER OF
LOCATIONS TO BE SAVED AND RESTORED, LCPYS AND LCPYR, MUST BE
SET TO 58.
C      IT IS ASSUMED THAT IN ADDITION TO THE VARIABLES IN THE ARRAY A
SAVED BY CALLING LDASAV, THE USER ALSO SAVES THE ARRAYS SAVE,
YLSV, YMAX, ESV, AND PW.
C
C      TO RESTART THE USER FIRST CALLS LDARST TO RESTORE THE VALUES SAVED
BY LDASAV, THEN RE-ENTERS LDASUB WITH JSTART < 0, AND WITH THE
OTHER PARAMETERS THE SAME AS RETURNED FROM THE LAST ENTRY TO
LDASUB, PARTICULARLY THOSE ARRAYS MENTIONED ABOVE.
C
C      ENTRY LDASAV(SAV)
LCOPYS = 29
CALL COPYZ (SAV,A,LCOPYS)
CALL COPYZ (SAVE,Y,LCOPYY)
CALL COPYZ (YLSV,YL,LCOPYL)
RETURN
C

```

```

LDA 4820
LDA 4830
LDA 4840
LDA 4850
LDA 4860
LDA 4870
LDA 4880
LDA 4890
LDA 4900
LDA 4910
LDA 4920
LDA 4930
LDA 4940
LDA 4950
LDA 4960
LDA 4970
LDA 4980
LDA 4990
LDA 5000
LDA 5010
LDA 5020
LDA 5030
LDA 5040
LDA 5050
LDA 5060
LDA 5070
LDA 5080
LDA 5090
LDA 5100
LDA 5110
LDA 5120
LDA 5130
LDA 5140
LDA 5150
LDA 5160
LDA 5170
LDA 5180
LDA 5190
LDA 5200
LDA 5210
LDA 5220
LDA 5230
LDA 5240
LDA 5250
LDA 5260
LDA 5270
LDA 5280
LDA 5290

```



```

ENTRY LDARST(SAV)
LCOPYR = 29
CALL COPYZ (A, SAV, LCOPYR)
RETURN
-----
1 FORMAT (2I5, I2, 1P2E10.2, 7E14.6 / (32X, 7E14.6))
2 FORMAT (32X, 1P7E14.6)
3 FORMAT (11N = , I3, , NL = , I3, , RMSEPS = , 1PE9.2, , TEND = ,
1,E9.2, , H = , E9.2 //)
4 FORMAT (11NS NW Q H , 8X, , T , 8X, , Y(1,*) AND YL(*)' //)
END
-----
SUBROUTINE COPYZ(S, Y, L)
DIMENSION S(1), Y(1)
-----
THIS SUBROUTINE COPIES THE ARRAY Y, OF LENGTH L, INTO THE ARRAY S
-----
IF(L.LE.0) RETURN
DO 100 J=1, L
S(J) = Y(J)
100 RETURN
END
-----
SUBROUTINE DERVAL (Y, YL, T, N, NY, W, KERET)
10 DER
20 DER
30 DER
40 DER
50 DER
60 DER
70 DER
80 DER
90 DER
100 DER
110 DER
120 DER
130 DER
140 DER
150 DER
160 DER
170 DER
180 DER
190 DER
200 DER

THIS SUBROUTINE CALCULATES THE INITIAL VALUES OF THE DERIVATIVES
IN THE GENERAL CASE. IT IS WRITTEN SO THAT IT SHOULD WORK IF THE
FIRST NY EQUATIONS ALL INVOLVE DERIVATIVES. IT ATTEMPTS TO SOLVE
THE FIRST NY EQUATIONS USING NEWTON'S METHOD, BUT SINCE IT TRIES
TO EVALUATE DF/DY BY CALLING JACMAT IN SUCH A WAY AS TO MAKE THE
DF/DY TERM INSIGNIFICANT, IT IS POSSIBLE THAT IT MAY FAIL FOR THAT
REASON. IT MAY FAIL FOR OTHER REASONS, AS WELL. IF IT DOES FAIL
THE USER CAN SUPPLY HIS OWN VERSION OF DERVAL, OR MODIFY THIS
ROUTINE IN SUITABLE FASHION. THIS ROUTINE ASSUMES THAT VALUES OF
THE LINEAR VARIABLES HAVE BEEN SUPPLIED PREVIOUSLY. IF THOSE
MUST BE SOLVED FOR SIMULTANEOUSLY WITH THE DERIVATIVES, THE USER
MUST SUPPLY HIS OWN VERSION OF DERVAL.

THE CALLING SEQUENCE FOR THIS SUBROUTINE IS
CALL Derval(Y, YL, T, N, NY, W, KERET)
WHERE THE PARAMETERS ARE DEFINED AS FOLLOWS

```











```

      IF (ER.LT.EPS2) GO TO 150
140  CONTINUE
      GO TO 170
      DO 160 I=1,NY
160  Y(2,I) = Y(3,I)
      RETURN
170  KERET = 1
      RETURN
      END
DER 690
DER 700
DER 710
DER 720
DER 730
DER 740
DER 750
DER 760
DER 770
DER 780
DER 790
DER 800

```

-----  
 THIS SUBROUTINE IS REQUIRED BY THE TIME INTEGRATION PACKAGE AND  
 MUST BE SUPPLIED BY THE USER. ITS PURPOSE IS TO EVALUATE THE  
 FUNCTION AT CURRENT VALUES OF THE VARIABLES.  
 -----

```

SUBROUTINE DIFFUN(Y,YL,T,HINV,DY)
REAL*8 C1,C2,C3,C4,C5
REAL*8 BIGG,BIGG,BIGH
INTEGER*2 NAME,JA,JB,NNZ,JC,NUMNP,NUMEL,NELDOF,CONN
COMMON/GTRY2/NAME(4615),JA(183),JB(184),NNZ,JC,NUMNP,NUMEL,NELDOF,
1  CONN(128,15)
DIMENSION Y(7,1),YL(1),DY(1)
COMMON/DELAY/SUM(183),PSI(93)
COMMON/COEF/C1(183),C2(183),C4(183),C5(183)
COMMON/SYSMT1/BIGG(4615)
COMMON/SYSMT2/BIGG(4615)
COMMON/SYSMT3/BIGH(4615)
DATA TOLD/-1.0E-30/
IF(T.EQ.TOLD) GO TO 97
CALL SUMT(Y,TOLD,T-TOLD)
TOLD=T
96  CONTINUE
97  CALL FEEDBK(Y,TOLD)
DO 110 I=1,NNZ
IF(I.GT.93) GO TO 401
PSI(I)=Y(1,I)
401 CONTINUE
DY(I)=0.0
JBB=JB(I)
JE=JB(I)+JA(I)-1
DO 120 J=JBB,JE
LL=NAME(J)
D-005
D-010
D-015
D-020
D-025
D-035
D-050
D-055
D-065
D-075
D-080
D-085
D-090
D-100
D-110
D-115
D-120
D-125
D-130
D-135
D-140

```



D-145  
D-180  
D-185  
D-195  
D-200

```

IF(LL-GT.NNZ) GO TO 120
DY(I)=DY(I)+BIGG(J)*(HINV*Y(2,LL)+C2(LL)*Y(1,LL)+C5(LL)*SUM(LL))+C
11(LL)*BIGG(J)*Y(1,LL)+C4(LL)*BIGH(J)*Y(1,LL)
120 CONTINUE
110 CONTINUE
RETURN
END

```

THIS SUBROUTINE IS REQUIRED BY THE TIME INTEGRATION PACKAGE AND  
MUST BE SUPPLIED BY THE USER. ITS PURPOSE IS TO EVALUATE THE J  
MATRIX NEEDED WHEN THE CORRECTOR EQUATION IS BEING SOLVED.

```

SUBROUTINE JACMAT(Y,YL,T,HINV,A2,N,NY,EPS,DY,F1,PW)
REAL*8 C1,C2,C4,C5
REAL*8 BIGG,BIGGG,BIGH
INTEGER*2 NAME,JA,J8,NNZ,J2,NUMNP,NUMEL,NELDOF,CONN
COMMON/GTRY2/NAME{4615},JA{183},JB{184},NNZ,J2,NUMNP,NUMEL,NELDOF,
1CONN{128,15}
DIMENSION Y(7,1),YL(1),F1(1),DY(1),PW(1)
COMMON/COEF/C1{183},C2{183},C4{183},C5{183}
COMMON/SYSMT1/BIGG(4615)
COMMON/SYSMT2/BIGGG(4615)
COMMON/SYSMT3/BIGH(4615)
AH=-A2*HINV
DO 300 I=1,NNZ
JB8=JB(I)
JE=JB(I)+JA(I)-1
DO 310 J=JB8,JE
LL=NAME(J)
PW(J)=BIGG(J)*(AH+C2(LL))+C1(LL)*BIGGG(J)+C4(LL)*BIGH(J)
310 CONTINUE
300 CONTINUE
RETURN
END

```

U-115  
U-120  
U-130  
U-135

```
CC-----
C      THIS SUBROUTINE IS REQUIRED BY THE TIME INTEGRATION PACKAGE AND
C      MUST BE SUPPLIED BY THE USER. ITS PURPOSE IS TO SOLVE A LINEAR
C      SYSTEM OF EQUATIONS FOR THE NEWTON ITERATES WHEN THE CORRECTOR
C      EQUATION IS BEING SOLVED.
C-----CC
```

SUBROUTINE NUTSL(PW,DY,FL,N,VY,EPS,YMAX,NEWPW,KRET)  
INTEGER\*2 K,JA,JB,NNZ,JC,NUMNP,NUMEL,NELDOF,CONN  
N-005  
N-010



```

COMMON/GTRY2/((4615),JA(183),JB(184),VVZ,JC,NUMNP,NUMEL,NELDOF,CON
1N(128,15)
DIMENSION PW(1),DY(1),F1(1),YMAX(1)
DATA SPD,SPDM1/1.05,.05/
KRET = 0
EPSS = EPSS**2
EPSA2 = EPSS*.0001
NOIT = N
280 DO 281 I=1,NY
    LL=JB(I)
281 F1(I)=DY(I)/PW(LL)
    DO 287 IT=1,NOIT
        RCH = 0.
        CH = 0.
        DO 285 I=1,NY
            LL=JB(I)
            JBB=JB(I)+1
            JE=JB(I)+JA(I)-1
            FN = DY(I)
            DO 284 J=JBB,JE
                IF(K(J).LE.0.OR.K(J).GT.NY) GO TO 284
                FN=FN-PW(J)*F1(K(J))
                CONTINUE
284 FN=FN/PW(LL)
                FN = FN*SPD - SPDM1*F1(I)
                ACH = F1(I) - FN
                CH = CH + (ACH/YMAX(I))**2
                RCH = RCH + (ACH/AMAX1(ABS(FN),EPS))**2
282 F1(I) = FN
285 IF(RCH.LT.EPSS)GO TO 288
                IF(CH.LE.EPSA2)GO TO 288
                CONTINUE
287 KRET=3
288 CONTINUE
    RETURN
END

```

```

C-----
C THIS SUBROUTINE CALCULATES THE GGG ELEMENT MATRIX USING 21 POINTS
C OF INTEGRATION. IT THEN PUTS THE GGG ELEMENT MATRIX INTO THE BIGH
C SYSTEM MATRIX.
C-----
C
C
C
C

```

```

SUBROUTINE FEEDBK(Y,TOLD)
REAL*8 GGG,BIGH
REAL*8 WL,WS,S,CL1,CL2,CL3,CL4,CL5,CL6,CL7,CL8,CL9,CL10,CL11,CL12,CL13,CL14,CL15,CL16,CL17,CL18,CL19,CL20,CL21,CL22,CL23,CL24,CL25,CL26,CL27,CL28,CL29,CL30,CL31,CL32,CL33,CL34,CL35,CL36,CL37,CL38,CL39,CL40,CL41,CL42,CL43,CL44,CL45,CL46,CL47,CL48,CL49,CL50,CL51,CL52,CL53,CL54,CL55,CL56,CL57,CL58,CL59,CL60,CL61,CL62,CL63,CL64,CL65,CL66,CL67,CL68,CL69,CL70,CL71,CL72,CL73,CL74,CL75,CL76,CL77,CL78,CL79,CL80,CL81,CL82,CL83,CL84,CL85,CL86,CL87,CL88,CL89,CL90,CL91,CL92,CL93,CL94,CL95,CL96,CL97,CL98,CL99,CL100,CL101,CL102,CL103,CL104,CL105,CL106,CL107,CL108,CL109,CL110,CL111,CL112,CL113,CL114,CL115,CL116,CL117,CL118,CL119,CL120,CL121,CL122,CL123,CL124,CL125,CL126,CL127,CL128,CL129,CL130,CL131,CL132,CL133,CL134,CL135,CL136,CL137,CL138,CL139,CL140,CL141,CL142,CL143,CL144,CL145,CL146,CL147,CL148,CL149,CL150,CL151,CL152,CL153,CL154,CL155,CL156,CL157,CL158,CL159,CL160,CL161,CL162,CL163,CL164,CL165,CL166,CL167,CL168,CL169,CL170,CL171,CL172,CL173,CL174,CL175,CL176,CL177,CL178,CL179,CL180,CL181,CL182,CL183,CL184,CL185,CL186,CL187,CL188,CL189,CL190,CL191,CL192,CL193,CL194,CL195,CL196,CL197,CL198,CL199,CL200,CL201,CL202,CL203,CL204,CL205,CL206,CL207,CL208,CL209,CL210,CL211,CL212,CL213,CL214,CL215,CL216,CL217,CL218,CL219,CL220,CL221,CL222,CL223,CL224,CL225,CL226,CL227,CL228,CL229,CL230,CL231,CL232,CL233,CL234,CL235,CL236,CL237,CL238,CL239,CL240,CL241,CL242,CL243,CL244,CL245,CL246,CL247,CL248,CL249,CL250,CL251,CL252,CL253,CL254,CL255,CL256,CL257,CL258,CL259,CL260,CL261,CL262,CL263,CL264,CL265,CL266,CL267,CL268,CL269,CL270,CL271,CL272,CL273,CL274,CL275,CL276,CL277,CL278,CL279,CL280,CL281,CL282,CL283,CL284,CL285,CL286,CL287,CL288,CL289,CL290,CL291,CL292,CL293,CL294,CL295,CL296,CL297,CL298,CL299,CL300,CL301,CL302,CL303,CL304,CL305,CL306,CL307,CL308,CL309,CL310,CL311,CL312,CL313,CL314,CL315,CL316,CL317,CL318,CL319,CL320,CL321,CL322,CL323,CL324,CL325,CL326,CL327,CL328,CL329,CL330,CL331,CL332,CL333,CL334,CL335,CL336,CL337,CL338,CL339,CL340,CL341,CL342,CL343,CL344,CL345,CL346,CL347,CL348,CL349,CL350,CL351,CL352,CL353,CL354,CL355,CL356,CL357,CL358,CL359,CL360,CL361,CL362,CL363,CL364,CL365,CL366,CL367,CL368,CL369,CL370,CL371,CL372,CL373,CL374,CL375,CL376,CL377,CL378,CL379,CL380,CL381,CL382,CL383,CL384,CL385,CL386,CL387,CL388,CL389,CL390,CL391,CL392,CL393,CL394,CL395,CL396,CL397,CL398,CL399,CL400,CL401,CL402,CL403,CL404,CL405,CL406,CL407,CL408,CL409,CL410,CL411,CL412,CL413,CL414,CL415,CL416,CL417,CL418,CL419,CL420,CL421,CL422,CL423,CL424,CL425,CL426,CL427,CL428,CL429,CL430,CL431,CL432,CL433,CL434,CL435,CL436,CL437,CL438,CL439,CL440,CL441,CL442,CL443,CL444,CL445,CL446,CL447,CL448,CL449,CL450,CL451,CL452,CL453,CL454,CL455,CL456,CL457,CL458,CL459,CL460,CL461,CL462,CL463,CL464,CL465,CL466,CL467,CL468,CL469,CL470,CL471,CL472,CL473,CL474,CL475,CL476,CL477,CL478,CL479,CL480,CL481,CL482,CL483,CL484,CL485,CL486,CL487,CL488,CL489,CL490,CL491,CL492,CL493,CL494,CL495,CL496,CL497,CL498,CL499,CL500,CL501,CL502,CL503,CL504,CL505,CL506,CL507,CL508,CL509,CL510,CL511,CL512,CL513,CL514,CL515,CL516,CL517,CL518,CL519,CL520,CL521,CL522,CL523,CL524,CL525,CL526,CL527,CL528,CL529,CL530,CL531,CL532,CL533,CL534,CL535,CL536,CL537,CL538,CL539,CL540,CL541,CL542,CL543,CL544,CL545,CL546,CL547,CL548,CL549,CL550,CL551,CL552,CL553,CL554,CL555,CL556,CL557,CL558,CL559,CL560,CL561,CL562,CL563,CL564,CL565,CL566,CL567,CL568,CL569,CL570,CL571,CL572,CL573,CL574,CL575,CL576,CL577,CL578,CL579,CL580,CL581,CL582,CL583,CL584,CL585,CL586,CL587,CL588,CL589,CL590,CL591,CL592,CL593,CL594,CL595,CL596,CL597,CL598,CL599,CL600,CL601,CL602,CL603,CL604,CL605,CL606,CL607,CL608,CL609,CL610,CL611,CL612,CL613,CL614,CL615,CL616,CL617,CL618,CL619,CL620,CL621,CL622,CL623,CL624,CL625,CL626,CL627,CL628,CL629,CL630,CL631,CL632,CL633,CL634,CL635,CL636,CL637,CL638,CL639,CL640,CL641,CL642,CL643,CL644,CL645,CL646,CL647,CL648,CL649,CL650,CL651,CL652,CL653,CL654,CL655,CL656,CL657,CL658,CL659,CL660,CL661,CL662,CL663,CL664,CL665,CL666,CL667,CL668,CL669,CL670,CL671,CL672,CL673,CL674,CL675,CL676,CL677,CL678,CL679,CL680,CL681,CL682,CL683,CL684,CL685,CL686,CL687,CL688,CL689,CL690,CL691,CL692,CL693,CL694,CL695,CL696,CL697,CL698,CL699,CL700,CL701,CL702,CL703,CL704,CL705,CL706,CL707,CL708,CL709,CL710,CL711,CL712,CL713,CL714,CL715,CL716,CL717,CL718,CL719,CL720,CL721,CL722,CL723,CL724,CL725,CL726,CL727,CL728,CL729,CL730,CL731,CL732,CL733,CL734,CL735,CL736,CL737,CL738,CL739,CL740,CL741,CL742,CL743,CL744,CL745,CL746,CL747,CL748,CL749,CL750,CL751,CL752,CL753,CL754,CL755,CL756,CL757,CL758,CL759,CL760,CL761,CL762,CL763,CL764,CL765,CL766,CL767,CL768,CL769,CL770,CL771,CL772,CL773,CL774,CL775,CL776,CL777,CL778,CL779,CL780,CL781,CL782,CL783,CL784,CL785,CL786,CL787,CL788,CL789,CL790,CL791,CL792,CL793,CL794,CL795,CL796,CL797,CL798,CL799,CL800,CL801,CL802,CL803,CL804,CL805,CL806,CL807,CL808,CL809,CL810,CL811,CL812,CL813,CL814,CL815,CL816,CL817,CL818,CL819,CL820,CL821,CL822,CL823,CL824,CL825,CL826,CL827,CL828,CL829,CL830,CL831,CL832,CL833,CL834,CL835,CL836,CL837,CL838,CL839,CL840,CL841,CL842,CL843,CL844,CL845,CL846,CL847,CL848,CL849,CL850,CL851,CL852,CL853,CL854,CL855,CL856,CL857,CL858,CL859,CL860,CL861,CL862,CL863,CL864,CL865,CL866,CL867,CL868,CL869,CL870,CL871,CL872,CL873,CL874,CL875,CL876,CL877,CL878,CL879,CL880,CL881,CL882,CL883,CL884,CL885,CL886,CL887,CL888,CL889,CL890,CL891,CL892,CL893,CL894,CL895,CL896,CL897,CL898,CL899,CL900,CL901,CL902,CL903,CL904,CL905,CL906,CL907,CL908,CL909,CL910,CL911,CL912,CL913,CL914,CL915,CL916,CL917,CL918,CL919,CL920,CL921,CL922,CL923,CL924,CL925,CL926,CL927,CL928,CL929,CL930,CL931,CL932,CL933,CL934,CL935,CL936,CL937,CL938,CL939,CL940,CL941,CL942,CL943,CL944,CL945,CL946,CL947,CL948,CL949,CL950,CL951,CL952,CL953,CL954,CL955,CL956,CL957,CL958,CL959,CL960,CL961,CL962,CL963,CL964,CL965,CL966,CL967,CL968,CL969,CL970,CL971,CL972,CL973,CL974,CL975,CL976,CL977,CL978,CL979,CL980,CL981,CL982,CL983,CL984,CL985,CL986,CL987,CL988,CL989,CL990,CL991,CL992,CL993,CL994,CL995,CL996,CL997,CL998,CL999,CL1000,CL1001,CL1002,CL1003,CL1004,CL1005,CL1006,CL1007,CL1008,CL1009,CL1010,CL1011,CL1012,CL1013,CL1014,CL1015,CL1016,CL1017,CL1018,CL1019,CL1020,CL1021,CL1022,CL1023,CL1024,CL1025,CL1026,CL1027,CL1028,CL1029,CL1030,CL1031,CL1032,CL1033,CL1034,CL1035,CL1036,CL1037,CL1038,CL1039,CL1040,CL1041,CL1042,CL1043,CL1044,CL1045,CL1046,CL1047,CL1048,CL1049,CL1050,CL1051,CL1052,CL1053,CL1054,CL1055,CL1056,CL1057,CL1058,CL1059,CL1060,CL1061,CL1062,CL1063,CL1064,CL1065,CL1066,CL1067,CL1068,CL1069,CL1070,CL1071,CL1072,CL1073,CL1074,CL1075,CL1076,CL1077,CL1078,CL1079,CL1080,CL1081,CL1082,CL1083,CL1084,CL1085,CL1086,CL1087,CL1088,CL1089,CL1090,CL1091,CL1092,CL1093,CL1094,CL1095,CL1096,CL1097,CL1098,CL1099,CL1100,CL1101,CL1102,CL1103,CL1104,CL1105,CL1106,CL1107,CL1108,CL1109,CL1110,CL1111,CL1112,CL1113,CL1114,CL1115,CL1116,CL1117,CL1118,CL1119,CL1120,CL1121,CL1122,CL1123,CL1124,CL1125,CL1126,CL1127,CL1128,CL1129,CL1130,CL1131,CL1132,CL1133,CL1134,CL1135,CL1136,CL1137,CL1138,CL1139,CL1140,CL1141,CL1142,CL1143,CL1144,CL1145,CL1146,CL1147,CL1148,CL1149,CL1150,CL1151,CL1152,CL1153,CL1154,CL1155,CL1156,CL1157,CL1158,CL1159,CL1160,CL1161,CL1162,CL1163,CL1164,CL1165,CL1166,CL1167,CL1168,CL1169,CL1170,CL1171,CL1172,CL1173,CL1174,CL1175,CL1176,CL1177,CL1178,CL1179,CL1180,CL1181,CL1182,CL1183,CL1184,CL1185,CL1186,CL1187,CL1188,CL1189,CL1190,CL1191,CL1192,CL1193,CL1194,CL1195,CL1196,CL1197,CL1198,CL1199,CL1200,CL1201,CL1202,CL1203,CL1204,CL1205,CL1206,CL1207,CL1208,CL1209,CL1210,CL1211,CL1212,CL1213,CL1214,CL1215,CL1216,CL1217,CL1218,CL1219,CL1220,CL1221,CL1222,CL1223,CL1224,CL1225,CL1226,CL1227,CL1228,CL1229,CL1230,CL1231,CL1232,CL1233,CL1234,CL1235,CL1236,CL1237,CL1238,CL1239,CL1240,CL1241,CL1242,CL1243,CL1244,CL1245,CL1246,CL1247,CL1248,CL1249,CL1250,CL1251,CL1252,CL1253,CL1254,CL1255,CL1256,CL1257,CL1258,CL1259,CL1260,CL1261,CL1262,CL1263,CL1264,CL1265,CL1266,CL1267,CL1268,CL1269,CL1270,CL1271,CL1272,CL1273,CL1274,CL1275,CL1276,CL1277,CL1278,CL1279,CL1280,CL1281,CL1282,CL1283,CL1284,CL1285,CL1286,CL1287,CL1288,CL1289,CL1290,CL1291,CL1292,CL1293,CL1294,CL1295,CL1296,CL1297,CL1298,CL1299,CL1300,CL1301,CL1302,CL1303,CL1304,CL1305,CL1306,CL1307,CL1308,CL1309,CL1310,CL1311,CL1312,CL1313,CL1314,CL1315,CL1316,CL1317,CL1318,CL1319,CL1320,CL1321,CL1322,CL1323,CL1324,CL1325,CL1326,CL1327,CL1328,CL1329,CL1330,CL1331,CL1332,CL1333,CL1334,CL1335,CL1336,CL1337,CL1338,CL1339,CL1340,CL1341,CL1342,CL1343,CL1344,CL1345,CL1346,CL1347,CL1348,CL1349,CL1350,CL1351,CL1352,CL1353,CL1354,CL1355,CL1356,CL1357,CL1358,CL1359,CL1360,CL1361,CL1362,CL1363,CL1364,CL1365,CL1366,CL1367,CL1368,CL1369,CL1370,CL1371,CL1372,CL1373,CL1374,CL1375,CL1376,CL1377,CL1378,CL1379,CL1380,CL1381,CL1382,CL1383,CL1384,CL1385,CL1386,CL1387,CL1388,CL1389,CL1390,CL1391,CL1392,CL1393,CL1394,CL1395,CL1396,CL1397,CL1398,CL1399,CL1400,CL1401,CL1402,CL1403,CL1404,CL1405,CL1406,CL1407,CL1408,CL1409,CL1410,CL1411,CL1412,CL1413,CL1414,CL1415,CL1416,CL1417,CL1418,CL1419,CL1420,CL1421,CL1422,CL1423,CL1424,CL1425,CL1426,CL1427,CL1428,CL1429,CL1430,CL1431,CL1432,CL1433,CL1434,CL1435,CL1436,CL1437,CL1438,CL1439,CL1440,CL1441,CL1442,CL1443,CL1444,CL1445,CL1446,CL1447,CL1448,CL1449,CL1450,CL1451,CL1452,CL1453,CL1454,CL1455,CL1456,CL1457,CL1458,CL1459,CL1460,CL1461,CL1462,CL1463,CL1464,CL1465,CL1466,CL1467,CL1468,CL1469,CL1470,CL1471,CL1472,CL1473,CL1474,CL1475,CL1476,CL1477,CL1478,CL1479,CL1480,CL1481,CL1482,CL1483,CL1484,CL1485,CL1486,CL1487,CL1488,CL1489,CL1490,CL1491,CL1492,CL1493,CL1494,CL1495,CL1496,CL1497,CL1498,CL1499,CL1500,CL1501,CL1502,CL1503,CL1504,CL1505,CL1506,CL1507,CL1508,CL1509,CL1510,CL1511,CL1512,CL1513,CL1514,CL1515,CL1516,CL1517,CL1518,CL1519,CL1520,CL1521,CL1522,CL1523,CL1524,CL1525,CL1526,CL1527,CL1528,CL1529,CL1530,CL1531,CL1532,CL1533,CL1534,CL1535,CL1536,CL1537,CL1538,CL1539,CL1540,CL1541,CL1542,CL1543,CL1544,CL1545,CL1546,CL1547,CL1548,CL1549,CL1550,CL1551,CL1552,CL1553,CL1554,CL1555,CL1556,CL1557,CL1558,CL1559,CL1560,CL1561,CL1562,CL1563,CL1564,CL1565,CL1566,CL1567,CL1568,CL1569,CL1570,CL1571,CL1572,CL1573,CL1574,CL1575,CL1576,CL1577,CL1578,CL1579,CL1580,CL1581,CL1582,CL1583,CL1584,CL1585,CL1586,CL1587,CL1588,CL1589,CL1590,CL1591,CL1592,CL1593,CL1594,CL1595,CL1596,CL1597,CL1598,CL1599,CL1600,CL1601,CL1602,CL1603,CL1604,CL1605,CL1606,CL1607,CL1608,CL1609,CL1610,CL1611,CL1612,CL1613,CL1614,CL1615,CL1616,CL1617,CL1618,CL1619,CL1620,CL1621,CL1622,CL1623,CL1624,CL1625,CL1626,CL1627,CL1628,CL1629,CL1630,CL1631,CL1632,CL1633,CL1634,CL1635,CL1636,CL1637,CL1638,CL1639,CL1640,CL1641,CL1642,CL1643,CL1644,CL1645,CL1646,CL1647,CL1648,CL1649,CL1650,CL1651,CL1652,CL1653,CL1654,CL1655,CL1656,CL1657,CL1658,CL1659,CL1660,CL1661,CL1662,CL1663,CL1664,CL1665,CL1666,CL1667,CL1668,CL1669,CL1670,CL1671,CL1672,CL1673,CL1674,CL1675,CL1676,CL1677,CL1678,CL1679,CL1680,CL1681,CL1682,CL1683,CL1684,CL1685,CL1686,CL1687,CL1688,CL1689,CL1690,CL1691,CL1692,CL1693,CL1694,CL1695,CL1696,CL1697,CL1698,CL1699,CL1700,CL1701,CL1702,CL1703,CL1704,CL1705,CL1706,CL1707,CL1708,CL1709,CL1710,CL1711,CL1712,CL1713,CL1714,CL1715,CL1716,CL1717,CL1718,CL1719,CL1720,CL1721,CL1722,CL1723,CL1724,CL1725,CL1726,CL1727,CL1728,CL1729,CL1730,CL1731,CL1732,CL1733,CL1734,CL1735,CL1736,CL1737,CL1738,CL1739,CL1740,CL1741,CL1742,CL1743,CL1744,CL1745,CL1746,CL1747,CL1748,CL1749,CL1750,CL1751,CL1752,CL1753,CL1754,CL1755,CL1756,CL1757,CL1758,CL1759,CL1760,CL1761,CL1762,CL1763,CL1764,CL1765,CL1766,CL1767,CL1768,CL1769,CL1770,CL1771,CL1772,CL1773,CL1774,CL1775,CL1776,CL1777,CL1778,CL1779,CL1780,CL1781,CL1782,CL1783,CL1784,CL1785,CL1786,CL1787,CL1788,CL1789,CL1790,CL1791,CL1792,CL1793,CL1794,CL1795,CL1796,CL1797,CL1798,CL1799,CL1800,CL1801,CL1802,CL1803,CL1804,CL1805,CL1806,CL1807,CL1808,CL1809,CL1810,CL1811,CL1812,CL1813,CL1814,CL1815,CL1816,CL1817,CL1818,CL1819,CL1820,CL1821,CL1822,CL1823,CL1824,CL1825,CL1826,CL1827,CL1828,CL1829,CL1830,CL1831,CL1832,CL1833,CL1834,CL1835,CL1836,CL1837,CL1838,CL1839,CL1840,CL1841,CL1842,CL1843,CL1844,CL1845,CL1846,CL1847,CL1848,CL1849,CL1850,CL1851,CL1852,CL1853,CL1854,CL1855,CL1856,CL1857,CL1858,CL1859,CL1860,CL1861,CL1862,CL1863,CL1864,CL1865,CL1866,CL1867,CL1868,CL1869,CL1870,CL1871,CL1872,CL1873,CL1874,CL1875,CL1876,CL1877,CL1878,CL1879,CL1880,CL1881,CL1882,CL1883,CL1884,CL1885,CL1886,CL1887,CL1888,CL1889,CL1890,CL1891,CL1892,CL1893,CL1894,CL1895,CL1896,CL1897,CL1898,CL1899,CL1900,CL1901,CL1902,CL1903,CL1904,CL1905,CL1906,CL1907,CL1908,CL1909,CL1910,CL1911,CL1912,CL1913,CL1914,CL1915,CL1916,CL1917,CL1918,CL1919,CL1920,CL1921,CL1922,CL1923,CL1924,CL1925,CL1926,CL1927,CL1928,CL1929,CL1930,CL1931,CL1932,CL1933,CL1934,CL1935,CL1936,CL1937,CL1938,CL1939,CL1940,CL1941,CL1942,CL1943,CL1944,CL1945,CL1946,CL1947,CL1948,CL1949,CL1950,CL1951,CL1952,CL1953,CL1954,CL1955,CL1956,CL1957,CL1958,CL1959,CL1960,CL1961,CL1962,CL1963,CL1964,CL1965,CL1966,CL1967,CL1968,CL1969,CL1970,CL1971,CL1972,CL1973,CL1974,CL1975,CL1976,CL1977,CL1978,CL1979,CL1980,CL1981,CL1982,CL1983,CL1984,CL1985,CL1986,CL1987,CL1988,CL1989,CL1990,CL1991,CL1992,CL1993,CL1994,CL1995,CL1996,CL1997,CL1998,CL1999,CL2000,CL2001,CL2002,CL2003,CL2004,CL2005,CL2006,CL2007,CL2008,CL2009,CL2010,CL2011,CL2012,CL2013,CL2014,CL2015,CL2016,CL2017,CL2018,CL2019,CL2020,CL2021,CL2022,CL2023,CL2024,CL2025,CL2026,CL2027,CL2028,CL2029,CL2030,CL2031,CL2032,CL2033,CL2034,CL2035,CL2036,CL2037,CL2038,CL2039,CL2040,CL2041,CL2042,CL2043,CL2044,CL2045,CL2046,CL2047,CL2048,CL2049,CL2050,CL2051,CL2052,CL2053,CL2054,CL2055,CL2056,CL2057,CL2058,CL2059,CL2060,CL2061,CL2062,CL2063,CL2064,CL2065,CL2066,CL2067,CL2068,CL2069,CL2070,CL2071,CL2072,CL2073,CL2074,CL2075,CL2076,CL2077,CL2078,CL2079,CL2080,CL2081,CL2082,CL2083,CL2084,CL2085,CL2086,CL2087,CL2088,CL2089,CL2090,CL2091,CL2092,CL2093,CL2094,CL2095,CL2096,CL2097,CL2098,CL2099,CL2100,CL2101,CL2102,CL2103,CL2104,CL2105,CL2106,CL2107,CL2108,CL2109,CL2110,CL2111,CL2112,CL2113,CL2114,CL2115,CL2116,CL2117,CL2118,CL2119,CL2120,CL2121,CL2122,CL2123,CL2124,CL2125,CL2126,CL2127,CL2128,CL2129,CL2130,CL2131,CL2132,CL2133,CL2134,CL2135,CL2136,CL2137,CL2138,CL2139,CL2140,CL2141,CL2142,CL2143,CL2144,CL2145,CL2146
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COMMON/GTRY2/VAME(4615),JA(183),JB(184),NNZ,JC,NUMNP,NUMEL,NELDOF,	
1CONN(128,15)	T-025
COMMON/GTRY1/X(505),P(505),Z(505)	T-030
COMMON/XOCAL/XX(15),YY(15),ZZ(15),PPSI(15)	T-035
DIMENSION Y(7,1)	T-045
COMMON/TEMP/C3(183),C6(15)	T-040
COMMON/VGGGMAT/GGG(15,15)	T-055
COMMON/SY SMT3/BIGH(4615)	
COMMON/GAUSS/WL(7),WS(5),S(5),CL1(7),CL2(7),CL3(7),FN(15,21),DL1(1	T-065
15,21),DL2(15,21),DL3(15,21),DS(15,21),DETJ(21)	T-070
DO 300 I=1,JC	
BIGH(I)=0.0	T-085
CONTINUE	T-090
NNAAA=7	T-100
NNSS=3	
DO 800 L=1,NUMEL	T-120
DO 600 I=1,15	T-125
DO 610 J=1,15	T-130
GGG(I,J)=0.0	T-135
CONTINUE	T-140
CONTINUE	T-145
DO 805 JJ=1,15	T-155
NN=CONN(L,JJ)	T-160
XX(JJ)=X(NN)	T-165
YY(JJ)=P(NN)	T-170
ZZ(JJ)=Z(NN)	T-175
IF(NN.GT.NNZ) GO TO 803	
C6(JJ)=C3(NN)	T-180
PPSI(JJ)=Y(1,NN)	T-185
GO TO 805	T-190
PPSI(JJ)=0.0	
C6(JJ)=0.0	T-195
CONTINUE	T-200
A5=2.0/(ZZ(1)-ZZ(15))	T-205
DETJJ=0.0	T-210
DO 802 K=1,NNSS	T-215
DETJL=0.0	T-220
DO 798 M=1,NNAAA	T-225
IF(K.EQ.1) GO TO 111	T-230
IF(K.EQ.2) GO TO 222	T-235
IF(K.EQ.3) GO TO 333	T-240
GO TO 788	T-245
111 N=M	T-250
GO TO 788	T-255
222 N=M+7	T-260
GO TO 788	T-265
333 N=M+14	T-270
788 CONTINUE	





```

T1=DL1(1,N)*XX(1)+DL1(2,N)*XX(2)-DL3(4,N)*XX(4)-DL3(5,N)*XX(5)
T2=DL1(7,N)*XX(7)-DL3(9,N)*XX(9)+DL1(10,N)*XX(10)+DL1(11,N)*XX(11)
T3=(DL1(6,N)-DL3(6,N))*XX(6)-DL3(13,N)*XX(13)-DL3(14,N)*XX(14)
T4=(DL1(15,N)-DL3(15,N))*XX(15)
DJ1=T1+T2+T3+T4
D1=DL1(1,N)*YY(1)+DL1(2,N)*YY(2)-DL3(4,N)*YY(4)-DL3(5,N)*YY(5)
D2=DL1(7,N)*YY(7)-DL3(9,N)*YY(9)+DL1(10,N)*YY(10)+DL1(11,N)*YY(11)
D3=(DL1(6,N)-DL3(6,N))*YY(6)-DL3(13,N)*YY(13)-DL3(14,N)*YY(14)
D4=(DL1(15,N)-DL3(15,N))*YY(15)
DJ12=D1+D2+D3+D4
T5=DL2(2,N)*XX(2)+DL2(3,N)*XX(3)-DL3(5,N)*XX(5)-DL3(6,N)*XX(6)
T6=DL2(8,N)*XX(8)-DL3(9,N)*XX(9)+(DL2(4,N)-DL3(4,N))*XX(4)
T7=-DL2(11,N)*XX(11)+DL2(12,N)*XX(12)+(DL2(13,N)-DL3(13,N))*XX(13)
T8=-DL3(14,N)*XX(14)-DL3(15,N)*XX(15)
DJ21=T5+T6+T7+T8
D5=DL2(2,N)*YY(2)+DL2(3,N)*YY(3)-DL3(5,N)*YY(5)-DL3(6,N)*YY(6)
D6=DL2(8,N)*YY(8)-DL3(9,N)*YY(9)+(DL2(4,N)-DL3(4,N))*YY(4)
D7=DL2(11,N)*YY(11)+DL2(12,N)*YY(12)+(DL2(13,N)-DL3(13,N))*YY(13)
D8=-DL3(14,N)*YY(14)-DL3(15,N)*YY(15)
DJ22=D5+D6+D7+D8
DETJ(N)=DJ11*DJ22-DJ12*DJ21
DETJ1=DETJ(N)*WL(M)+DETJ1
CONTINUE
798 DETJJ=DETJJ+WS(K)*DETJ1
802 CONTINUE
DO 700 J=1,15
DO 710 I=1,15
FHH=0.0
DO 806 K=1,NNSSS
FH=0.0
DO 799 M=1,NNAAA
TT1=0.0
TT2=0.0
TT3=0.0
IF(K.EQ.1) GO TO 1111
IF(K.EQ.2) GO TO 2222
IF(K.EQ.3) GO TO 3333
GO TO 7888
1111 N=M
GO TO 7888
2222 N=M+7
GO TO 7888
3333 N=M+14
7888 CONTINUE
1(3)+FN(4,N)*PPSI(4)*C6(4)+FN(5,N)*PPSI(5)*C6(5)+FN(6,N)*PPSI(6)*C6
1(6)+FN(7,N)*PPSI(7)*C6(7)+FN(8,N)*PPSI(8)*C6(8)
TT2=FN(9,N)*PPSI(9)*C6(9)+FN(10,N)*PPSI(10)*C6(10)+FN(11,N)*PPSI(11

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11)*C6(11)+FN(12,N)*PPSI(12)*C6(12)+FN(13,N)*PPSI(13)*C6(13)+FN(14,
1N)*PPSI(14)*C6(14)+FN(15,N)*PPSI(15)*C6(15)
TT3=TT1+TT2+1.0
FH=FH+WL(M)*FN(J,N)*FN(I,N)*ALOG(TT3)*DETJ(N)
799 CONTINUE
FHH=WS(K)*FH+FHH
806 CONTINUE
GGG(J,I)=FHH/A5
710 CONTINUE
700 CONTINUE
C INSERT INTO SYSTEM MATRIX
DO 20 K=1,15
KK=CONN(L,K)
IF(KK.GT.NNZ) GO TO 20
KKK=JA(KK)
LLL=JB(KK)-1
DO 10 I=1,15
II=CONN(L,I)
DO 91 M=1,KKK
MM=LLL+M
KKM=NAME(MM)
IF(II.EQ.KKM)GO TO 92
91 CONTINUE
GO TO 10
92 CONTINUE
BIGH(MM)=BIGH(MM)+GGG(K,I)
10 CONTINUE
20 CONTINUE
800 CONTINUE
RETURN
END
C-----
C THE CUMULATIVE CONTRIBUTION OF THE DELAYED NEUTRON FLUX IS
C CALCULATED BY THIS SUBROUTINE.
C-----
C
SUBROUTINE SUMT(Y,T,H)
DIMENSION Y(7,1)
COMMON/DELAY/SUM(183),PSI(93)
Q=0.4349710
TC=-Q*H
DO 100 I=1,93
SUM(I)=SUM(I)*EXP(-Q*H)+0.50*H*(PSI(I)*EXP(TC)+Y(1,I))
100 CONTINUE
RETURN
END
S-005
S-010
S-030
S-040
S-050
S-055
S-060
T-580
T-585
T-590
T-595
T-600
T-605
T-610
T-615
T-620
T-625
T-630
T-635
T-640
T-645
T-650
T-655
T-660
T-665
T-670
T-675
T-680
T-685
T-690
T-695
T-700
T-705
T-710
T-715
T-720

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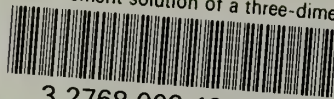
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